

## Chapter 2: Waves and Diffusions

### Sec 2.1 The Wave Equation

Ex 1 (similar to #10, p.37) Solve  $u_{xx} + u_{xt} - 20u_{tt} = 0$  in the  $xt$ -plane subject to the initial conditions

$$u(x, 0) = 25x^2 + 64x^3 \text{ and } u_t(x, 0) = -10x + 48x^2 \text{ for all } -\infty < x < \infty.$$

Soln: From previous work, <sup>(a) in Sec. 1.6</sup> we know that the general solution to the PDE in the  $xt$ -plane is  $u(x, t) = f(5x-t) + g(4x+t)$  where  $f$  and  $g$  are  $C^2$ -functions of a single real variable. We must determine  $f$  and  $g$  so the I.C.'s are satisfied.

$$(1) \quad 25x^2 + 64x^3 = u(x, 0) = f(5x) + g(4x) \text{ for all } -\infty < x < \infty.$$

$$u_t(x, 0) = -f'(5x) + g'(4x)$$

$$(2) \quad \begin{cases} -10x + 48x^2 = u_t(x, 0) = -f'(5x) + g'(4x) \text{ for all } -\infty < x < \infty. \end{cases}$$

$$(1') \quad \begin{cases} 50x + 192x^2 = 5f'(5x) + 4g'(4x) \text{ " " " " } \end{cases}$$

Multiplying (2) by 5 and adding to (1') gives

$$432x^2 = 9g'(4x) \Rightarrow 48x^2 = g'(4x)$$

$$\text{Letting } z = 4x \text{ gives } g'(z) = 3z^2$$

$$\Rightarrow g(z) = z^3 + c_1 \text{ for all } -\infty < z < \infty$$

Multiplying (2) by -4 and adding to (1') gives

$$90x = 9f'(5x) \Rightarrow 10x = f'(5x)$$

$$\text{(Let } z = 5x) \Rightarrow f'(z) = 2z$$

$$\Rightarrow f(z) = z^2 + c_2 \text{ for all } -\infty < z < \infty$$

$$\therefore u(x,t) = f(5x-t) + g(4x+t)$$

$$u(x,t) = (5x-t)^2 + c_2 + (4x+t)^3 + c_1$$

However, the constants  $c_1$  and  $c_2$  are not independent. For example,

$$\begin{aligned} 25x^2 + 64x^3 = u(x,0) &= (5x)^2 + c_2 + (4x)^3 + c_1 \\ &= 25x^2 + c_2 + 64x^3 + c_1 \end{aligned}$$

so we need  $c_1 + c_2 = 0$ . Therefore the particular solution to initial value our problem is

$$u(x,t) = (5x-t)^2 + (4x+t)^3$$

Similar techniques give the following result:

The solution to  $u_{xt} - c^2 u_{xx} = 0$  in the  $xt$ -plane satisfying the initial conditions

$$u(x,0) = \varphi(x) \quad \text{and} \quad u_t(x,0) = \psi(x) \quad \text{for all } -\infty < x < \infty$$

is given by

d'Alembert's formula  $\longrightarrow u(x,t) = \frac{1}{2} [\varphi(x+ct) + \varphi(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(\xi) d\xi.$

(Here we assume that  $\varphi$  is a  $C^2$ -function and  $\psi$  is a  $C^1$ -function of a single real variable.)

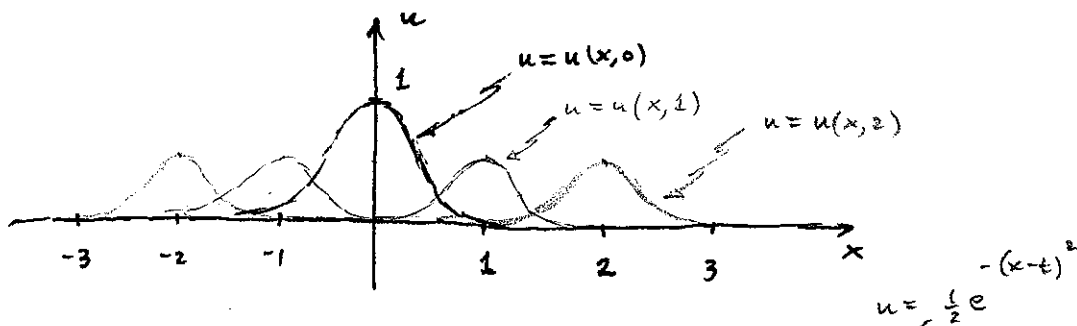
Note: You are responsible for ~~the derivation of~~ <sup>the derivation of</sup> this result on Exam I.

