

1.(25 pts.) Find all second-degree polynomial functions of two real variables,

$$u(x,t) = ax^2 + bxt + ct^2 + dx + et + f$$

where $a, b, c, d, e,$ and f are real constants, which are solutions in the xt -plane of the one-dimensional diffusion equation

$$u_t - ku_{xx} = 0.$$

$$u_t = bx + 2ct + e$$

$$u_x = 2ax + bt + d$$

$$u_{xx} = 2a$$

We want $u_t = ku_{xx}$ so substituting gives

$$bx + 2ct + e = k2a. \quad 12 \text{ pts. to here.}$$

This is to hold for all $-\infty < x < \infty$ and $-\infty < t < \infty$.

Therefore we must have "like" coefficients equal on the left and right sides of this identity. That is, $b=0$, $2c=0$, and $e=2ka$. Thus

$$u(x,t) = ax^2 + dx + 2kat + f$$

$$u(x,t) = a(x^2 + 2kt) + dx + f$$

25 pts. to here.

where $a, d,$ and f are arbitrary real constants.

6 pts.
to here.

18 pts. to here.

2.(25 pts.) Find the general solution of

$$yu_x - xu_y = 0$$

in the xy -plane. Sketch several characteristic curves of this partial differential equation.

The solution $u = u(x, y)$ is constant along all curves in the xy -plane which satisfy the characteristic equation of this pde: $\frac{dy}{dx} = \frac{b(x, y)}{a(x, y)} = \frac{-x}{y}$. 6 pts. to here.
Separating variables and integrating yields

$$\frac{y^2}{2} = \int y dy = \int -x dx = -\frac{x^2}{2} + c_1$$

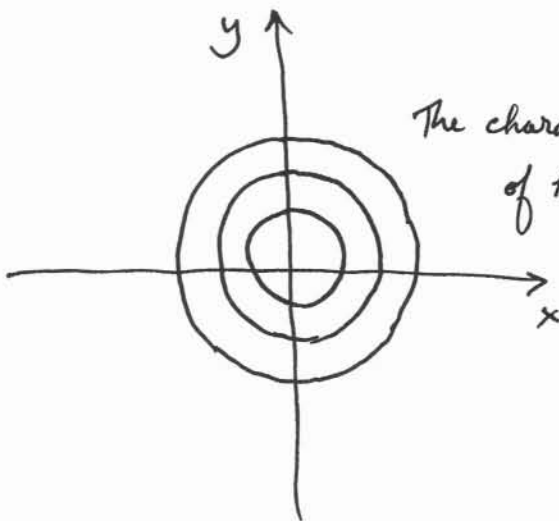
$$\Rightarrow x^2 + y^2 = c. \quad (\text{where } c = 2c_1)$$

12 pts. to here.

Along each such curve we have

$$u(x, y) = u(x, \pm\sqrt{c-x^2}) = u(0, \pm\sqrt{c}) = f(c).$$

Therefore $u(x, y) = f(x^2 + y^2)$ where f is an ^{arbitrary} differentiable function of a single real variable. 20 pts. to here.



The characteristic curves $x^2 + y^2 = c$ of the pde are circles centered at the origin. 25 pts. to here.

3.(25 pts.) Consider the linearized gas dynamics equations

$$\frac{\partial \mathbf{v}}{\partial t} + \frac{c_0^2}{\rho_0} \text{grad}(\rho) = \mathbf{0}$$

$$\frac{\partial \rho}{\partial t} + \rho_0 \text{div}(\mathbf{v}) = 0$$

where ρ_0 is the density and c_0 is the speed of sound in still air. Verify that if $\text{curl}(\mathbf{v}) = \mathbf{0}$ when $t = 0$, then $\text{curl}(\mathbf{v}) = \mathbf{0}$ at all later times.

Suppose $\vec{v} = \vec{v}(x, y, z, t)$ and $\rho = \rho(t)$ are solutions to the gas dynamics equation. Consider the function

$$\vec{w} = \text{curl}(\vec{v}(x, y, z, t)).$$

Then $\frac{\partial \vec{w}}{\partial t} = \frac{\partial}{\partial t} \text{curl}(\vec{v}(x, y, z, t)) = \text{curl}\left(\frac{\partial \vec{v}}{\partial t}\right)$ because $\vec{v} = \vec{v}(x, y, z, t)$

is a C^2 -function so $\frac{\partial}{\partial t} \left(\frac{\partial v_3}{\partial y} - \frac{\partial v_2}{\partial z} \right) = \frac{\partial}{\partial y} \left(\frac{\partial v_3}{\partial t} \right) - \frac{\partial}{\partial z} \left(\frac{\partial v_2}{\partial t} \right)$, etc. But then

