

.(25 pts.) Consider

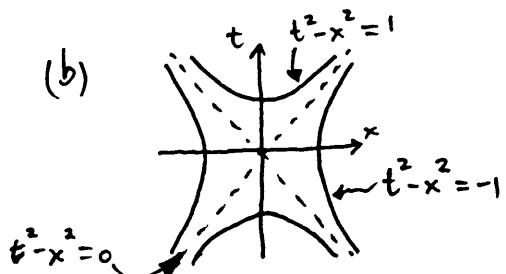
$$(*) \quad tu_x + xu_t = 0.$$

- (a) Find the characteristic curves of (*).
- (b) Sketch and identify two characteristic curves of (*).
- (c) Write the general solution of (*).
- (d) Determine the particular solution of (*) that satisfies the

auxiliary condition $u(x, 0) = e^{-x^2}$ for $-\infty < x < \infty$.

- (e) In what region of the xt -plane is the solution in part (d) uniquely determined?

$$(a) \quad \frac{dt}{dx} = \frac{x}{t} \Rightarrow t dt = x dx \Rightarrow \frac{1}{2}t^2 = \frac{1}{2}x^2 + c_1 \Rightarrow \boxed{t^2 - x^2 = c}$$

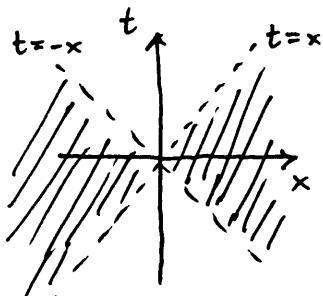


$$(c) \quad \boxed{u(x, t) = f(t^2 - x^2)} \text{ where } f \text{ is any } C^1 \text{-function of a single real variable.}$$

$$(d) \quad e^{-x^2} = u(x, 0) = f(-x^2) \text{ for all real } x \Rightarrow f(w) = e^w \text{ for all } w \leq 0.$$

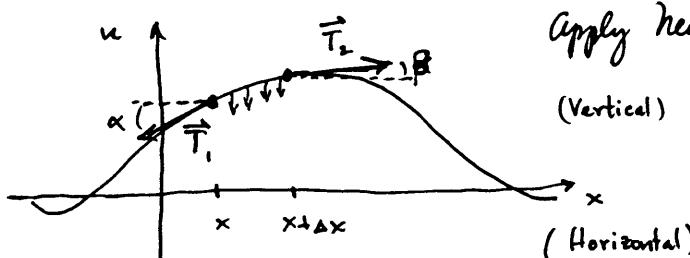
$$\therefore \boxed{u(x, t) = e^{t^2 - x^2}}$$

- (e) The solution in part (d) is uniquely determined in the region of the xt -plane given by $\{(x, t) \in \mathbb{R}^2 : t^2 - x^2 \leq 0\}$



2.(25 pts.) Carefully derive the equation of motion of a string in a medium which resists motion with a force proportional to the velocity of the motion.

Let $u(x,t)$ denote the vertical displacement of the string at position x and time t . Fix $t > 0$ and consider the segment of string between x and $x+\Delta x$.



Apply Newton's second law of motion to this segment:

$$(Vertical) \quad \int_x^{x+\Delta x} \rho(\xi) u_{tt}(\xi, t) d\xi \stackrel{(1)}{=} |\vec{T}_2| \sin(\beta) - |\vec{T}_1| \sin(\alpha) - \int_x^{x+\Delta x} r u_t(\xi, t) d\xi$$

$$0 \stackrel{(2)}{=} |\vec{T}_2| \cos(\beta) - |\vec{T}_1| \cos(\alpha)$$

(Assume motion is upward for sake of argument.)

(Here $\rho = \rho(x)$ denotes the mass density of the string at x , r is a proportionality constant, and

\vec{T}_1 and \vec{T}_2 are the tensions exerted at x and $x+\Delta x$, respectively.) Since $u_x(x, t) = \tan(\alpha)$,

it follows that $\sin(\alpha) \stackrel{(3)}{=} \frac{u_x(x, t)}{\sqrt{1+u_x^2(x, t)}}$ and $\cos(\alpha) \stackrel{(4)}{=} \frac{1}{\sqrt{1+u_x^2(x, t)}}$. Similarly $\sin(\beta) \stackrel{(5)}{=}$

