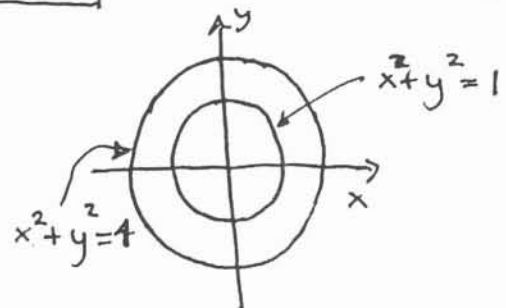


- 1.(25 pts.) (a) Find the characteristic curves of $yu_x - xu_y = 0$ in the xy -plane.
 (b) Sketch and identify two characteristic curves of this partial differential equation.
 (c) Write the general solution of this partial differential equation in the xy -plane.
 (d) Find the solution of this partial differential equation which satisfies $u(x,0) = x^6$ for all real x .
 (e) In what region of the xy -plane is the solution in part (d) uniquely determined?

6 pts. (a) The characteristic curves of the p.d.e. satisfy $\frac{dy}{dx} = \frac{b(x,y)}{a(x,y)} = \frac{-x}{y}$.
 Separating variables gives $ydy = -x dx$ and integrating both sides yields
 $\frac{1}{2}y^2 = -\frac{1}{2}x^2 + c_1$, or equivalently $\boxed{x^2 + y^2 = c}$. ($c = 2c_1$ is an arbitrary constant)

3 pts. (b) The characteristic curves are circles.



8 pts. (c) The solutions to the p.d.e. are constant along characteristic curves. Along such a curve $x^2 + y^2 = c$ we have

$$u(x,y) = u(x, \pm\sqrt{c-x^2}) = u(0, \pm\sqrt{c}) = f(c).$$

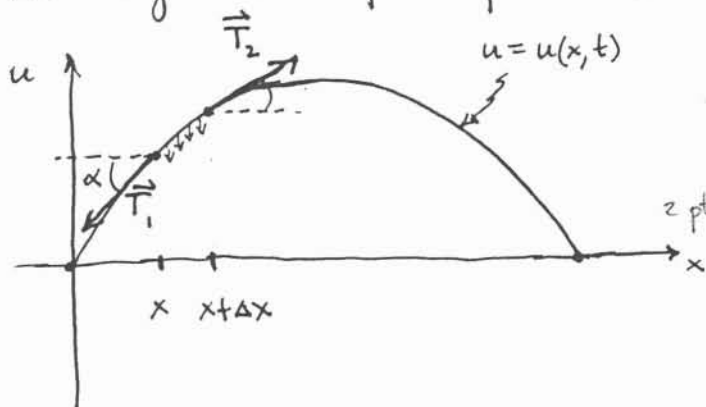
Therefore the general solution to the p.d.e. in the xy -plane is $\boxed{u(x,y) = f(x^2 + y^2)}$
 where f is any C^1 -function of a single real variable.

6 pts. (d) $x^6 = u(x,0) = f(x^2 + 0^2) = f(x^2)$ for all real x . Thus $f(t) = t^3$
 for every $t \geq 0$, so $\boxed{u(x,y) = f(x^2 + y^2) = (x^2 + y^2)^3}$ for all (x,y) in the plane.

2 pts. (e) The solution in part (d) is uniquely determined in the entire xy -plane.

2.(25 pts.) An elastic string is held fixed at the endpoints and plucked. Assuming that air exerts a resistance at each point of the string which is proportional to the string's velocity at that point, give a careful derivation of the partial differential equation that governs the small vibrations of the string.

Fix $t > 0$ and suppose the profile of the string at this instant is as indicated in the diagram. Consider the segment of string between x and $x+\Delta x$. The contact forces \vec{T}_1 and \vec{T}_2 that act at the endpoints of the segment are tangential to the string at the respective points. If α and β are the angles these contact forces make with respect to the horizontal, then



forces make with respect to the horizontal, then

2 pts. (1) $\tan(\alpha) = \text{slope of the tangent line to the profile at } x = u_x(x, t)$

so

4 pts. to here. $\sin(\alpha) \stackrel{(2)}{=} \frac{u(x, t)}{\sqrt{1 + u_x^2(x, t)}}$