Ex 2 Solve  $u_{tt} - u_{xx} = 0$  in the  $xt$ -plane subject to

$$u(x,0) = e^{-x^2} \text{ and } u_t(x,0) = 0 \text{ for } -\infty < x < \infty.$$

Solu: 
$$u(x,t) = \frac{1}{2} \left[ e^{-\frac{(x-t)^2}{2}} + e^{-\frac{(x+t)^2}{2}} \right] + \frac{1}{2} \int_{x-t}^{x+t} 0 d\xi$$

$$= \frac{1}{2} e^{-\frac{(x-t)^2}{2}} + \frac{1}{2} e^{-\frac{(x+t)^2}{2}}$$



Note: The solution consists of two "traveling" waves, one that moves to the right along the  $x$ -axis with speed 1 and one that moves to the left with speed 1.

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Ex 3 (#5, p.37 The Hammer Blow) Find the solution of  $u_{tt} - c^2 u_{xx} = 0$  satisfying  $u(x,0) = \varphi(x) = 0$  and  $u_t(x,0) = \psi(x) = \begin{cases} 1 & \text{if } -a < x < a, \\ 0 & \text{o.w.} \end{cases}$

Sketch the profile of the string at each of the successive instants

$$t = \frac{a}{2c}, \frac{a}{c}, \frac{3a}{2c}, \frac{2a}{c}, \text{ and } \frac{5a}{c}.$$

Solu: 
$$u(x,t) = \frac{1}{2} \left[ \varphi(x-ct) + \varphi(x+ct) \right] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(\xi) d\xi$$

$$= \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(\xi) d\xi$$

We will make use of the following facts to evaluate the integral.

Characteristic (or indicator) function of A

$$0. \text{ If } A \subseteq \mathbb{R} \text{ then } \chi_A(\xi) = \begin{cases} 1 & \text{if } \xi \in A, \\ 0 & \text{o.w.} \end{cases}$$

$$1. \int_c^d f(\xi) d\xi = \int_{-\infty}^{\infty} \chi_{(c,d)}(\xi) f(\xi) d\xi$$

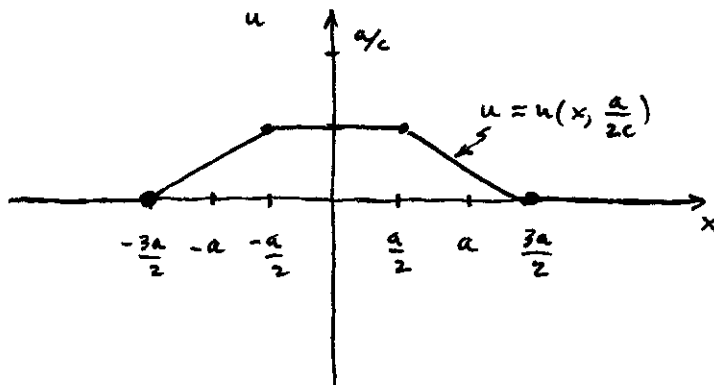
$$2. \chi_A(\xi) \chi_B(\xi) = \chi_{A \cap B}(\xi) \text{ for all } A, B \subseteq \mathbb{R} \text{ and all } \xi \in \mathbb{R}.$$

$$3. \int_{-\infty}^{\infty} \chi_{(c,d)}(\xi) d\xi = \text{length of } (c,d) (= d-c).$$

$$\begin{aligned} \therefore u(x,t) &= \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(\xi) d\xi \stackrel{0.}{=} \frac{1}{2c} \int_{x-ct}^{x+ct} \chi_{(-a,a)}(\xi) d\xi \\ &\stackrel{1.}{=} \frac{1}{2c} \int_{-\infty}^{\infty} \chi_{(x-ct, x+ct)}(\xi) \chi_{(-a,a)}(\xi) d\xi \\ &\stackrel{2.}{=} \frac{1}{2c} \int_{-\infty}^{\infty} \chi_{(x-ct, x+ct) \cap (-a,a)}(\xi) d\xi \\ &\stackrel{3.}{=} \frac{1}{2c} \text{length of } \{(x-ct, x+ct) \cap (-a,a)\}. \end{aligned}$$

$$\boxed{t = \frac{a}{2c}}$$

$$u(x, \frac{a}{2c}) = \frac{1}{2c} \text{length of } \{(x - \frac{a}{2}, x + \frac{a}{2}) \cap (-a, a)\}$$



$$t = \frac{a}{c}$$

$$u(x, \frac{a}{c}) = \frac{1}{2c} \text{ length of } \left\{ (x-a, x+a) \cap (-a, a) \right\}$$

