

Role of Fluoride Doping in Low-Temperature Combustion-Synthesized ZrO_x Dielectric Films

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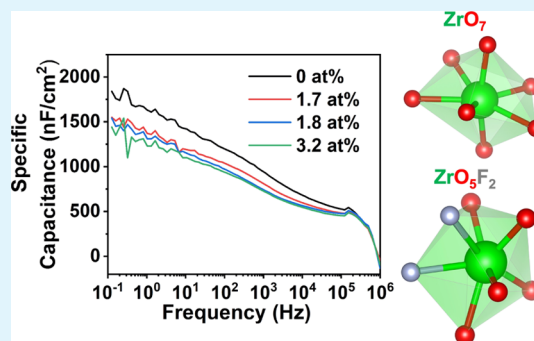
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ABSTRACT: Zirconium oxide (ZrO_x) is an attractive metal oxide dielectric material for low-voltage, optically transparent, and mechanically flexible electronic applications due to the high dielectric constant ($\kappa \sim 14\text{--}30$), negligible visible light absorption, and, as a thin film, good mechanical flexibility. In this contribution, we explore the effect of fluoride doping on structure–property–function relationships in low-temperature solution-processed amorphous ZrO_x . Fluoride-doped zirconium oxide (F:ZrO_x) films with a fluoride content between 1.7 and 3.2 in atomic (at) % were synthesized by a combustion synthesis procedure. Irrespective of the fluoride content, grazing incidence X-ray diffraction, atomic-force microscopy, and UV–vis spectroscopy data indicate that all F:ZrO_x films are amorphous, atomically smooth, and transparent in visible light. Impedance spectroscopy measurements reveal that unlike solution-processed fluoride-doped aluminum oxide (F:AlO_x), fluoride doping minimally affects the frequency-dependent capacitance instability of solution-processed F:ZrO_x films. This result can be rationalized by the relatively weak Zr–F versus Zr–O bonds and the large ionic radius of Zr^{+4} , as corroborated by EXAFS analysis and MD simulations. Nevertheless, the performance of pentacene thin-film transistors (TFTs) with F:ZrO_x gate dielectrics indicates that fluoride incorporation reduces $I\text{--}V$ hysteresis in the transfer curves and enhances bias stress stability versus TFTs fabricated with analogous, but undoped ZrO_x films as gate dielectrics, due to reduced trap density.

KEYWORDS: high- κ dielectrics, fluoride doping, zirconium oxide, combustion synthesis, metal oxides



INTRODUCTION

The past decade has witnessed major advances in large-area, low-power, mechanically flexible, and wearable electronic devices enabled using thin-film transistors (TFTs) of diverse types.^{1–5} In this regard, high-dielectric constant (κ) metal oxide dielectric materials such as aluminum oxide (AlO_x)^{5–7} and zirconium oxide (ZrO_x)^{8–12} and other unconventional dielectrics such as self-assembled nanodielectrics and composites^{1,13} hold the key to TFT operation at low driving voltages.^{14,15} Compared to conventional low- κ silicon dioxide (SiO_2 , $\kappa \sim 3.9$), high- κ oxide materials ($\kappa \sim 8.8\text{--}30$) yield significantly higher specific capacitance values for equivalent SiO_2 film thickness and thereby can allow comparable source–drain currents at far lower operating voltages, essential for low-power electronics.^{14,15} Furthermore, high- κ dielectrics can circumvent the high leakage current of ultrathin SiO_2 (<2 nm) dielectrics, thereby extending Moore’s Law.¹⁴

At present, most high- κ metal oxide dielectric films are deposited using vapor-phase-based techniques such as atomic layer deposition,^{16–19} sputtering,^{20–22} and pulsed laser deposition.^{23,24} Although vapor-phase-based methods produce high-quality films, they are time-inefficient and capital- and energy-intensive.¹⁵ In contrast, solution-based techniques are

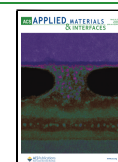
more cost-effective and compatible with roll-to-roll manufacturing and employ less-toxic/reactive chemical precursors.^{14,25,26} However, achieving complete precursor conversion to metal oxides having high chemical purity and good film densification typically requires post-film deposition annealing at temperatures of >300–400 °C.^{27,28} Using lower processing temperatures for solution-processed dielectric films often results in unstable capacitance–frequency characteristics, particularly at low frequencies, due to the presence of impurities, mobile ions (e.g., H^+), and dipolar groups (e.g., –OH).^{29–31}

Recently, this laboratory utilized fluoride doping to stabilize the capacitance at low frequencies of low-temperature solution-processed AlO_x dielectric films.³² Upon incorporation of as little as ~ 3.7 atomic (at) % of fluoride in AlO_x , the specific capacitance of F:AlO_x films, annealed at only 120 °C for 1 min

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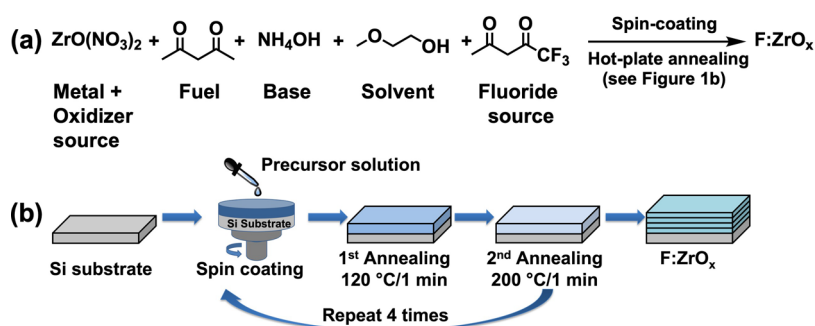


Figure 1. (a) Precursor chemical structures and their role in the combustion synthesis processing of F:ZrO_x. (b) Schematic of the spin-coating procedure for the fabrication of F:ZrO_x thin films.

followed by 300 °C for 1 min, remains stable at $\sim 166 \pm 11$ nF/cm² over a wide frequency range (10^{-1} to 10^4 Hz) compared to that of the undoped AlO_x samples, which changes from 781 ± 85 nF/cm² at 10^{-1} Hz to 104 ± 4 nF/cm² at 10^4 Hz. As shown by solid-state NMR spectroscopy, fluoride doping achieves capacitance stabilization by generating AlOF, which strongly reduces the mobile hydrogen content, suppressing polarization mechanisms at low frequencies. When employed for TFT fabrication and when utilizing both p-type and n-type semiconductors, this capacitance stability translates into reliable TFT performance with minimum *I*–*V* hysteresis, therefore enabling low-temperature solution-processed, but stable, device performance. Nevertheless, the dielectric constant of AlO_x is among the lowest (6.1–9.7) of high- κ metal oxide dielectrics.¹⁵