$$\frac{\partial \vec{w}}{\partial t} = \text{curl}\left(-\frac{c_0^2}{\rho_0} \text{grad}(\rho)\right) = -\frac{c_0^2}{\rho_0} \text{curl}(\text{grad}(\rho)) = \vec{0}$$

from the first gas dynamic equation and the fact that the curl of a gradient of a C^2 -function is zero. It follows that the function \vec{w} is independent of t ; i.e. for all $t \geq 0$ and all $(x, y, z) \in \mathbb{R}^3$,

$$\text{curl}(\vec{v}(x, y, z, t)) = \vec{w}(x, y, z, t) = \vec{w}(x, y, z, 0) = \text{curl}(\vec{v}(x, y, z, 0)) = \vec{0}.$$

4. (25 pts.) Consider the partial differential equation

(*) $u_{xx} - 3u_{xt} - 4u_{tt} = 0.$

- (a) Classify (*) as elliptic, hyperbolic, or parabolic.
- (b) Find the general solution of (*) in the xt -plane.
- (c) If ϕ and ψ are C^2 and C^1 functions of a single real variable, respectively, find the solution of (*) which satisfies the initial conditions

$$u(x,0) = \phi(x) \text{ and } u_t(x,0) = \psi(x) \text{ for all } -\infty < x < \infty.$$

5 pts.

(a) $B^2 - 4AC = (-3)^2 - 4(1)(-4) = 9 + 16 = 25 > 0.$ (*) is hyperbolic.

(b) (*) is equivalent to $(\frac{\partial^2}{\partial x^2} - 3\frac{\partial^2}{\partial x \partial t} - 4\frac{\partial^2}{\partial t^2})u = 0$ and thus to

$(\frac{\partial}{\partial x} - 4\frac{\partial}{\partial t})(\frac{\partial}{\partial x} + \frac{\partial}{\partial t})u = 0.$ The change-of-coordinates $\begin{cases} \xi = -(3x - 4t) = -(-4x - t) \\ \eta = 5x - 4t = x - t \end{cases}$

transforms (*) into $\frac{\partial^2 u}{\partial \xi \partial \eta} = 0$ so the general solution of (*) is 10 pts. to here.

$u = f(\xi) + g(\eta) = f(4x + t) + g(x - t)$ where f and g are arbitrary twice-differentiable functions of a single real variable.

(c) Using (b) and the initial conditions in (c), we have

(1) $\phi(x) = u(x,0) = f(4x) + g(x) \quad (-\infty < x < \infty)$

(2) $\psi(x) = u_t(x,0) = f'(4x) - g'(x) \quad (-\infty < x < \infty)$

Differentiating (1) yields

(1') $\phi'(x) = 4f'(4x) + g'(x) \quad (-\infty < x < \infty).$

Adding (2) and (1') produces

(3) $\phi'(x) + \psi(x) = 5f'(4x) \quad (-\infty < x < \infty)$

or equivalently

(4) $f'(z) = \frac{1}{5}\phi'(\frac{z}{4}) + \frac{1}{5}\psi(\frac{z}{4}) \quad (-\infty < z < \infty).$

Integrating yields

(OVER)

15 pts. to here.

$$(5) \quad f(z) = \frac{1}{5} \varphi\left(\frac{z}{4}\right) + \frac{1}{5} \int_0^{z/4} \psi(\xi) d\xi + c, \quad (-\infty < z < \infty).$$

Therefore, using (5) and (1) we have

$$(b) \quad g(x) = \varphi(x) - f(4x) = \varphi(x) - \left(\frac{1}{5} \varphi(x) + \frac{1}{5} \int_0^x \psi(\xi) d\xi + c, \right) \\ = \frac{1}{5} \varphi(x) - \frac{1}{5} \int_0^x \psi(\xi) d\xi - c, \quad (-\infty < x < \infty)$$

It follows from (b), (5), and (6) that

$$u(x,t) = f(4x+t) + g(x-t) \\ = \frac{1}{5} \varphi\left(\frac{4x+t}{4}\right) + \frac{1}{5} \int_0^{(4x+t)/4} \psi(\xi) d\xi + \cancel{c} + \frac{1}{5} \varphi(x-t) - \frac{1}{5} \int_0^{x-t} \psi(\xi) d\xi - \cancel{c}$$

$$u(x,t) = \frac{1}{5} \left[\varphi\left(x + \frac{1}{4}t\right) + \varphi(x-t) \right] + \frac{1}{5} \int_{x-t}^{x + \frac{1}{4}t} \psi(\xi) d\xi$$

25 pts.
to here.

