

1. (30 pts.) Find a solution to

$$u_t - u_{xx} = 0 \quad \text{in } -\infty < x < \infty, 0 < t < \infty,$$

satisfying

$$u(x, 0) = e^{-x^2} \quad \text{for } -\infty < x < \infty.$$

$$u(x, t) = \int_{-\infty}^{\infty} \frac{e^{-\frac{(x-y)^2}{4t}}}{\sqrt{4\pi t}} e^{-y^2} dy$$

$$= \int_{-\infty}^{\infty} \frac{e^{-\frac{x^2-2xy+y^2}{4t}}}{\sqrt{4\pi t}} \cdot e^{-y^2} dy$$

$$= \int_{-\infty}^{\infty} e^{-\frac{(4t+1)y^2 - 2y\sqrt{4t+1}\left(\frac{x}{\sqrt{4t+1}}\right) + \frac{x^2}{4t+1}}{4t}} \cdot e^{-\frac{-x^2 + \frac{x^2}{4t+1}}{4t}} dy$$

$$= e^{-\frac{(4t+1)x^2 + x^2}{4t(4t+1)}} \int_{-\infty}^{\infty} e^{-\frac{(\sqrt{4t+1}y - \frac{x}{\sqrt{4t+1}})^2}{4t}} dy$$

$$= \frac{e^{-\frac{x^2}{4t+1}}}{\sqrt{\pi} \cdot \sqrt{4t+1}} \int_{-\infty}^{\infty} e^{-p^2} dp$$

$$= \boxed{\frac{e^{-\frac{x^2}{4t+1}}}{\sqrt{4t+1}}}$$

Complete square and
"compensate".

Let $p = \frac{\sqrt{4t+1}y - \frac{x}{\sqrt{4t+1}}}{\sqrt{4t}}$
Then $dp = \frac{\sqrt{4t+1}}{\sqrt{4t}} dy$

Valid for $t > -\frac{1}{4}$ and $-\infty < x < \infty$.

2. (30 pts.) (a) Let $f = f(x)$ be an absolutely integrable function on $(-\infty, \infty)$, let c be a real constant, and define the translate of f by c according to the formula

$$f_c(x) = f(x-c) \quad \text{for all } -\infty < x < \infty.$$

Show that the Fourier transform of f_c satisfies the relation

$$\mathcal{F}(f_c)(\xi) = e^{-i\xi c} \mathcal{F}(f)(\xi) \quad \text{for all } -\infty < \xi < \infty.$$

(b) Use the Fourier transform method to solve

$$u_t - u_{xx} + u_x = 0 \quad \text{in } -\infty < x < \infty, \quad 0 < t < \infty,$$

subject to the initial condition

$$u(x, 0) = \phi(x) \quad \text{for all } -\infty < x < \infty.$$

$$\begin{aligned} (a) \quad \mathcal{F}(f_c)(\xi) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f_c(x) e^{-i\xi x} dx = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x-c) e^{-i\xi x} dx \leftarrow \begin{cases} \text{Let } y = x-c. \\ \text{Then } dy = dx \end{cases} \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(y) e^{-i\xi(y+c)} dy = e^{-i\xi c} \cdot \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(y) e^{-i\xi y} dy = e^{-i\xi c} \mathcal{F}(f)(\xi). \end{aligned}$$

$$(b) \quad \mathcal{F}(u_t - u_{xx} + u_x)(\xi) = \mathcal{F}(0)(\xi)$$

$$\frac{d}{dt} \mathcal{F}(u)(\xi) - (i\xi)^2 \mathcal{F}(u)(\xi) + i\xi \mathcal{F}(u)(\xi) = 0$$

$$\frac{d}{dt} \mathcal{F}(u)(\xi) + (i\xi^2 + i\xi) \mathcal{F}(u)(\xi) = 0 \quad \leftarrow \text{linear, first order DE. An integrating factor is } \mu(t) = e^{\int (i\xi^2 + i\xi) dt} = e^{(i\xi^2 + i\xi)t}.$$