Zirconium oxide (ZrO_x) is a widely used high- κ oxide dielectric with a high dielectric constant ($\kappa \sim 14$ – 30) and large band gap ($E_g \sim 5.5$ eV), making it an attractive candidate for low-power electronics.^{15,33,34} However, to ensure its applicability in solution-processed TFTs, the frequency-dependent capacitance stability of solution-processed ZrO_x films should be investigated and possibly addressed. In this contribution, we investigate structure–property–function relationships in fluoride-doped ZrO_x (F:ZrO_x) films. Specifically, F:ZrO_x films with different fluoride doping contents (0–3.2 at %) are synthesized using low-temperature combustion techniques using zirconyl nitrate (ZrO(NO₃)₂) and 1,1,1-trifluoroacetylacetone (FACAc) as the Zr and fluoride sources, respectively. X-ray photoelectron spectroscopy (XPS) data corroborate fluoride incorporation and that this element ionically binds to Zr. Atomic force microscopy (AFM) and grazing incidence X-ray diffraction (GIXRD) measurements reveal that F:ZrO_x films are atomically smooth and amorphous, while UV–Vis spectroscopy confirms optical transparency in the visible region, independent of the fluoride content. Impedance spectroscopic measurements reveal that unlike AlO_x, fluoride doping does not affect capacitance instability at low frequencies in ZrO_x. It will be seen that the aforementioned observations can be convincingly explained by the low bond energy of the Zr–F bonds versus Zr–O bonds and the unchanged Zr coordination upon fluoride doping, as supported by the extended X-ray absorption fine structure (EXAFS) and density function theory (DFT) + molecular dynamics (MD) simulations. Nevertheless, F[–] doping of ZrO_x reduces hysteresis and increases bias stability in pentacene TFTs due to the reduction of interfacial charge density. Thus, this work provides an informative picture of how F[–] influences

the structure, dielectric properties, and TFT performance in amorphous ZrO_x (a-ZrO_x) films.

RESULTS AND DISCUSSION

Combustion Synthesis and Composition of Fluoride-Doped Zirconium Oxide Films. To study the effect of fluoride doping in ZrO_x solution-based combustion synthesis of metal oxides was utilized for F:ZrO_x film growth. This method is known to produce oxide thin films of superior electronic properties due to the strong exotherm of the reaction lowering the required annealing time and temperature for the film growth process.^{27,32,35} In this work, zirconyl nitrate (ZrO(NO₃)₂) was used as the Zr and oxidizer source, with acetylacetone (AcAcH) as the fuel, ammonium hydroxide (NH₄OH) as the base, and 2-methoxyethanol (2-Me) as the solvent (Figure 1a).²⁷ Note that small amounts of water [9.09% (v/v)] were used to enhance the solubility of ZrO(NO₃)₂ in 2-methoxyethanol (see the Experimental Section for details). After screening several fluoride sources including methyl 4,4,4-trifluoroacetylacetone (MFACAc) and 1,1,1-trifluoroacetylacetone (FACAc), the latter was selected as it enables higher fluoride incorporation in the ZrO_x matrix (Table S1). Viscosity values of the combustion precursor solutions without and with 20 wt % FACAc are 2.28 and 2.32 cP, respectively, which are dominated by the solvent (2-methoxyethanol). For F:ZrO_x film fabrication, the precursor solutions (Figure 1a) [ZrO(NO₃)₂ + *x* wt % FACAc, *x* = 0, 5, 10, 20] were spin-coated on n⁺⁺ Si substrates, and the resulting films were subjected to a two-step fast annealing procedure (Figure 1b) consisting of a first step that removes most of the solvent from the spin-coated film precursor, thus increasing the efficiency of heat generation in the second step.^{35,36} The second step is fast and was carried out at greatly reduced annealing times and temperatures versus conventional combustion synthesis (~ 0.5 – 1.0 h, ~ 250 – 350 °C),^{27,28} therefore enabling substantial fluoride incorporation in the ZrO_x matrix.³⁵ Specifically, here, the annealing time and temperature were carefully optimized to 120 °C/1 min (first step) + 200 °C/1 min (second step) since increasing these parameters further dramatically reduces fluoride incorporation (Table S2). This is likely due to HF elimination by the reaction of F:ZrO_x with atmospheric humidity.^{36,37} A similar reaction has been documented for other metal oxides including fluoride-doped In₂O₃.^{37,38} The entire process was repeated four times to yield ~ 19 nm-thick F:ZrO_x films (Figure 1b).

XPS measurements indicate that when the FACAc quantity in the precursor solution is increased from 0 to 20 wt % [vs the total weight of ZrO(NO₃)₂], the fluoride content in the ZrO_x

Table 1. Fluoride Incorporation in the F:ZrO_x Films and Related Properties

no.	FACAc [wt %] ^a	F ⁻ [at %] ^b	R _q [nm] ^c	avg. ρ _m [g/cc] ^d	E _g [eV] ^e
1	0	0	0.328	3.48	5.53
2	5	1.7	0.298	3.21	5.63
3	10	1.8	0.336	2.85	5.58
4	20	3.2	0.354	3.03	5.56

^aVersus total weight of ZrO(NO₃)₂. ^bFrom XPS, calculated by comparing the F 1s peak with Zr 3d, O 1s, and C 1s peaks, utilizing Advantage software. ^cRMS surface roughness (R_q) determined by AFM. ^dMass density (ρ_m) determined by XRR. ^eDirect band gap (E_g) determined by the Tauc plot method from UV-vis spectra.