Bonus.(25 pts.) A homogeneous solid material occupying $D = \{(x, y, z) \in \mathbb{R}^3 : 4 \leq x^2 + y^2 + z^2 \leq 100\}$ is completely insulated and its initial temperature at position (x, y, z) in D is $200/\sqrt{x^2 + y^2 + z^2}$.

(a) Write (without proof or derivation) the partial differential equation and initial/boundary conditions that completely govern the temperature $u(x, y, z, t)$ at position (x, y, z) in D and time $t \geq 0$.

(b) Use Gauss' divergence theorem to help show that the heat energy $H(t) = \iiint_D c\rho u(x, y, z, t) dV$ of the material in D at time t is a constant function of time. Here c and ρ denote the (constant) specific heat and mass density, respectively, of the material in D .

(c) Compute the (constant) steady-state temperature that the material in D reaches after a long time.

9 pts.

$$(a) \begin{cases} u_t - k(\overbrace{u_{xx} + u_{yy} + u_{zz}}^{\nabla^2 u}) \stackrel{\textcircled{1}}{=} 0 & \text{if } 4 < x^2 + y^2 + z^2 < 100 \text{ and } t > 0 \\ \nabla u \cdot \vec{n} \stackrel{\textcircled{2}}{=} 0 & \text{if } x^2 + y^2 + z^2 = 4 \text{ or } x^2 + y^2 + z^2 = 100 \text{ and } t \geq 0 \\ u(x, y, z, 0) \stackrel{\textcircled{3}}{=} \frac{200}{\sqrt{x^2 + y^2 + z^2}} & \text{if } 4 \leq x^2 + y^2 + z^2 \leq 100 \end{cases}$$

8 pts.

$$(b) H'(t) = \frac{d}{dt} \iiint_D c\rho u(x, y, z, t) dV = \iiint_D \frac{\partial}{\partial t} (c\rho u) dV = c\rho \iiint_D u_t dV$$

$$\stackrel{\textcircled{1}}{=} k c \rho \iiint_D \nabla^2 u dV = k c \rho \iiint_D \nabla \cdot (\nabla u) dV \stackrel{\text{Gauss' Divergence Thm.}}{=} k c \rho \iint_{\partial D} \nabla u \cdot \vec{n} dV \stackrel{\textcircled{2}}{=} 0.$$

Therefore the heat energy $H = H(t)$ is a constant function of t .

8 pts.

(c) Let $U = \lim_{t \rightarrow \infty} u(x, y, z, t)$ be the constant steady-state temperature

that the material in D reaches after a long time. From part (b), $H(0) = H(t)$

for all $t \geq 0$. Therefore

$$H(0) = \lim_{t \rightarrow \infty} H(t) = \lim_{t \rightarrow \infty} \iiint_D c\rho u(x, y, z, t) dV = c\rho \iiint_D \lim_{t \rightarrow \infty} u(x, y, z, t) dV$$

$$= c\rho \iiint_D U dV = c\rho U \text{ vol}(D). \quad (\text{OVER})$$

Solving for U we find

$$U = \frac{H(0)}{c \text{vol}(\mathcal{D})}$$

But computations yield

$$\text{vol}(\mathcal{D}) = \frac{4}{3}\pi r_1^3 - \frac{4}{3}\pi r_2^3 = \frac{4}{3}\pi(10^3 - 2^3) = \frac{3968}{3}\pi$$

and

$$\begin{aligned} H(0) &= \iiint_{\mathcal{D}} c\rho u(x,y,z,0) dV = c\rho \iiint_{\mathcal{D}} \frac{200}{\sqrt{x^2+y^2+z^2}} dV \\ &= c\rho \int_0^{2\pi} \int_0^{\pi} \int_2^{10} \frac{200}{r} r^2 \sin\phi dr d\phi d\theta \\ &= c\rho (2\pi) (-\cos\phi) \Big|_0^{\pi} (100r^2) \Big|_2^{10} \\ &= 38,400\pi c\rho. \end{aligned}$$

Therefore

$$U = \frac{38,400\pi c\rho}{c\rho\left(\frac{3968}{3}\pi\right)} = \frac{900}{31} \doteq 29.0$$

Math 325

Exam I

Summer 2009

n : 15

standard deviation : 20.7

mean : 72.6

Distribution of Scores :

frequency

87 - 100

A

3

73 - 86

B

8

60 - 72

C (graduate), B (undergraduate)

1

50 - 59

C

1

0 - 49

F (graduate), D (undergraduate)

2