$$\therefore e^{(i\xi^2 + i\xi)t} \frac{d}{dt} \mathcal{F}(u)(\xi) + (i\xi^2 + i\xi)e^{(i\xi^2 + i\xi)t} \mathcal{F}(u)(\xi) = 0$$

$$\frac{d}{dt} \left(e^{(i\xi^2 + i\xi)t} \mathcal{F}(u)(\xi) \right) = 0$$

$$e^{(i\xi^2 + i\xi)t} \mathcal{F}(u)(\xi) = c(\xi)$$

$$\mathcal{F}(u)(\xi) = c(\xi) e^{-i\xi t}$$

$$\therefore \mathcal{F}(\varphi)(\xi) = \mathcal{F}(u(\cdot, 0))(\xi) = c(\xi) e^{-i\xi \cdot 0} = c(\xi)$$

$$\text{so } \mathcal{F}(u)(\xi) = \mathcal{F}(\varphi)(\xi) e^{-i\xi t} \cdot e$$

Table entry I:

$$\mathcal{F}\left(e^{-at}\right)(\xi) = \frac{e^{-\frac{a^2}{4}\xi^2}}{\sqrt{2\pi}}$$

Take $a = \frac{1}{4t}$ and rearrange:

$$\mathcal{F}\left(\frac{1}{\sqrt{2t}} e^{-\frac{(c+\xi)^2}{4t}}\right)(\xi) = e^{-\frac{\xi^2}{4t}}$$

Therefore $\mathcal{F}(u)(\xi) = \mathcal{F}(\varphi)(\xi) \mathcal{F}\left(\frac{1}{\sqrt{2t}} e^{-\frac{(\cdot)^2}{4t}}\right)(\xi) e^{-i\xi t}$

$= \frac{1}{\sqrt{2\pi}} \mathcal{F}\left(\varphi * \frac{1}{\sqrt{2t}} e^{-\frac{(\cdot)^2}{4t}}\right)(\xi) e^{-i\xi t}$

$= \mathcal{F}\left(\varphi * \frac{1}{\sqrt{4\pi t}} e^{-\frac{(\cdot)^2}{4t}}\right)(\xi) e^{-i\xi t}$

\downarrow $\mathcal{F}\left[\left(\varphi * \frac{1}{\sqrt{4\pi t}} e^{-\frac{(\cdot)^2}{4t}}\right)_t\right](\xi)$

$\mathcal{F}(f*g)(\xi) = \sqrt{2\pi} \mathcal{F}(f)(\xi) \mathcal{F}(g)(\xi)$

Part (a) with $c = t$

Therefore the inversion theorem implies

$$u(x,t) = \left(\varphi * \frac{1}{\sqrt{4\pi t}} e^{-\frac{(\cdot)^2}{4t}} \right)_t(x) \quad \text{for all } -\infty < x < \infty \text{ and} \\ \text{all } 0 < t < \infty.$$

$$\text{I.e. } u(x,t) = \left(\varphi * \frac{1}{\sqrt{4\pi t}} e^{-\frac{(x-t)^2}{4t}} \right) (x-t)$$

$$= \boxed{\int_{-\infty}^{\infty} \frac{e^{-\frac{(x-t-y)^2}{4t}}}{\sqrt{4\pi t}} \varphi(y) dy}$$

for all $-\infty < x < \infty$ and
all $0 < t < \infty$.

3. (40 pts.) (a) Find a solution to

$$(1) \quad u_{tt} - u_{xx} + 2u_t = 0 \quad \text{in } 0 < x < \pi, 0 < t < \infty,$$

satisfying

$$(2) \quad u(0, t) = 0 = u(\pi, t) \quad \text{for } t \geq 0,$$

$$(3) \quad u(x, 0) = 0 \quad \text{for } 0 \leq x \leq \pi,$$

$$(4) \quad u_t(x, 0) = \frac{3}{4} \sin(x) - \frac{1}{4} \sin(3x) \quad \text{for } 0 \leq x \leq \pi.$$

(b) If $u = u(x, t)$ satisfies (1)-(2), show that its energy

$$E(t) = \frac{1}{2} \int_0^\pi [u_t^2(x, t) + u_x^2(x, t)] dx$$

is nonincreasing on $0 \leq t < \infty$.

