

- (25 pts.) (a) State and prove the theorem known as the (weak) maximum-minimum principle for (solutions to) Laplace's equation.  
(b) Give a physical interpretation for the result in part (a).

(a) Theorem: If  $u = u(x, y)$  solves  $u_{xx} + u_{yy} = 0$  in an open bounded set  $D$  in the plane and is continuous on the closure  $\bar{D} = D \cup \partial D$  of  $D$ , then the maximum and minimum values of  $u$  on  $\bar{D}$  are attained on the boundary  $\partial D$  of  $D$ .

Proof: Let  $\epsilon > 0$  and let  $v(x, y) = u(x, y) + \epsilon(x^2 + y^2)$  for  $(x, y)$  in  $\bar{D}$ . We claim that  $v$  attains its maximum on  $\bar{D}$  at a point of  $\partial D$ . For suppose not; then  $v$  attains its maximum on  $\bar{D}$  at (an interior point)  $(x_0, y_0)$  in  $D$ . This implies  $v_x(x_0, y_0) = v_y(x_0, y_0) = 0$ ,  $v_{xx}(x_0, y_0) \leq 0$ , and  $v_{yy}(x_0, y_0) \leq 0$ . But then

$$0 \geq v_{xx}(x_0, y_0) + v_{yy}(x_0, y_0) = u_{xx}(x_0, y_0) + u_{yy}(x_0, y_0) + 4\epsilon = 4\epsilon,$$

which contradicts  $\epsilon > 0$ . This proves the claim.

Note that  $M = \max_{(x, y) \in \bar{D}} (x^2 + y^2)$  is finite since  $D$  is bounded. Thus

$$\begin{aligned} (+) \quad \max_{(x, y) \in \bar{D}} u(x, y) &\leq \max_{(x, y) \in \bar{D}} [u(x, y) + \epsilon(x^2 + y^2)] = \max_{(x, y) \in \bar{D}} v(x, y) = \max_{(x, y) \in \partial D} v(x, y) \\ &= \max_{(x, y) \in \partial D} [u(x, y) + \epsilon(x^2 + y^2)] \leq \epsilon M + \max_{(x, y) \in \partial D} u(x, y). \end{aligned}$$

Because  $\epsilon > 0$  is arbitrary, it follows from (+) that  $\max_{(x, y) \in \bar{D}} u(x, y) \leq \max_{(x, y) \in \partial D} u(x, y)$ .

Since  $\partial D \subseteq \bar{D}$ , the reverse inequality clearly holds, and hence

$$\max_{(x, y) \in \bar{D}} u(x, y) = \max_{(x, y) \in \partial D} u(x, y); \quad \text{i.e. the maximum value of } u \text{ on the closure}$$

of  $D$  is attained on the boundary of  $D$ .

To prove the analogous minimum value result, note that  $-u$  solves the Laplace equation in  $D$  and is continuous on  $\bar{D}$ . Thus, the maximum principle holds for  $-u$ , so

$$\min_{(x, y) \in \bar{D}} u(x, y) = - \max_{(x, y) \in \bar{D}} [-u(x, y)] = - \max_{(x, y) \in \partial D} [-u(x, y)] = \min_{(x, y) \in \partial D} u(x, y).$$

[See back side for (b).]

#1 (b) The steady-state temperature  $u(x,y)$  at position  $(x,y)$  of a laminar region  $\bar{D}$  satisfies the Laplace equation  $u_{xx} + u_{yy} = 0$  in  $D$ . If  $D$  is bounded then the maximum-minimum principle says that the hottest and coldest temperatures in  $\bar{D}$  occur on the boundary of  $D$ .

2. (25 pts.) Find the solution to

$$u_t - u_{xx} = 0 \quad \text{for } -\infty < x < \infty, 0 < t < \infty,$$

which satisfies

$$u(x, 0) = x^2 \quad \text{for } -\infty < x < \infty.$$

You may find the following identities useful:

$$\int_{-\infty}^{\infty} e^{-p^2} dp = \sqrt{\pi}, \quad \int_{-\infty}^{\infty} p e^{-p^2} dp = 0, \quad \int_{-\infty}^{\infty} p^2 e^{-p^2} dp = \frac{\sqrt{\pi}}{2}.$$

By a formula from Section 2.4, a candidate for the solution is

$$(*) \quad u(x, t) = \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4t}} y^2 dy \quad \text{for } -\infty < x < \infty \text{ and } 0 < t < \infty.$$

[Since  $\varphi(y) = y^2$  is not bounded on  $-\infty < y < \infty$ , we are not guaranteed that this formula actually gives the solution to the problem; we will need to check our final form.]