matrices increases from 0 to 3.2 at % (Table 1 and Figure 2a). Note that addition of FACAc to the precursor solution beyond 20 wt % does not significantly increase the fluoride concentration in the ZrO_x films (e.g., 30 wt % FACAc affords ~3.1 at % F:ZrO_x). This is presumably due to the solubility limit of the fluoride anion in ZrO_x matrices, which is ubiquitous in various solid-state systems.^{39,40} From XPS measurements, the binding energy of F 1s is ~685.15 eV. This indicates that F is present as fluoride bonded to Zr (Figure 2a) since the binding energy of ionic fluorides is 684–685.5 eV, whereas that of organic fluorides is 688–689 eV.⁴¹ The carbon content under the optimized annealing conditions (120 °C/1 min + 200 °C/1 min) is around 2.62–3.56 at % (Table S3). Next, XPS was further employed to study metal oxide lattice formation in the F:ZrO_x films. To this purpose, the O1s binding region was scanned, and the oxygen spectra were deconvoluted into three distinct oxygen environments: (1) metal oxide lattice (M–O–M) at 529.9 ± 0.1 eV; (2) metal hydroxide (M–OH) species at 531.1 ± 0.1 eV; and (3) weakly bound CO₂ and H₂O, species (M–OR) at 532.2 ± 0.1 eV.^{36,42} The metal oxide lattice (M–O–M) content is associated with MO matrix densification, whereas metal hydroxide (M–OH) and weakly bound CO₂ and H₂O species (M–OR) relate to porosity, trapping sites, and chemical impurities. Therefore, the M–O–M content (η_{M–O–M}), given by the ratio between the O 1s M–O–M area and the total O 1s peak area, is a figure of merit for the metal oxide quality. Figures 2b and S1 and Table S4 report the M–O–M, M–OH,

and M–OR contents in the present F:ZrO_x films. As shown in Figure 2b, the η_{M–O–M} in undoped ZrO_x is around 67.8%, consistent with other reports.^{12,35,36} Upon fluoride doping, η_{M–O–M} decreases almost linearly from 64.4% for 1.7 at % F:ZrO_x to 63.9% for 1.8 at % F:ZrO_x and to 62.4% for 3.2 at % F:ZrO_x.

Thin-Film Morphology, Microstructure, and Optical Properties. Next, the film microstructure and morphology were investigated by GIXRD and AFM. GIXRD reveals that under the present annealing conditions (120 °C/1 min + 200 °C/1 min), all the F:ZrO_x films are amorphous (Figure 2c), which is an important characteristic of a dielectric to avoid interfacial defects in the semiconductor, minimize leakage current, and enhance dielectric strength.⁴³ AFM imaging (Figure 2d and Table 1) indicates that the RMS surface roughness (R_q) of the undoped ZrO_x film is ~0.328 nm. Upon fluoride incorporation, R_q remains in a narrow range of 0.298 nm for 1.7 at % F:ZrO_x, 0.336 nm for 1.8 at % F:ZrO_x, and 0.354 nm for 3.2 at % F:ZrO_x. Therefore, atomically smooth (R_q < 1 nm) fluoride-doped ZrO_x thin films were obtained utilizing combustion synthesis, which is consistent with previous combustion synthesis results.^{27,32}

Next, X-ray reflectivity (XRR) measurements were carried out to assess the thickness and the mass density of the F:ZrO_x films. The thickness was calculated from the first derivative of the electron density profile. The average electron density (ρ_e) was obtained from the area under the curve of electron density vs thickness graph and then dividing it by the film thickness. The average electron density (ρ_e) can be converted to mass density (ρ_m) by eq 1 where *M* is the molar mass, *N_A* is Avogadro's

$$\rho_m = \frac{\rho_e M}{N_A N_e} \quad (1)$$

number, and *N_e* is the total number of electrons in the chemical species. Figure S2, Tables 1 and S5 report the XRR data including the electron density profile, thickness, and mass density of the F:ZrO_x films. From Figure S2 and Table S5, the thickness of all F:ZrO_x samples is ~18–19 nm. The resulting mass density of the undoped ZrO_x film is 3.48 g/cc (Table 1), which is considerably lower than that of combustion-synthesized ZrO_x (4.45 g/cc) reported in the literature and processed at a temperature of 300 °C/1 h.²⁷ Thus, the current

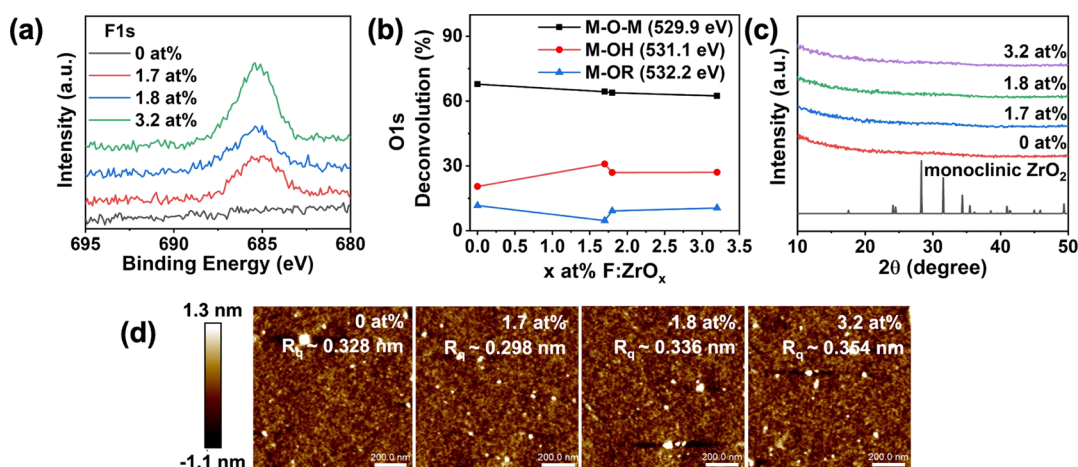


Figure 2. (a) F1s XPS spectra, (b) deconvoluted O1s environment, (c) GIXRD patterns, and (d) AFM images of the indicated *x* at % F:ZrO_x films.