(c) Is there only one solution to the problem in part (a)? Why or why not?

(a) (via separation of variables) We seek nontrivial solutions of (1)-(2)-(3) of the form $u(x, t) = \Xi(x)T(t)$. Substituting in (1) gives $\Xi''(x)T'(t) - \Xi''(x)T(t) + 2\Xi(x)T'(t) = 0$, and hence $\frac{T''(t) + 2T'(t)}{T(t)} = \frac{\Xi''(x)}{\Xi(x)} = \text{constant} = -\lambda$. Substituting in (2)-(3) gives $\Xi(0)T(t) = 0 = \Xi(\pi)T(t)$ for $t \geq 0$ and $\Xi(x)T(0) = 0$ for $0 \leq x \leq \pi$. Thus:

$$\begin{cases} \Xi''(x) + \lambda \Xi(x) = 0 \text{ for } 0 < x < \pi, \quad \Xi(0) = 0 = \Xi(\pi); \\ T''(t) + 2T'(t) + \lambda T(t) = 0 \text{ for } t > 0, \quad T(0) = 0. \end{cases}$$

The eigenvalues are $\lambda_n = n^2$ ($n=1, 2, 3, \dots$) and the corresponding eigenfunctions are $\Xi_n(x) = \sin(nx)$ ($n=1, 2, 3, \dots$). The functions of t satisfy

$$T_n''(t) + 2T_n'(t) + n^2 T_n(t) = 0. \quad \text{Then } T_n(t) = e^{\alpha_n t} \text{ leads to } \alpha_n^2 + 2\alpha_n + n^2 = 0$$

$$\Rightarrow \alpha_n = \frac{-2 \pm \sqrt{4 - 4n^2}}{2} \quad (n=1, 2, \dots)$$

$$\therefore \alpha_1 = -1 \text{ and } \alpha_n = -1 \pm i\sqrt{n^2 - 1} \quad \text{for } n=2, 3, \dots$$

$$\text{and } T_1(t) = a_1 e^{-t} + b_1 t e^{-t}, \quad T_n(t) = e^{-t} \left[a_n \cos(t\sqrt{n^2 - 1}) + b_n \sin(t\sqrt{n^2 - 1}) \right] \quad (n \geq 2)$$

The initial condition $T_n(0) = 0$ implies $a_1 = 0$ ($n=1$) and $a_n = 0$ ($n \geq 2$).

Thus (upto a constant factor) $u_1(x, t) = t e^{-t} \sin(x)$ and

$$u_n(x, t) = e^{-t} \sin(t\sqrt{n^2 - 1}) \sin(nx) \quad \text{for } n \geq 2. \quad \text{Consequently,}$$

$$u(x,t) = b_1 t e^{-t} \sin(x) + \sum_{n=1}^{\infty} b_n e^{-t} \sin(t\sqrt{n^2-1}) \sin(nx)$$

is a (formal) solution of (1)-(2)-(3) for any choice of constants b_1, b_2, \dots . Differentiating with respect to t (holding x fixed) yields

$$u_t(x,t) = b_1 (1 e^{-t} - t e^{-t}) \sin(x) + \sum_{n=1}^{\infty} b_n (\sqrt{n^2-1} \cos(t\sqrt{n^2-1}) - \sin(t\sqrt{n^2-1})) e^{-t} \sin(nx)$$

so we want to choose the constants b_1, b_2, \dots so that, for all $0 \leq x \leq \pi$,

$$\frac{3}{4} \sin(x) - \frac{1}{4} \sin(3x) = u_t(x,0) = b_1 \sin(x) + \sum_{n=1}^{\infty} \sqrt{n^2-1} b_n \sin(nx).$$