Make the substitution  $p = \frac{y-x}{\sqrt{4t}}$  in the integral of (\*). Then  $dp = \frac{dy}{\sqrt{4t}}$  so

$$\begin{aligned} u(x, t) &= \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} (x + p\sqrt{4t})^2 dp \\ &= \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} (x^2 + 2xp\sqrt{4t} + 4tp^2) dp \\ &= \frac{x^2}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} dp + \frac{2x\sqrt{4t}}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} p dp + \frac{4t}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} p^2 dp \\ &= x^2 \cdot 1 + \frac{2x\sqrt{4t}}{\sqrt{\pi}} \cdot 0 + 4t \cdot \frac{1}{2} \end{aligned}$$

$$\boxed{u(x, t) = x^2 + 2t}$$

Check:  $u_t - u_{xx} = 2 - 2 = 0$

$u(x, 0) = x^2$



3. (25 pts.) Let  $f$  and  $\psi$  be piecewise-continuous, absolutely integrable functions on  $(-\infty, \infty)$ . Use Fourier transform methods to solve

$$u_{xx} + u_{yy} = 0 \quad \text{for } -\infty < x < \infty, 0 < y < \infty,$$

subject to the boundary condition

$$u(x, 0) = f(x) \quad \text{for } -\infty < x < \infty$$

and the decay conditions

$$\lim_{y \rightarrow \infty} u(x, y) = 0 \quad \text{for each } x \in (-\infty, \infty)$$

and, for each  $y > 0$ ,

$$|u(x, y)| \leq |\psi(x)| \quad \text{for all } -\infty < x < \infty.$$

Bonus (10 pts.): Compute an explicit formula for the solution to the above problem if the function  $f$  is given by  $f(x) = 1$  for  $|x| < 1$ , and  $f(x) = 0$  otherwise.

$$\mathcal{F}(u_{xx} + u_{yy})(\xi) = \mathcal{F}(0)(\xi)$$

$$\mathcal{F}(u_{xx})(\xi) + \mathcal{F}(u_{yy})(\xi) = 0$$

$$(i\xi)^2 \mathcal{F}(u)(\xi) + \frac{\partial^2}{\partial y^2} \mathcal{F}(u)(\xi) = 0$$

$$\frac{\partial^2}{\partial y^2} \mathcal{F}(u)(\xi) - \xi^2 \mathcal{F}(u)(\xi) = 0$$

$$\mathcal{F}(u)(\xi) = c_1(\xi) e^{\xi y} + c_2(\xi) e^{-\xi y}$$

The decay conditions imply

$$\begin{aligned} (**) \lim_{y \rightarrow \infty} \mathcal{F}(u(\cdot, y))(\xi) &= \lim_{y \rightarrow \infty} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u(x, y) e^{-i\xi x} dx \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \lim_{y \rightarrow \infty} u(x, y) e^{-i\xi x} dx \\ &= 0 \end{aligned}$$

for all  $-\infty < \xi < \infty$ . On the other hand,

$$(***) \begin{cases} \lim_{y \rightarrow \infty} e^{\xi y} = +\infty & \text{if } \xi > 0, \\ \lim_{y \rightarrow \infty} e^{-\xi y} = +\infty & \text{if } \xi < 0. \end{cases}$$

Therefore ~~(\*)~~ - ~~(\*\*)~~ - ~~(\*\*\*)~~ imply  $c_1(\xi) = 0$  if  $\xi > 0$  and  $c_2(\xi) = 0$  if  $\xi < 0$ . Thus

$$\mathcal{F}(u)(\xi) = \begin{cases} c_2(\xi) e^{-\xi y} & \text{if } \xi > 0 \\ c_1(\xi) e^{\xi y} & \text{if } \xi < 0 \end{cases} = c(\xi) e^{-|\xi| y}$$