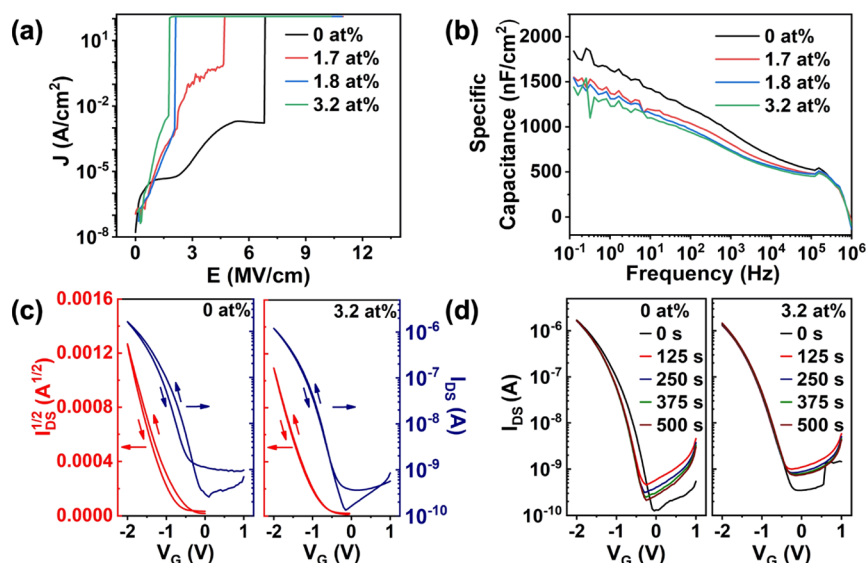


Figure 3. (a) Leakage current density–electric field (J – E) characteristics and (b) specific capacitance–frequency (C – f) characteristics of the indicated x at % F:ZrO _{x} dielectrics. (c) Transfer curves ($V_{DS} = -2$ V) and (d) bias stability ($V_{DS} = -2$ V, $V_{G, stress} = -1$ V) of pentacene (P5)-TFTs based on the indicated x -x at % F:ZrO _{x} dielectrics.

procedure enables fluoride incorporation at the expense of substantial densification of the oxide matrix.²⁷ Upon fluoride doping, the mass density falls further to 3.21 g/cc for 1.7 at % F:ZrO _{x} , to 2.85 g/cc for 1.8 at % F:ZrO _{x} , and to 3.03 g/cc for 3.2 at % F:ZrO _{x} (Table 1), which can be explained by the MD simulation (vide infra). Regarding the changes in AFM images and XRR data between 1.7 and 1.8 at % F:ZrO _{x} , these could be rationalized by the very different amounts of reactant (AcAcF) decomposing (5 vs 10 wt %), which could very reasonably affect the film morphology and density.

Finally, the optical absorption characteristics of the F:ZrO _{x} films were investigated by UV–vis spectroscopy (Figure S3 and Table 1). The optical transmittance spectra of F:ZrO _{x} films in Figure S3a indicate that all samples are highly transparent (>90%) in the visible light region (380–750 nm). Therefore, fluoride doping does not significantly diminish the optical transparency of ZrO _{x} in the visible light region, making it suitable for transparent electronics. The resulting band gap (E_g) values were calculated utilizing Tauc's method (Figure S3b and Table 1). The E_g of ZrO _{x} is 5.53 eV, in agreement with DFT calculations (vide infra) and literature reports.²⁷ Upon fluoride incorporation, E_g slightly changes to 5.63 eV for 1.7 at % F:ZrO _{x} and then falls to 5.58 eV for 1.8 at % F:ZrO _{x} and to 5.56 eV for 3.2 at % F:ZrO _{x} (Table 1).

Dielectric Properties. The dielectric properties of F:ZrO _{x} were evaluated utilizing a metal–insulator–semiconductor (MIS) architecture where the F:ZrO _{x} films (~19 nm) were spin-coated on n⁺⁺ Si wafers, and the device was completed by depositing 200 × 200 μm thermally evaporated Au electrodes through a shadow mask (Figure S4). At first, the electric field (E) dependence of the leakage current density (J) of F:ZrO _{x} films was investigated. As shown in Figure 3a, the breakdown electric field of undoped ZrO _{x} is ~6.6 MV/cm, in good agreement with literature data (~3.5–9.5 MV/cm).^{8,15,28}

Upon fluoride doping, the breakdown electric field decreases from ~4.8 MV/cm for 1.7 at % F:ZrO _{x} to ~2.0 MV/cm for 1.8 at % F:ZrO _{x} and to ~1.9 MV/cm for 3.2 at % F:ZrO _{x} (Table S6). The leakage current density (J) of undoped ZrO _{x} at 1 MV/cm is ~4.36 × 10⁻⁶ A/cm² (Figure 3a), and when the

fluoride concentration in the ZrO _{x} matrices increases, it falls to 5.69 × 10⁻⁶ A/cm² for 1.7 at % F:ZrO _{x} , to 9.20 × 10⁻⁶ A/cm² for 1.8 at % F:ZrO _{x} , and to 2.17 × 10⁻⁵ A/cm² for 3.2 at % F:ZrO _{x} (Table S6). Erosion of the dielectric strength upon fluoride incorporation is likely the result of the reduced M–O–M content, as assessed by XPS and lower densification as assessed by XRR (vide infra).^{14,44}

Next, capacitance–frequency (C – f) plots of F:ZrO _{x} MIS devices were measured in the frequency range from 10⁻¹ to 10⁶ using an oscillating frequency of 50 mV and a bias voltage of 2 V. As shown in Figure 3b, the specific capacitance of undoped ZrO _{x} increases almost linearly from 528 nF/cm² in the 10⁵ Hz region to 955 nF/cm² in the 10³ Hz region, to 1430 nF/cm² in the 10 Hz region, and to 1930 nF/cm² in the 10⁻¹ Hz region, in complete agreement with the literature report for films grown by sol–gel methods.⁴⁵ Upon fluoride doping, the present 1.7–3.2 at % F:ZrO _{x} films exhibit a similar trend, with C increasing almost linearly from 454–484 nF/cm² in the 10⁵ Hz region to 738–822 nF/cm² in the 10³ Hz region, to 1100–1200 nF/cm² in the 10 Hz region, and to 1410–1550 nF/cm² in the 10⁻¹ Hz region. The dielectric constant of 0, 1.7, 1.8, and 3.2 at % F:ZrO _{x} at 1000 Hz is 28.67, 26.79, 23.48, and 24.06, respectively. Thus, F⁻ doping does not affect κ , as expected from the C – f plots. This result is consistent with our computational results (vide infra) since the calculated real part of the dielectric function, ϵ_0 , is 2.44 and 2.39 for amorphous undoped and fluoride-doped ZrO₂ (average of 10 realizations at optimal density), respectively. Finally, C – V plots of x at % F:ZrO _{x} are shown in Figure S6, which also shows unchanged C – V characteristics with fluoride doping, in excellent agreement with C – f data (Figure 3a).