$\frac{u_x(x+\Delta x, t)}{\sqrt{1+u_x^2(x+\Delta x, t)}}$ and $\cos(\beta) \stackrel{(6)}{=} \frac{1}{\sqrt{1+u_x^2(x+\Delta x, t)}}$. Substituting from (4) and (6) into (2) gives

$$|\vec{T}_2| = \frac{|\vec{T}_1| \cos(\alpha)}{\cos(\beta)} \stackrel{(7)}{=} \frac{|\vec{T}_1| \sqrt{1+u_x^2(x+\Delta x, t)}}{\sqrt{1+u_x^2(x, t)}}. \text{ Substituting (3), (5), and (7) into (1) yields}$$

$$\int_x^{x+\Delta x} \rho(\xi) u_{tt}(\xi, t) d\xi \stackrel{(8)}{=} \frac{|\vec{T}_1|}{\sqrt{1+u_x^2(x, t)}} \left[u_x(x+\Delta x, t) - u_x(x, t) \right] - \int_x^{x+\Delta x} r u_t(\xi, t) d\xi. \text{ Dividing (8) by}$$

Δx and letting $\Delta x \rightarrow 0$ produces

$$\boxed{\rho(x) u_{tt}(x, t) = \frac{|\vec{T}_1|}{\sqrt{1+u_x^2(x, t)}} u_{xx}(x, t) - r u_t(x, t)}.$$

For "small" vibrations of the string $u_x^2 \ll 1$ so $\sqrt{1+u_x^2(x, t)} \approx 1 \approx \sqrt{1+u_x^2(x+\Delta x, t)}$

and $|\vec{T}_1| = |\vec{T}_2| = \text{constant} = T_0$. Also, for homogeneous string $\rho(x) = \text{constant} = \rho$,

so the equation of motion becomes

$$\boxed{\rho u_{tt} - T_0 u_{xx} + r u_t = 0}.$$

3. (25 pts.) Consider

$$(+) \quad u_{xx} + 4u_{yy} - 4u_{xy} + u = (x - 2y)^2.$$

- (a) Classify the order and type (nonlinear, linear, homogeneous, inhomogeneous, elliptic, hyperbolic, parabolic) of (+).
 (b) Find, if possible, the general solution of (+) in the xy -plane.

(a) 2nd-order, linear, inhomogeneous. $B^2 - 4AC = (-4)^2 - 4(1)(1) = 0$ parabolic.

$$(b) \left(\frac{\partial^2}{\partial x^2} - 4 \frac{\partial^2}{\partial x \partial y} + 4 \frac{\partial^2}{\partial y^2} \right) u + u = (x - 2y)^2.$$

$$\left(\frac{\partial}{\partial x} - 2 \frac{\partial}{\partial y} \right) \left(\frac{\partial}{\partial x} - 2 \frac{\partial}{\partial y} \right) u + u = (x - 2y)^2.$$

Let $\begin{cases} \xi = 2x + y, \\ \eta = x - 2y. \end{cases}$ Then $\frac{\partial}{\partial x} = \frac{2}{\partial \xi} + \frac{1}{\partial \eta}$ and $\frac{\partial}{\partial y} = \frac{1}{\partial \xi} - 2 \frac{\partial}{\partial \eta}$
 so $\frac{\partial}{\partial x} - 2 \frac{\partial}{\partial y} = 5 \frac{\partial}{\partial \eta}.$

Hence (+) is equivalent to $25 \frac{\partial^2 u}{\partial \eta^2} + u = \eta^2 \Rightarrow \frac{\partial^2 u}{\partial \eta^2} + \frac{1}{25} u = \frac{\eta^2}{25}. \quad (++)$

$$\therefore u(\xi, \eta) = \underbrace{c_1(\xi) \cos\left(\frac{\eta}{5}\right) + c_2(\xi) \sin\left(\frac{\eta}{5}\right)}_{\text{general solution to } u_{\eta\eta} + \frac{1}{25}u = 0} + u_p \quad \text{a particular solution of (++)}.$$

We use undetermined coefficients to find u_p . A trial solution is $u_p = A\eta^2 + B\eta + C$

Then $\frac{\partial u_p}{\partial \eta} = 2A\eta + B$ and $\frac{\partial^2 u_p}{\partial \eta^2} = 2A$. Substituting in (++) gives

$$2A + \frac{1}{25}(A\eta^2 + B\eta + C) = \frac{\eta^2}{25}$$

$$\Rightarrow \frac{A}{25} = \frac{1}{25}, \quad \frac{B}{25} = 0, \quad \text{and} \quad 2A + \frac{C}{25} = 0, \quad \text{i.e.} \quad A = 1, \quad B = 0, \quad C = -50.$$

Thus $u_p = \eta^2 - 50$, and hence $u(\xi, \eta) = c_1(\xi) \cos\left(\frac{\eta}{5}\right) + c_2(\xi) \sin\left(\frac{\eta}{5}\right) + \eta^2 - 50$.