Applying the boundary condition yields

$$\mathcal{F}(f)(\xi) = \mathcal{F}(u(\cdot, 0))(\xi) = c(\xi) e^0 = c(\xi).$$

Thus,

$$\begin{aligned} \mathcal{F}(u)(\xi) &= \mathcal{F}(f)(\xi) e^{-|\xi| y} \stackrel{\text{Table, Entry C (with } a=y)}{=} \mathcal{F}(f)(\xi) \mathcal{F}\left(\sqrt{\frac{2}{\pi}} \frac{y}{(\cdot)^2 + y^2}\right)(\xi) \\ &= \frac{1}{\sqrt{2\pi}} \mathcal{F}\left(f * \sqrt{\frac{2}{\pi}} \frac{y}{(\cdot)^2 + y^2}\right)(\xi). \end{aligned}$$

Therefore

$$u(x, y) = \frac{1}{\pi} \left( f * \frac{y}{(\cdot)^2 + y^2} \right)(x)$$

$$u(x, y) = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{y f(s) ds}{(x-s)^2 + y^2}$$

for  $-\infty < x < \infty, 0 < y < \infty$ .

(OVER FOR BONUS)

#3 BONUS: If  $f(x) = \begin{cases} 1 & \text{for } |x| < 1, \\ 0 & \text{otherwise,} \end{cases}$

then the solution to #3 is

$$u(x,y) = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{f(s)y ds}{(x-s)^2 + y^2} = \frac{1}{\pi} \int_{-1}^1 \frac{1 \cdot y ds}{(x-s)^2 + y^2}$$

Let  $v = \frac{s-x}{y}$   
Then  $dv = \frac{ds}{y}$

$$\therefore u(x,y) = \frac{1}{\pi} \int_{\frac{-1-x}{y}}^{\frac{1-x}{y}} \frac{y^2 dv}{y^2 v^2 + y^2}$$

$$= \frac{1}{\pi} \int_{\frac{-1-x}{y}}^{\frac{1-x}{y}} \frac{dv}{v^2 + 1}$$

$$= \frac{1}{\pi} \left[ \text{Arctan} \left( \frac{1-x}{y} \right) - \text{Arctan} \left( \frac{-1-x}{y} \right) \right]$$

$$\boxed{u(x,y) = \frac{1}{\pi} \left[ \text{Arctan} \left( \frac{x+1}{y} \right) - \text{Arctan} \left( \frac{x-1}{y} \right) \right]}$$

for  $-\infty < x < \infty$ ,  $0 < y < \infty$ .

4. (25 pts.) Use separation of variables to find a formal solution to

$$u_t - u_{xx} = 0 \quad \text{for } 0 < x < 1, 0 < t < \infty,$$

subject to

$$u_x(0,t) = 0 = u_x(1,t) - u(1,t) \quad \text{for } 0 \leq t < \infty,$$

and

$$u(x,0) = \phi(x) \quad \text{for } 0 \leq x \leq 1$$

where  $\phi$  is a (given) continuous function on  $[0,1]$ .

Bonus (10 pts.): Is there at most one solution to the above problem? Why or why not?

$$u(x,t) = X(x)T(t) \Rightarrow X(x)T'(t) - X''(x)T(t) = 0 \Rightarrow \frac{X''(x)}{X(x)} = \frac{T'(t)}{T(t)} = \text{const.} = -\lambda.$$

$$X'(0)T(t) = 0 \text{ for } 0 \leq t \Rightarrow X'(0) = 0; X'(1)T(t) - X(1)T(t) = 0 \text{ for } 0 \leq t \Rightarrow X'(1) = X(1).$$

$$\begin{cases} X''(x) + \lambda X(x) = 0, & X'(0) = 0 = X'(1) - X(1) \\ T'(t) + \lambda T(t) = 0 \end{cases} \quad (\text{eigenvalue problem})$$

Case  $\lambda > 0$  (say  $\lambda = \beta^2$  where  $\beta > 0$ ):  $X(x) = c_1 \cos(\beta x) + c_2 \sin(\beta x)$

$$X'(x) = -\beta c_1 \sin(\beta x) + \beta c_2 \cos(\beta x)$$

$$0 = X'(0) = \beta c_2 \Rightarrow c_2 = 0$$

$$0 = X'(1) - X(1) = [-\beta \sin(\beta) - \cos(\beta)] c_1$$

$$\Rightarrow \tan(\beta) = -\frac{1}{\beta} \quad (*)$$

There is an infinite sequence of <sup>positive</sup> eigenvalues  $\lambda_1 = \beta_1^2, \lambda_2 = \beta_2^2, \dots$  where  $\beta_n$  is the  $n^{\text{th}}$  positive solution to (\*). (See figure 1.)