As a control experiment, capacitance–frequency (C – f) plots of combustion-synthesized F:ZrO _{x} films deposited at higher temperatures (120 °C/1 min + 300 °C/1 min) were measured. Note that due to increased annealing temperature, only very less (0.8 at %) fluoride ions were incorporated in the ZrO _{x} matrices (vide supra; Table S2). Therefore, as expected, both ZrO _{x} and 0.8 at % F:ZrO _{x} films exhibit almost identical C – f responses (Figure S5) due to low fluoride incorporation.

Table 2. Pentacene (P5) TFT Device Performance Data Based on the Indicated a -x at % F:ZrO_x Gate Dielectric Films

dielectric material	I_{on} [A] ^a	I_{off} [A] ^a	$I_{\text{on}}/I_{\text{off}}$	V_{T} [V] ^a	μ [cm ² V ⁻¹ s ⁻¹] ^a	SS [V/dec] ^a	N_{it} ($\times 10^{13}$) [V ⁻¹ cm ⁻²] ^a
ZrO _x	$1.20 \times 10^{-6} \pm 4.70 \times 10^{-7}$	$1.98 \times 10^{-10} \pm 9.15 \times 10^{-11}$	103	-0.75 ± 0.04	0.075 ± 0.028	-0.29 ± 0.02	3.51 ± 0.25
3.2 at % F:ZrO _x	$1.36 \times 10^{-6} \pm 1.79 \times 10^{-7}$	$1.47 \times 10^{-10} \pm 3.32 \times 10^{-11}$	10 ³	-0.80 ± 0.02	0.120 ± 0.013	-0.30 ± 0.02	2.81 ± 0.18

^aAverage of ~ 5 devices.

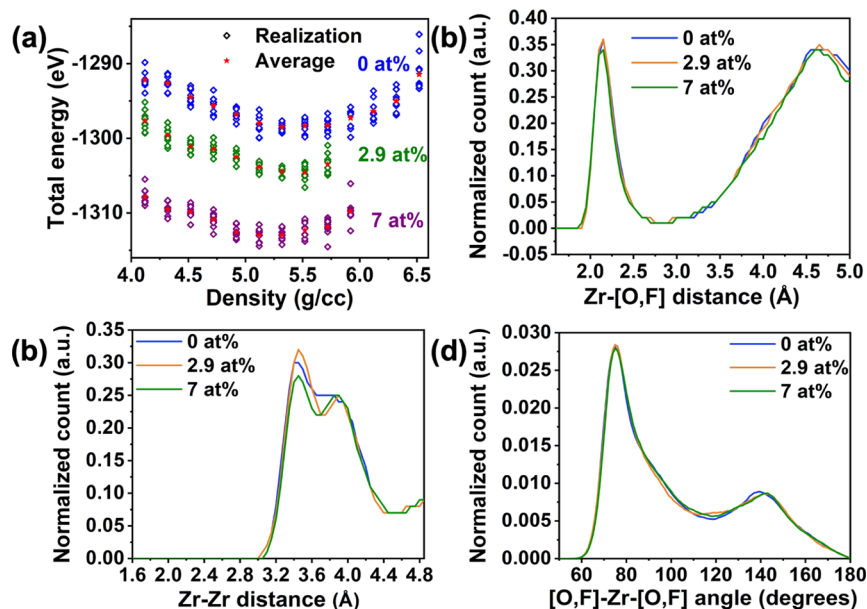


Figure 4. (a) Total energy of undoped and F⁻-doped a -ZrO₂ as calculated from *ab initio* molecular dynamics simulations as a function of density. At least 10 MD liquid quenching realizations were obtained for each density; each diamond symbol represents total energy average over 500 MD steps at 300 K, whereas the star is the average value of 10 realizations. (b,c) Distance and (d) angle distributions for undoped and F⁻-doped a -ZrO₂ as calculated from *ab initio* molecular dynamics simulations for the structures at optimal density. The results are based on 10 MD liquid quenching realizations and represent an average over 3000 MD steps at 300 K in each case.

The dielectric response of the present samples contrasts that of fluoride-doped AlO_x (F:AlO_x) films where fluoride doping stabilizes the capacitance over a wide range of frequencies (10⁻¹ to 10⁶ Hz) compared to undoped AlO_x.³² This result can be explained by the relatively weak Zr–F bond energy (627.1 kJ/mol) compared with that of the Zr–O bond (766.1 kJ/mol),⁴⁶ implying that F-doping replaces stronger Zr–O bonds with weaker Zr–F bonds.⁴⁶ Thus, this phenomenon acts as an impurity in the F:ZrO_x system and has no effect on the C – f responses. Note that Al–F bonds (675 kJ/mol) are stronger than Al–O bonds (501.9 kJ/mol),⁴⁶ meaning that F⁻ doping replaces the weaker Al–O bonds with stronger Al–F bonds, which suppresses H⁺ mobility and defects, thereby enhancing dielectric properties.³² Compositionally, with fluoride doping in AlO_x, the electron density falls very little from 0.89 e/Å³ (AlO_x) to 0.88 e/Å³ (F:AlO_x).³² In contrast, in F:ZrO_x, the electron density falls more drastically from 0.954 e/Å³ for ZrO_x to 0.829 e/Å³ for 3.2 at % F:ZrO_x (Table S5). This can be convincingly explained by MD simulation results (*vide infra*), and may rationalize the different C – f responses and dielectric breakdown behaviors in F:ZrO_x versus F:AlO_x.

TFT Measurements. Next, to evaluate the applicability of low-temperature combustion-synthesized F:ZrO_x films as gate dielectrics in TFTs, bottom-gate top-contact TFTs with 0 at % (undoped) and 3.2 at % (largest F content) F:ZrO_x dielectrics were fabricated. Since these F:ZrO_x films were annealed at low temperature and metal oxide TFTs typically require higher