Therefore

$$u(x, y) = f(2x+y) \cos\left(\frac{x-2y}{5}\right) + g(2x+y) \sin\left(\frac{x-2y}{5}\right) + (x-2y)^2 - 50$$

where f and g are any C^2 -functions of a single real variable.

4.(25 pts.) (a) Write (no proof or derivation required) and simplify an expression for the solution to

$$(*) \quad u_{tt} - 4u_{xx} = 0 \quad \text{for } -\infty < x < \infty, 0 < t < \infty,$$

subject to the initial conditions

$$u(x,0) = e^{-x^2} \quad \text{and} \quad u_t(x,0) = 4xe^{-x^2} \quad \text{for } -\infty < x < \infty.$$

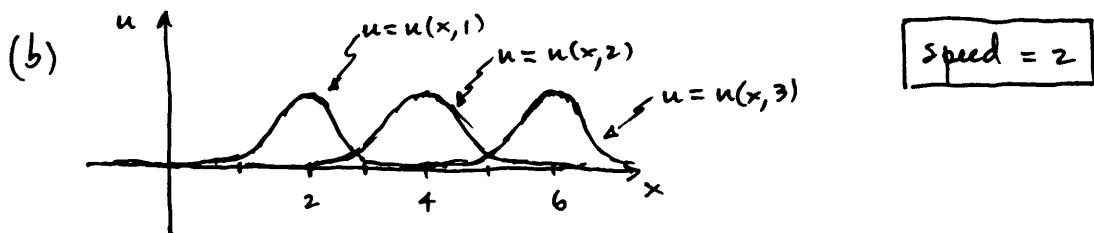
(b) Sketch profiles of the solution in part (a), $u = u(x,t)$, at $t = 1, 2, 3$ in order to illustrate that the solution is a wave traveling to the right along the x -axis. What is its speed?

(c) Derive a general nontrivial relation between ϕ and ψ which will produce a solution to $(*)$ in the xt -plane satisfying

$$u(x,0) = \phi(x) \quad \text{and} \quad u_t(x,0) = \psi(x) \quad \text{for } -\infty < x < \infty$$

and such that u consists solely of a wave traveling to the right along the x -axis.

$$\begin{aligned} (a) \quad u(x,t) &= \frac{1}{2} \left[e^{-(x-2t)^2} + e^{-(x+2t)^2} \right] + \frac{1}{4} \int_{x-2t}^{x+2t} 4se^{-s^2} ds \quad \leftarrow \begin{array}{l} \text{Let } w = -s^2. \text{ Then } dw = -2sds \\ \text{so } \int 4se^{-s^2} ds = \int -2e^w dw = -2e^w \\ = -2e^{-s^2} \end{array} \\ &= \frac{1}{2} \left[e^{-(x-2t)^2} + e^{-(x+2t)^2} \right] - \frac{1}{2} \left[e^{-(x+2t)^2} - e^{-(x-2t)^2} \right] \\ &= \boxed{e^{-(x-2t)^2}} \end{aligned}$$



$$(c) \quad u(x,t) = \frac{1}{2} \left[\phi(x-2t) + \phi(x+2t) \right] + \frac{1}{4} \int_{x-2t}^{x+2t} \psi(s) ds$$

We need $\frac{1}{2}\phi(x+2t) + \frac{1}{4} \int_0^{x+2t} \psi(s) ds = 0$ for all x and t in order for

$u = u(x,t)$ to consist solely of a wave traveling to the right along the x -axis.

I.e. $\phi(z) = -\frac{1}{2} \int_0^z \psi(s) ds$ for all real $z \Rightarrow \boxed{\phi'(z) = -\frac{1}{2} \psi(z) \text{ for all real } z}.$