and

$\cos(\alpha) \stackrel{(3)}{=} \frac{1}{\sqrt{1 + u_x^2(x, t)}}$

Similarly,

6 pts. to here. $\sin(\beta) \stackrel{(4)}{=} \frac{u_x(x + \Delta x, t)}{\sqrt{1 + u_x^2(x + \Delta x, t)}}$

and

$\cos(\beta) \stackrel{(5)}{=} \frac{1}{\sqrt{1 + u_x^2(x + \Delta x, t)}}$

Applying Newton's second law, $m\vec{a} = \vec{F}_{\text{net}}$, to the system of particles comprising the segment of string from x to $x + \Delta x$, we have upon resolving into components:

8 pts. to here.

(Horizontal) $|\vec{T}_2| \cos(\beta) - |\vec{T}_1| \cos(\alpha) \stackrel{(6)}{=} 0$

15 pts. to here
(2 + 3 + 2)

(Vertical) $|\vec{T}_2| \sin(\beta) - |\vec{T}_1| \sin(\alpha) - \int_x^{x+\Delta x} r u_t(\xi, t) d\xi \stackrel{(7)}{=} \int_x^{x+\Delta x} \rho(\xi) u_{tt}(\xi, t) d\xi$

(OVER)

where r is a positive constant and $\rho(\xi)$ is the linear density of the string at ξ in the interval $[x, x+\Delta x]$. (Here ^{we} have made use of the assumption that the air exerts a resistive force at each point that is proportional to the string's velocity u_t at that point.) Equations (3), (5), and (6) imply

$$(8) \quad |\vec{T}_2| = |\vec{T}_1| \frac{\sqrt{1 + u_x^2(x+\Delta x, t)}}{\sqrt{1 + u_x^2(x, t)}}$$

so substituting from (8) into (7) and using (2) and (4) yields

$$(9) \quad \frac{|\vec{T}_1|}{\sqrt{1 + u_x^2(x, t)}} \left(u_x(x+\Delta x, t) - u_x(x, t) \right) - \int_x^{x+\Delta x} r u_t(\xi, t) d\xi = \int_x^{x+\Delta x} \rho(\xi) u_{tt}(\xi, t) d\xi.$$

Dividing (9) by Δx and letting $\Delta x \rightarrow 0$ produces

$$(10) \quad \frac{|\vec{T}(x, u(x, t), u_x(x, t))|}{\sqrt{1 + u_x^2(x, t)}} u_{xx}(x, t) - r u_t(x, t) = \rho(x) u_{tt}(x, t).$$

For small vibrations, $u(x, t)$ and $u_x(x, t)$ are quite small so $\sqrt{1 + u_x^2(x, t)} \approx 1$ for all $0 \leq x \leq L$. Thus (8) implies the magnitude of the contact force is approximately constant, say $|\vec{T}(x, u(x, t), u_x(x, t))| = \text{constant} = T_0$. Thus (10) becomes

$$(11) \quad \frac{T_0}{\sqrt{1 + u_x^2(x, t)}} u_{xx}(x, t) - r u_t(x, t) = \rho(x) u_{tt}(x, t).$$

If the string is homogeneous, say $\rho(x) = \text{constant} = \rho_0$, and we use the small vibrations approximation, $\sqrt{1 + u_x^2(x, t)} \approx 1$ for all $0 \leq x \leq L$, then

(11) becomes

$$(12) \quad \boxed{T_0 u_{xx}(x, t) - r u_t(x, t) = \rho_0 u_{tt}(x, t)}$$

where T_0 , r , and ρ_0 are positive constants.

3. (25 pts.) Classify the following second order partial differential equations as hyperbolic, parabolic, elliptic, or nonlinear. Find the general C^2 -solution in the plane whenever possible.

(a) $9u_{xx} + 6u_{xy} + u_{yy} + 100u = 0$

(b) $u_{xx} + u_{yy} + uu_x = 0$

(c) $u_{xx} - u_{xy} + 3u_{yy} - 3u_{yx} + 2u_x - 2u_y = 0 \iff u_{xx} - 4u_{xy} + 3u_{yy} + 2u_x - 2u_y = 0$.

2 pts. (a) $B^2 - 4AC = 6^2 - 4(9)(1) = 0$ so the p.d.e. is **parabolic**.