annealing temperatures (>250 °C) and times (>20 min) for optimum device performance,^{42,47} a semiconductor that can be deposited at room temperature was employed. Thus, p-type organic semiconductor pentacene (P5) films (~ 30 nm thick) were vacuum-deposited at room temperature on n⁺-Si/19 nm ZrO_x and n⁺-Si/3.2 at % F:ZrO_x platforms serving as the gate electrode and the gate dielectric, respectively. Vacuum deposition at room temperature also prevents fluoride diffusion into the semiconductor. On top of the P5, 40 nm-thick gold source and drain contacts were thermally deposited through a shadow mask, yielding a device channel length (L) of 50 μm and width (W) of 1000 μm (Figure S7). Figures 3c and S8 show representative transfer and output curves, respectively. As shown in Figure S8, clear saturation behavior can be observed after -1 V V_{DS} . From transfer curves (Figure 3c), on-current (I_{on}), off-current (I_{off}), threshold voltage (V_{T}), carrier mobility (μ), subthreshold swing (SS), and interfacial charge density (N_{it}) were extracted and are tabulated in Table 2. In the case of both ZrO_x- and F:ZrO_x-based P5 TFTs, I_{on} (1.20×10^{-6} A for ZrO_x and 1.36×10^{-6} A for 3.2 at % F:ZrO_x) and I_{off} (1.98×10^{-10} A for ZrO_x and 1.47×10^{-10} A for 3.2 at % F:ZrO_x) remain near-identical, yielding identical $I_{\text{on}}/I_{\text{off}}$ ratios ($\sim 10^3$). The hole mobility increases slightly from 0.075 cm²/Vs for ZrO_x-based P5 TFTs to 0.12 cm²/Vs for F:ZrO_x-based P5 TFTs. Note that the mobility values of P5 TFTs obtained here are in agreement with literature values using oxide dielectrics and vapor-depositing P5 at room temperature.³² Also, V_{T} values

Table 3. EXAFS Fitting Results of Experimental Samples^a

samples	CN (1st shell)	CN (2nd shell)	R [Å] (1st shell)	R [Å] (2nd shell)	σ^2 (1st shell)	σ^2 (2nd shell)
a-ZrO _x	7.68 ± 0.89	5.85 ± 3.85	2.253 ± 0.037	3.579 ± 0.067	0.0091 ± 0.0016	0.015 ± 0.0055
a-3.2 at % F:ZrO _x	7.53 ± 0.90	5.92 ± 3.88	2.257 ± 0.038	3.584 ± 0.068	0.0091 ± 0.0016	0.015 ± 0.0065

^aCN = coordination number, R = bond distance, and σ^2 = mean square disorder.

increase from -0.75 V for ZrO_x to -0.8 V for 3.2 at % F:ZrO_x (Table 2). Interestingly, the ZrO_x-based P5 TFTs exhibit anticlockwise hysteresis, whereas hysteresis is negligible for the 3.2 at % F:ZrO_x-based P5 TFTs. This is presumably due to reduction of both interfacial charge density (N_{it}) at the pentacene/dielectric interface [$SS \sim -0.29$ V/dec, $N_{it} \sim 3.51 \times 10^{13}$ V⁻¹ cm⁻² (ZrO_x); $SS \sim -0.30$ V/dec, 2.8×10^{13} V⁻¹ cm⁻² (3.2 at % F:ZrO_x)],⁴⁸ calculated from SS (see eqs 3 and 4 in the Experimental Section), and reduction of bulk charge density in the dielectric layer due to fluoride doping.^{48,49} As reported in the literature,⁴⁹ fluorides can form hydrogen bonds with hydroxyl groups (F··HO–M). This can passivate OH-induced trap sites, rationalizing the decrease in trap sites in the dielectric layer upon fluoride doping.⁴⁹

Finally, the F⁻-doped ZrO_x-based P5 TFTs were next subjected to a V_G constant gate bias of -1 V for 125 s intervals over 500 s to test for bias stability (Figure 3d). The 3.2 at % F:ZrO_x-based P5 TFTs exhibit significantly improved TFT bias stability, especially after 125 s (first cycle) compared to undoped ZrO_x-based P5 TFTs (vide supra).

Ab Initio MD Simulations. To understand the atomistic mechanisms of F⁻ incorporation in the high- κ oxide dielectric materials, the structural properties of undoped and F⁻-doped a-ZrO₂ obtained from *ab initio* molecular dynamics liquid quenching simulations were investigated. Specifically, total energy calculations for the structures with different mass densities show that the presence of F⁻ suppresses the total energy increase with decreasing density, Figure 4a: the energy increases by 12.1 meV/atom in a-ZrO₂, by 5.4 meV/atom in a-ZrO₂ with 2.9 at % F, and by 2.8 meV/atom in a-ZrO₂ with 7 at % F when the density is decreased by 8% from the optimal value in each case. Because the parabolic-like energy–density curve becomes shallower as the F⁻ concentration increases, a large drop in density becomes energetically favorable with F⁻ doping of a-ZrO₂. This finding is explained by the weaker Zr–F bonds as compared to the Zr–O bonds. At room temperature, the average time variance of the Zr–F bond length (3.3×10^{-3} Å²) is more than twice larger than that for Zr–O (1.3×10^{-3} Å²). These fluctuations in the Zr–F bond distances facilitate a more pronounced decrease in density in F⁻-doped a-ZrO₂ as compared to F⁻-doped amorphous Al₂O₃ with the opposite bond strengths for Al–O versus Al–F. This result is in agreement with the experimental mass density data obtained above from XRR data.

Figure 4b–d shows the calculated Zr–[O,F] distance and [O,F]–Zr–[O,F] angle distributions for a-ZrO₂ with 0, 2.9, and 7 at % F, each at optimal density (Figure 4a). Strikingly, we find that the presence of F⁻ has no effect on the local structure of Zr. The average first-shell coordination number (CN) (bond length) changes insignificantly, from 6.77 (2.21 Å) to 6.81 (2.21 Å) and to 6.68 (2.20 Å) for 0, 2.9, and 7 at % F⁻-doped structures, respectively. These values are calculated within the cutoff radius of 3 Å and are in agreement with those of prior computational work on a-ZrO₂.⁵⁰ Similarly, the Zr–Zr distance distribution varies insignificantly upon F⁻ doping, Figure 4c. The amorphous morphology remains predominantly

corner-shared (60%), whereas edge-shared and face-shared polyhedra contribute 37 and 3%, respectively, independently of the F⁻ content. Because little changes are found for the first- and second-shell structure upon F⁻ doping, the calculated band gap remains identical, ~ 5.6 eV, for undoped and F⁻-doped amorphous zirconia, in excellent agreement with our experimental measurements (Table 1).

Importantly, we find that the average Zr–F distance in F⁻-doped a-ZrO₂, 2.25 Å, is significantly longer than the Zr–F distance in crystalline ZrF₄ (2.11 Å) and also slightly longer than the Zr–O distance in a-ZrO₂ (2.21 Å). The resulting weak Zr–F bonding along with the aforementioned strong thermal variance in the Zr–F distances in a-F:ZrO₂ is likely to promote dynamic changes in the morphology of the disordered lattice that make low-density F⁻-doped configurations energetically favorable (Figure 4a), again in accordance with the observed lower densification (Table 1).