$$t = \frac{3a}{2c}$$

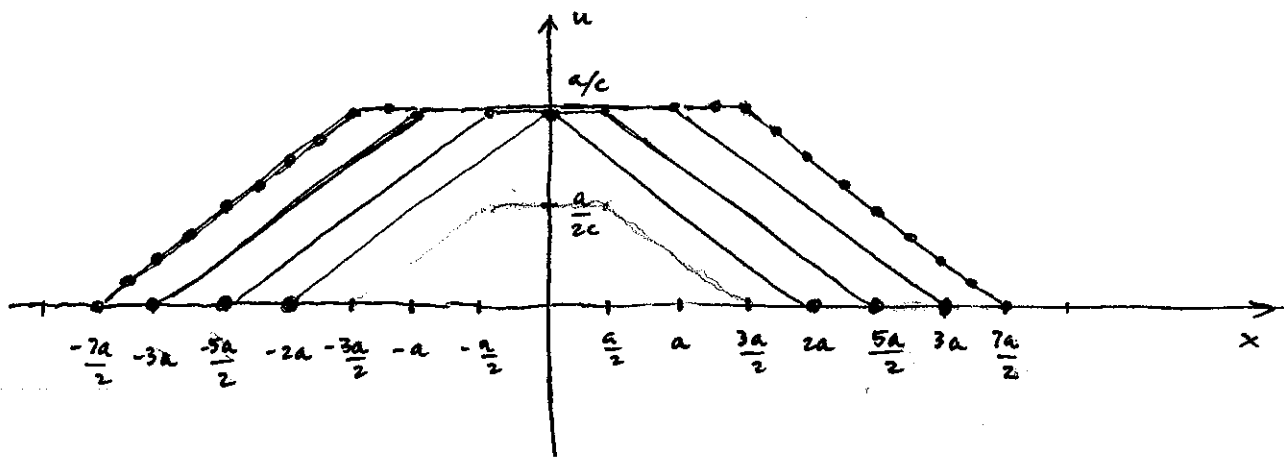
$$u(x, \frac{3a}{2c}) = \frac{1}{2c} \text{ length of } \left\{ (x - \frac{3a}{2}, x + \frac{3a}{2}) \cap (-a, a) \right\}$$

$$t = \frac{2a}{c}$$

$$u(x, \frac{2a}{c}) = \frac{1}{2c} \text{ length of } \left\{ (x-2a, x+2a) \cap (-a, a) \right\}$$

$$t = \frac{5a}{2c}$$

$$u(x, \frac{5a}{2c}) = \frac{1}{2c} \text{ length of } \left\{ (x - \frac{5a}{2}, x + \frac{5a}{2}) \cap (-a, a) \right\}$$



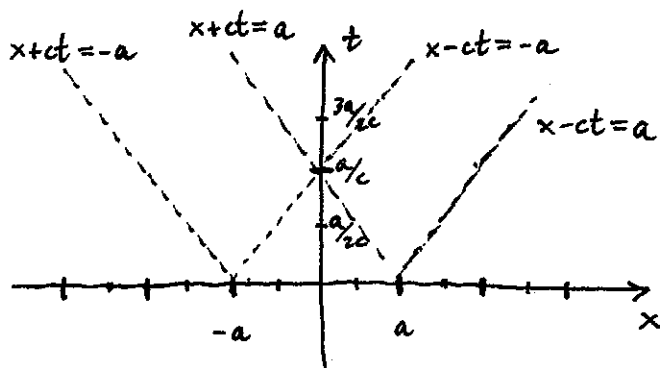
Notes: ① The function  $u(x,t) = \frac{1}{2c} \text{ length of } \left\{ (x-ct, x+ct) \cap (-a, a) \right\}$  is

called a "generalized solution" of the problem:  $u_{tt} - c^2 u_{xx} = 0$  in  $-\infty < x < \infty, 0 < t < \infty$ ,

and  $u(x,0) = 0$  and  $u_t(x,0) = \chi_{(-a,a)}(x)$  for  $-\infty < x < \infty$ . Note that

$u = u(x,t)$  fails to be differentiable at all points on the lines

$x-ct = \pm a$  and  $x+ct = \pm a$ , hence it is not a "classical solution"



of the problem (i.e. a function of  $x$  and  $t$  that:

(a) is  $C^2$  and satisfies the PDE at all points of the region  $-\infty < x < \infty, 0 < t < \infty$ ;

(b) is  $C^1$  on the closure of the region  $-\infty < x < \infty, 0 \leq t < \infty$ ;

and (c) satisfies the initial conditions on the boundary  $t=0$  of the region.)

② The singularities of the generalized solution (i.e. points where the function fails to be differentiable) propagate along those characteristics  $|x \pm ct| = a$  of the PDE  $u_{tt} - c^2 u_{xx} = 0$  which pass through the points  $(\pm a, 0)$  where the initial data is singular.