By inspection, $b_1 = \frac{3}{4}$, $\sqrt{9-1} b_3 = -\frac{1}{4}$ (i.e. $b_3 = -\frac{1}{8\sqrt{2}}$), and all other $b_n = 0$ suffices. That is,

$$u(x,t) = \frac{3}{4} t e^{-t} \sin(x) - \frac{1}{8\sqrt{2}} e^{-t} \sin(2t\sqrt{2}) \sin(3x)$$

$$\begin{aligned}
 (b) \quad E'(t) &= \frac{d}{dt} \left\{ \frac{1}{2} \int_0^\pi [u_t^2(x,t) + u_x^2(x,t)] dx \right\} = \frac{1}{2} \int_0^\pi \frac{\partial}{\partial t} [u_t^2(x,t) + u_x^2(x,t)] dx \\
 &= \int_0^\pi [u_t(x,t) u_{tt}(x,t) + u_x(x,t) u_{xt}(x,t)] dx \stackrel{\text{by (1)}}{=} \int_0^\pi u_t(x,t) [u_{xx}(x,t) - 2u_{xt}(x,t)] dx + \int_0^\pi u_x(x,t) u_{xt}(x,t) dx \\
 &= \int_0^\pi u_t(x,t) u_{xx}(x,t) dx - 2 \int_0^\pi u_t^2(x,t) dx + \left. u_x(x,t) u_t(x,t) \right|_{x=0}^{x=\pi} - \int_0^\pi u_t(x,t) u_{xx}(x,t) dx \\
 &= -2 \int_0^\pi u_t^2(x,t) dx + u_x(\pi,t) u_t(\pi,t) - u_x(0,t) u_t(0,t) \\
 &\leq 0. \quad \boxed{\text{Thus } E \downarrow \text{ on } 0 \leq t < \infty.}
 \end{aligned}$$

$$\begin{aligned}
 u_t(0,t) &= \lim_{k \rightarrow 0} \frac{u(0+t+k) - u(0,t)}{k} \\
 &= 0
 \end{aligned}$$

and similarly for $u_t(\pi,t)$.

(c) There is only one solution to (1)-(2)-(3)-(4). To see this, suppose that $u = u(x,t)$ denotes the solution found above in part (a) and $v = v(x,t)$ is another solution. Then $v(x,t) = u(x,t) - u_v(x,t)$ solves

$$(1') \quad v_{tt} - v_{xx} + 2v_t = 0 \quad \text{in } 0 < x < \pi, 0 < t < \infty,$$

$$(2') \quad v(0,t) = 0 = v(\pi,t) \quad \text{for } t \geq 0,$$

(OVER)

$$(3') \quad v(x,0) = 0 \quad \text{and} \quad (4') \quad v_t(x,0) = 0 \quad \text{for } 0 \leq x \leq \pi.$$

By part (b), the energy function of v ,

$$E(t) = \frac{1}{2} \int_0^\pi [v_t^2(x,t) + v_x^2(x,t)] dx,$$

is nonincreasing on $[0, \infty)$. Thus, for all $t \geq 0$,

$$0 \leq E(t) \leq E(0) = \frac{1}{2} \int_0^\pi [v_t^2(x,0) + v_x^2(x,0)] dx.$$

$$\text{According to } (3'), \quad v_x(x,0) = \lim_{h \rightarrow 0} \frac{v(x+h,0) - v(x,0)}{h} = \lim_{h \rightarrow 0} \frac{0 - 0}{h} = 0$$

for all $0 < x < \pi$, and (4') says $v_t(x,0) = 0$ for $0 \leq x \leq \pi$; thus $E(0) = 0$.

It follows that $E(t) = 0$ for all $t \geq 0$. The vanishing theorem then

gives $v_t(x,t) = 0 = v_x(x,t)$ for all $0 \leq x \leq \pi, 0 \leq t$, from which it

follows that $v(x,t) = \text{constant}$. In light of (2') and (3'), $v(x,t) = 0$

for all $0 \leq x \leq \pi$ and all $t \geq 0$; i.e. $u_i = u$.