2 pts. (b) The p.d.e. is **nonlinear** because of the term uu_x .

2 pts. (c) $B^2 - 4AC = (-4)^2 - 4(1)(3) = 4 > 0$ so the p.d.e. is **hyperbolic**.

8 pts. to here. To solve the p.d.e. in (a) note that it is equivalent to $(3\frac{\partial}{\partial x} + \frac{\partial}{\partial y})(3\frac{\partial}{\partial x} + \frac{\partial}{\partial y})u + 100u = 0$.

10 pts. to here. Let $\begin{cases} \xi = 3x+y \\ \eta = x-3y \end{cases}$. Then by the chain rule, as operators we have

$$\frac{\partial}{\partial x} = \frac{\partial \xi}{\partial x} \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial x} \frac{\partial}{\partial \eta} = 3\frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta} \quad \text{and} \quad \frac{\partial}{\partial y} = \frac{\partial \xi}{\partial y} \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial y} \frac{\partial}{\partial \eta} = \frac{\partial}{\partial \xi} - 3\frac{\partial}{\partial \eta}$$

Therefore $3\frac{\partial}{\partial x} + \frac{\partial}{\partial y} = 3(3\frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta}) + (\frac{\partial}{\partial \xi} - 3\frac{\partial}{\partial \eta}) = 10\frac{\partial}{\partial \xi}$. Thus the p.d.e.

12 pts. to here. in (a) is equivalent to $(10\frac{\partial}{\partial \xi})(10\frac{\partial}{\partial \xi})u + 100u = 0 \Rightarrow \frac{\partial^2 u}{\partial \xi^2} + u = 0$.

Consequently, $u = c_1(\eta)\cos(\xi) + c_2(\eta)\sin(\xi)$, or equivalently,

14 pts. to here.
$$u(x,y) = f(x-3y)\cos(3x+y) + g(x-3y)\sin(3x+y)$$

is the general solution of (a) with f and g arbitrary C^2 -functions of a single real variable.

16 pts. to here. To solve the p.d.e. in (c) note that it is equivalent to $0 = (\frac{\partial}{\partial x} - 3\frac{\partial}{\partial y})(\frac{\partial}{\partial x} - \frac{\partial}{\partial y})u + 2(\frac{\partial}{\partial x} - \frac{\partial}{\partial y})u$.

18 pts. to here. Let $\begin{cases} \xi = 3x+y \\ \eta = x+y \end{cases}$. Then the chain rule implies, as operators we have

(OVER)

$$\frac{\partial}{\partial x} = \frac{\partial \xi}{\partial x} \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial x} \frac{\partial}{\partial \eta} = 3 \frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta} \quad \text{and} \quad \frac{\partial}{\partial y} = \frac{\partial \xi}{\partial y} \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial y} \frac{\partial}{\partial \eta} = \frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta}.$$

Thus
$$\frac{\partial}{\partial x} - 3 \frac{\partial}{\partial y} = 3 \frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta} - 3 \left(\frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta} \right) = -2 \frac{\partial}{\partial \eta}$$

and
$$\frac{\partial}{\partial x} - \frac{\partial}{\partial y} = 3 \frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta} - \left(\frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta} \right) = 2 \frac{\partial}{\partial \xi}.$$

Consequently, the p.d.e. in (c) is equivalent to

$$\left(-2 \frac{\partial}{\partial \eta} \right) \left(2 \frac{\partial}{\partial \xi} \right) u + 2 \left(2 \frac{\partial}{\partial \xi} \right) u = 0$$

20 pts. to here.

$$- \frac{\partial}{\partial \eta} \left(\frac{\partial u}{\partial \xi} \right) + \frac{\partial u}{\partial \xi} = 0.$$

22 pts. to here. Letting $v = \frac{\partial u}{\partial \xi}$, this becomes $\frac{\partial v}{\partial \eta} - v = 0$. Therefore $v = c(\xi) e^{\eta}$.