EXAFS Measurements. Next, intrigued by the DFT-MD simulation results and to understand the coordination environment of Zr in the a-F:ZrO_x films, EXAFS measurements were carried out at the Zr K-edge (18 keV) on ~ 19 nm-thick undoped ZrO_x and 3.2 at % F:ZrO_x films. For both samples, two distinct peaks were present in the pseudoradial distribution function (*p*-RDF) produced from the Fourier transform of the EXAFS absorption coefficient $\chi(k)$ (Figure S9). The first peak corresponds to the nearest neighbor Zr–O shell ($1 \text{ Å} < R < 2 \text{ Å}$), while the second smaller peak corresponds to the second-nearest neighbor Zr–Zr shell ($2.5 \text{ Å} < R < 3.5 \text{ Å}$). Details of EXAFS fittings are summarized in Figure S9 and Table 3.

As shown in Table 3, the first-shell CN is 7.68 and 7.53, whereas the second shell CN is 5.85 and 5.92 for undoped ZrO_x and 3.2 at % F:ZrO_x, respectively, therefore exhibiting almost no change in CN in both shells of ZrO_x upon F⁻ doping. In addition, the bond distance (*R*) and the variances of bond distances (σ^2) for both the first and second shell of the samples are unchanged (Table 3). This result can be rationalized by the large radius of Zr⁴⁺ (0.78 Å), which can accommodate the addition of fluoride anions without significantly changing the CN of Zr.⁵¹ These results are in agreement with MD simulations that also indicate minimal change of the nearest and next-nearest CN in ZrO_x matrices with fluoride doping.

CONCLUSIONS

This contribution reports the first detailed investigation of F⁻ doping as a means to clarify structure–property–function relationships in a-F:ZrO_x films. The precursor chemistry, fluoride source, and annealing conditions were thoroughly optimized to enable significant fluoride incorporation (0.0–3.2 at %) in the ZrO_x matrix. The F 1s binding energies from XPS reveal that the organic precursor is fully thermolyzed and the fluoride ion is ionically bounded to Zr⁴⁺. AFM, GIXRD, XRR, and UV–vis spectroscopy studies confirm that atomically smooth, amorphous, visible light-transparent F:ZrO_x samples are produced via combustion synthesis. Impedance spectroscopy

copy measurements indicate that fluoride doping does not significantly affect capacitance instability at low frequencies in low-temperature solution-processed F:ZrO_x dielectric films. This seminal observation can be rationalized by the low bond energy of Zr–F compared to Zr–O and the large ionic radius of Zr⁴⁺, further supported by EXAFS and MD simulations. Furthermore, pentacene TFTs fabricated with F:ZrO_x gate dielectric films reveal that fluoride doping reduces *I*–*V* hysteresis compared to ZrO_x dielectrics due to the reduction of trap charge density, demonstrating its applicability in TFT performance optimization. This is the first study investigating anion doping in the structure–function relationship in application-relevant high- κ ZrO_x, paving the way for high-performing optoelectronic devices in future.

EXPERIMENTAL SECTION

Combustion Precursor Solution Preparation. All chemicals were purchased from Sigma-Aldrich and Alfa-Aesar and were used without further purification. Exactly 346.86 mg of ZrO(NO₃)₂ was dissolved in 3.0 mL of water [9.09% (v/v)] and sonicated for 30 min. To this, 30 mL of 2-methoxyethanol was added, followed by the addition of 150 μ L of acetylacetone and 67.5 μ L of ammonium hydroxide solution (14.5 M), and the resulting 0.05 M solution was stirred for 12–14 h. Then, the resulting solution was filtered using 0.45 μ m PVDF filters. Approximately 1 h prior to spin coating, the fluoride precursor 1,1,1-trifluoro-2,4-pentanedione (FACAc) in 0, 5, 10, and 20 wt % to the total weight of ZrO(NO₃)₂ was added to 5 mL of the abovementioned precursor solutions. The resulting solution was filtered again utilizing 0.20 μ m PTFE filters before spin coating.

Solution viscosity was measured at 25 °C, 60 rpm, utilizing a Brookfield viscometer.

Capacitor and Transistor Fabrication and Electrical Characterization. n⁺-Si wafers were used as the gate electrode. Before spin coating, the substrates were cleaned ultrasonically in a 1:1 mixture of isopropyl alcohol and acetone three times (10 min each), followed by 5 min in an O₂ plasma. Then, the filtered precursor solutions (using a 0.2 μ m syringe PTFE membrane) were spin-coated at 3500 rpm for 30 s and then annealed on a hot plate at 120 °C for 1 min followed by 200 °C for 1 min (relative humidity ~ 33%). This process was repeated four times to achieve the desired thickness (~18–19 nm). Finally, 40 nm Au electrodes (200 μ m \times 200 μ m) were thermally evaporated utilizing a shadow mask.

Impedance characterization was performed under ambient conditions in the dark on a custom probe station using an Agilent 1500 semiconductor parameter analyzer or a Bio-Logic SP-150 analyzer.

For TFT fabrication, n⁺-Si wafer/(F⁻-doped) ZrO_x were used as the gate electrode and dielectrics, respectively. On top of that, a 30 nm pentacene (P5) film was vacuum-deposited at room temperature at a rate of ~0.02–0.03 nm s⁻¹ (~5 \times 10⁻⁶ Torr). Then, 40 nm Au source/drain electrodes were thermally evaporated to form a channel length of 50 μ m and channel width of 1000 μ m utilizing shadow masks. TFT characterization was performed under ambient light and atmosphere on a custom probe station using an Agilent 1500 semiconductor parameter analyzer. Electron mobility (μ) was calculated in the saturation regime using the following equation

$$I_{DS} = \frac{\mu CW(V_G - V_T)^2}{2L} \quad (2)$$

Here, *C* is the capacitance per unit area of the dielectric layer, *V_T* is the threshold voltage, *V_G* is the gate voltage, *I_{DS}* is the source–drain current, and *W* and *L* are the channel width and length, respectively. SS was calculated using the following equation⁵²

$$SS = \frac{dV_G}{d \log_{10} I_{SD}} \quad (3)$$

Interfacial charge density (*N_{it}*) was calculated using the following equation⁵³

$$SS = \ln 10 \frac{kT}{q} \left(1 + \frac{qN_{it}}{C_i} \right) \quad (4)$$

Here, *k* is the Boltzmann constant, *T* is the absolute temperature (*T* = 298 K here), *q* is the charge of the carrier (1.602 \times 10⁻¹⁹ C), and *C_i* is capacitance per unit area of the dielectric.