Eigenfunctions:  $X_n(x) = \cos(\beta_n x) \quad (n=1,2,3,\dots)$

$$\therefore T_n'(t) + \lambda_n T_n(t) = 0 \Rightarrow T_n(t) = e^{-\lambda_n t} = e^{-\beta_n^2 t}$$

$$(\beta_1 \approx 2.79838604578, \beta_2 \approx 6.1212504669, \dots)$$

Case  $\lambda = 0$ :  $X(x) = c_1 x + c_2$   
 $X'(x) = c_1$

$$0 = X'(0) = c_1$$

$$0 = X'(1) - X(1) = 0 - c_2$$

No nontrivial solutions in this case.

Case  $\lambda < 0$  (say  $\lambda = -\beta^2$  where  $\beta > 0$ ):  $X(x) = c_1 \cosh(\beta x) + c_2 \sinh(\beta x)$

$$X'(x) = \beta c_1 \sinh(\beta x) + \beta c_2 \cosh(\beta x)$$

$$0 = X'(0) = \beta c_2 \Rightarrow c_2 = 0$$

$$0 = X'(1) - X(1) = [\beta \sinh(\beta) - \cosh(\beta)] c_1$$

$$\Rightarrow \tanh(\beta) = \frac{1}{\beta} \quad (**)$$

There is one <sup>negative</sup> eigenvalue  $\lambda_0 = -\beta_0^2$  where  $\beta_0$  is the unique positive solution to (\*\*). (See figure 2.) ( $\beta_0 \approx 1.19767864026$ )

Eigenfunction:  $X_0(x) = \cosh(\beta_0 x)$

$$T_0'(t) + \lambda_0 T_0(t) = 0 \Rightarrow T_0(t) = e^{-\lambda_0 t} = e^{\beta_0^2 t}$$

(cont.)

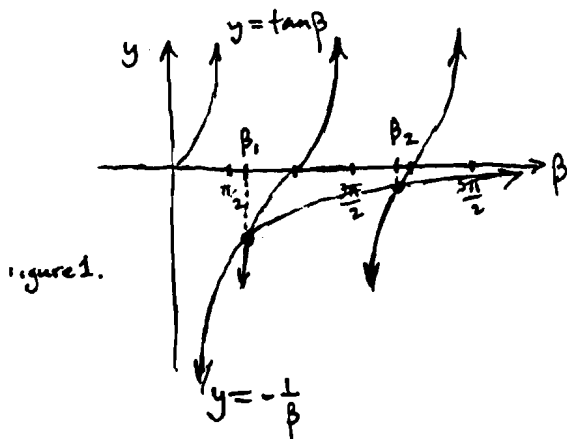


Figure 1.

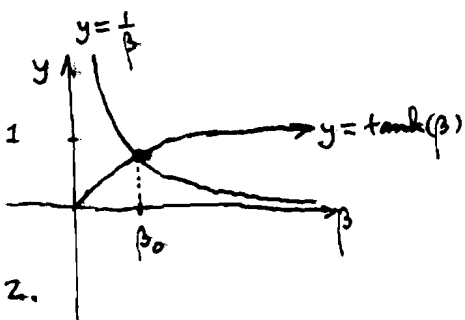


Figure 2.

#4 (cont.)

$$\therefore u(x,t) = c_0 \cosh(\beta_0 x) e^{\beta_0^2 t} + \sum_{n=1}^{\infty} c_n \cos(\beta_n x) e^{-\beta_n^2 t}$$

( $c_0, c_1, c_2, \dots$  "arbitrary" constants)

is a formal solution to the homogeneous part of #4. In order to satisfy the initial condition the constants must be chosen (if possible) to satisfy

$$\varphi(x) = u(x,0) = c_0 \cosh(\beta_0 x) + \sum_{n=1}^{\infty} c_n \cos(\beta_n x) \quad \text{for all } 0 \leq x \leq 1.$$