That is, $\frac{\partial u}{\partial \xi} = c(\xi) e^{\eta}$ so $u = \int c(\xi) e^{\eta} d\xi = e^{\eta} \int c(\xi) d\xi = e^{\eta} (f(\xi) + c_2(\eta))$

24 pts. to here. Consequently, $u = f(\xi) e^{\eta} + g(\eta) \quad (g(\eta) = e^{\eta} c_2(\eta)),$

or equivalently

$$u(x, y) = f(3x+y) e^{x+y} + g(x+y)$$

25 pts. to here.

is the general solution of (c) where f and g are arbitrary C^2 -functions of a single real variable.

4.(25 pts.) Let $u = u(x, t)$ be a nonconstant C^2 -solution to

$$u_{tt} - u_{xx} + u_t = 0$$

in the upper half-plane $-\infty < x < \infty$, $0 \leq t < \infty$, such that for each fixed $t \geq 0$, u_t , u_x , u_{xx} , u_{xt} , and u_{tt} are square-integrable functions of x on $(-\infty, \infty)$ and $\lim_{|x| \rightarrow \infty} u_x(x, t) u_t(x, t) = 0$.

(a) Show that $G(t) = \int_{-\infty}^{\infty} \left\{ [u_t(x, t)]^2 + [u_x(x, t)]^2 \right\} dx$ is a decreasing function for $t \geq 0$.

(b) Give a physical interpretation of the result in part (a).

(u)
$$\begin{aligned} G'(t) &= \frac{d}{dt} \int_{-\infty}^{\infty} \left\{ [u_t(x, t)]^2 + [u_x(x, t)]^2 \right\} dx \\ &= \int_{-\infty}^{\infty} \frac{\partial}{\partial t} \left\{ [u_t(x, t)]^2 + [u_x(x, t)]^2 \right\} dx \\ &= 2 \int_{-\infty}^{\infty} \left[u_t(x, t) u_{tt}(x, t) + u_x(x, t) u_{tx}(x, t) \right] dx \end{aligned}$$

3 pts to here.

But
$$\int_{-\infty}^{\infty} u_x(x, t) u_{tx}(x, t) dx = \lim_{\substack{M \rightarrow \infty \\ N \rightarrow -\infty}} \int_N^M \overbrace{u_x(x, t)}^u \overbrace{u_{tx}(x, t)}^{dV} dx \quad (\text{Integrate by parts})$$

$$= \lim_{\substack{M \rightarrow \infty \\ N \rightarrow -\infty}} \left(u_x(x, t) u_t(x, t) \Big|_{x=N}^M - \int_N^M u_t(x, t) u_{xx}(x, t) dx \right)$$

3 pts. to here.

$$= - \int_{-\infty}^{\infty} u_t(x, t) u_{xx}(x, t) dx.$$

Therefore,
$$G'(t) = 2 \int_{-\infty}^{\infty} \left[u_t(x, t) u_{tt}(x, t) - u_t(x, t) u_{xx}(x, t) \right] dx$$

$$= 2 \int_{-\infty}^{\infty} u_t(x, t) \left[u_{tt}(x, t) - u_{xx}(x, t) \right] dx$$

1 pt. to here.

$$= -2 \int_{-\infty}^{\infty} u_t(x, t) u_t(x, t) dx \quad (\text{since } u \text{ solves } u_{tt} - u_{xx} + u_t = 0)$$

(OVER)

$$\text{so } G'(t) = -2 \int_{-\infty}^{\infty} [u_t(x,t)]^2 dx \leq 0 \quad \text{and hence } G \text{ is a}$$

14 pts. to here. decreasing function (in the weak sense) on $t \geq 0$.

(b) The equation $u_{tt} - u_{xx} + u_t = 0$ models the small vibrations of a homogeneous elastic string with tension $T_0 = 1$, density $\rho_0 = 1$, and

20 pts. to here. damping constant $r = 1$. (See problem #2 on this exam.) The energy at time t of a solution $u = u(x,t)$ to $u_{tt} - u_{xx} + u_t = 0$

is

$$22 \text{ pts. to here. } E(t) = \int_{-\infty}^{\infty} \left[\frac{1}{2} \rho_0 u_t^2(x,t) + \frac{1}{2} T_0 u_x^2(x,t) \right] dx = \frac{1}{2} G(t).$$

25 pts. to here. Since damping transfers energy from the vibrating string to the environment, one expects $E(t)$ (and consequently $G(t)$) to decrease with time t .