Oxide Film Characterization. AFM data were acquired using a Bruker Dimension FastScan atomic force microscope in the tapping mode.

GIXRD and XRR measurements were carried out with a Rigaku SmartLab thin-film diffraction workstation using a high intensity 9 kW copper rotating anode X-ray source, which is coupled to a multilayer optic. XRR data were fit with the Igor-based MOTOFIT package. UV–vis samples were prepared on quartz substrates, and the UV–vis spectra were acquired using a PerkinElmer LAMBDA 1050 UV–vis–NIR spectrophotometer. XPS analyses were performed on a Thermo Scientific ESCALAB 250Xi at a base pressure of 4.5 \times 10⁻¹⁰ mbar. Spectra were obtained after the surface of the film was etched with an Ar etch gun for 20 s.

Extended X-ray absorption fine structure (EXAFS) measurements were performed at the beamline SBM-D at the Advanced Photon Source at Argonne National Laboratory. Data were collected at the Zr *K*-edge (18 keV). ~19 nm film samples of amorphous ZrO_x and F⁻-doped ZrO_x on quartz substrates were measured in the fluorescence mode. Crystalline ZrO₂ powder was used as a reference. The normalized linear EXAFS absorption coefficient $\chi(k)$ was fit by the equation

$$\chi(k) = \sum_i \frac{S_0^2 N_i f_i(k)}{k R_i^2} \sin[2kR_i + \delta(k)] e^{-\frac{2R_i}{\lambda(k)}} e^{-2k^2 \sigma_i^2} \quad (5)$$

where *S₀²* is the intrinsic loss factor, $\lambda(k)$ is the electron mean free path, *N_i* and *R_i* are the CN and bond distance of the *i*th shell of the absorbing atom, respectively, *f_i(k)* and $\delta(k)$ are the backscattering amplitude and the phase shift, respectively, and $e^{-2k^2 \sigma_i^2}$ is the Debye–Waller factor—a measure of the structural disorder or variation in *R_i*. The CN, bond distances, and Debye–Waller factors were refined for the nearest neighbor (O) and second-nearest neighbor (Zr) shells to determine the effect of fluoride doping on the local structure. Data reduction and analysis were performed using the Demeter software package.⁵⁴ The initial model was obtained using FEFF simulations based on crystalline ZrO₂. The intrinsic loss factor (*S₀²* = 0.98) was determined from the fit to the crystalline reference sample with the CN fixed to the expected value for ZrO₂. Fittings were carried out over the *R*-range 1.1–3.5 Å using data in the *k*-range of 3–11 Å⁻¹. Fitting was done in *R*-space, and a *k*-weight of 2 and 3 was used.

Computational Approach. A-ZrO₂ structures with and without F⁻ doping were calculated using the Vienna Ab Initio Simulation Package (VASP).^{55–58} The calculations are based on DFT^{59,60} with periodic boundary conditions and use the PBE functional^{61,62} within the projector augmented-wave method.^{63,64} To obtain amorphous structures, we employed the *ab initio* molecular dynamics (MD) liquid quenching approach as implemented in VASP. An initial Zr₄₆O₉₂ structure with a crystalline density of 5.68 g/cm³ was melted at 5000 K to eliminate any crystalline memory. For F:ZrO₂, fluoride atoms were randomly introduced into the supercell, and the number of oxygen atoms was adjusted to zero maintain charge neutrality, resulting in Zr₄₆O₉₀F₄ (2.9 at % F) and Zr₄₆O₈₇F₁₀ (7 at % F) configurations. The structures had additional melting at 5000 K for 10 ps to randomize the multicomponent configuration and stabilize the total energy. Next, each structure was rapidly quenched to 100 K at 200 K/ps rate. An energy cutoff of 260 eV and single Γ -point were used during melting and quenching processes. Finally, each structure was equilibrated at 300 K for 6 ps with a cutoff energy of 400 eV. All MD simulations were carried out in the NVT ensemble with a Nose–Hoover thermostat using an integration time step of 2 fs. For an

accurate structural analysis of the simulated amorphous oxides (distance and angle distributions), the room temperature structures were used.

To determine the supercell density, at least 10 melt–quench cycles were performed for each density value that ranged from 4.12 to 6.52 g/cm³ for undoped a-ZrO₂. Based on the total energy density calculations, the optimal density was found to be 5.32 g/cm³ for undoped amorphous zirconia, which is about 6% lower than crystalline density.

Finally, the atomic configurations obtained from the *ab initio* MD simulations were optimized within DFT using the PBE functional. For the optimization, a cutoff energy of 500 eV and the 4 × 4 × 4 Γ -centered *k*-point were used; the atomic positions were relaxed until the Hellmann–Feynman force on each atom was below 0.01 eV/Å. The electronic and optical properties of the optimized amorphous zirconia were calculated using the hybrid Heyd–Scuseria–Ernzerhof (HSE06) approach^{65,66} with a mixing parameter of 0.25 and a screening parameter α of 0.2 Å⁻¹. Optical absorption was derived from the frequency-dependent dielectric function, $\epsilon(\omega) = \epsilon_1(\omega) + i\epsilon_2(\omega)$, calculated within independent particle approximation in VASP. The imaginary part, $\epsilon_2(\omega)$, is related to the optical absorption at a given frequency, ω , and is determined based on the electronic transitions of the hybrid functional solution. The real part of the complex dielectric function is obtained using Kramers–Kronig relations. The resulting atomic structures were plotted using VESTA software.⁶⁷

■ ASSOCIATED CONTENT

SI Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsami.1c22853>.

AFM, XRR plots and data, TFT architecture, bias stress data, MD simulations, UV–Vis data, and XPS data (PDF)

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Notes

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■ ABBREVIATIONS

XPS, X-ray photoelectron spectroscopy
AFM, atomic force microscopy
GIXRD, grazing incidence X-ray diffraction
XRR, X-ray reflectivity
TFT, thin-film transistor
MD, molecular dynamics
DFT, density functional theory
EXAFS, extended X-ray absorption fine structure
UV/Vis/NIR, ultraviolet/visible/near infrared
soln., solution
w/o, without
at %, atom %
vs, versus
wt %, weight %

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