

- 1.(25 pts.) (a) Solve the equation $\cos(x)u_x + y \sin(x)u_y = 0$ subject to $u(0, y) = e^{-y^2}$ for $-\infty < y < \infty$.
 (b) In which region of the xy -plane is the solution uniquely determined?

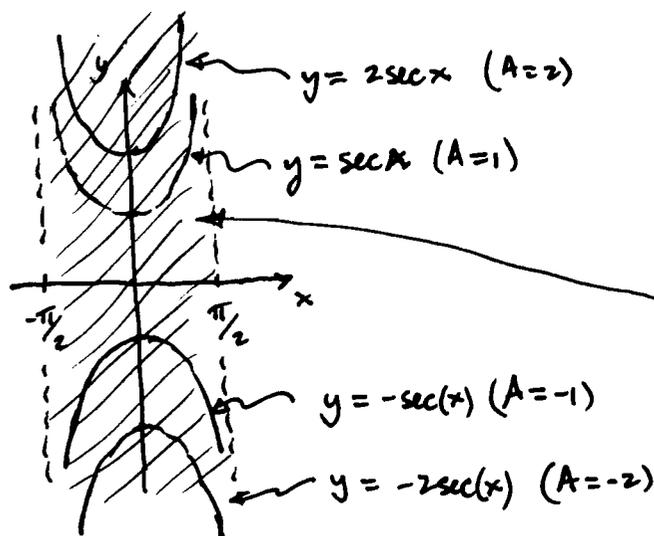
20 pts. (a) Along the characteristic curves of the PDE, described by $\frac{dy}{dx} = \frac{y \sin(x)}{\cos(x)}$, 3 pts. to here
 the solution u is constant. Then $\frac{1}{y} dy = \frac{\sin(x) dx}{\cos(x)} \Rightarrow \ln|y| = -\ln|\cos(x)|$
 $\Rightarrow \ln|y \cos(x)| = C \Rightarrow y \cos(x) = \pm e^C = A \Rightarrow y = A \sec(x)$. (Characteristic curves) 9 pts. to here.

Along a characteristic curve, $u(x, y) = u(x, A \sec(x)) = u(0, A) = f(A)$. Thus 15 pts. to here.
 the general solution is $u(x, y) = f(y \cos(x))$ where f is a C^1 -function of a single real variable. Applying the auxiliary condition yields

17 pts. to here.
 $e^{-y^2} = u(0, y) = f(y \cos(0)) = f(y)$ for all real y .

20 pts. to here.
 $\therefore u(x, y) = f(y \cos(x)) = e^{-y^2 \cos^2(x)}$

5 pts. (b)



The characteristic curves "give" the strip,

$S: -\pi/2 < x < \pi/2, -\infty < y < \infty$.

Therefore the solution is uniquely determined in S .

$$B^2 - 4AC = (-3)^2 - 4(1)(-4) > 0$$

2.(25 pts.) Consider the partial differential equation $u_{xx} - 3u_{xy} - 4u_{yy} = 0$.

(a) Classify its order and type (linear, nonlinear, homogeneous, inhomogeneous, parabolic, etc.).

(b) Find, if possible, its general solution in the xy -plane.

(c) Find, if possible, its solution which satisfies $u(x,0) = x^3$ and $u_y(x,0) = -3x^2$ if $-\infty < x < \infty$.

5 pts. (a) The PDE is second-order, linear, homogeneous, and of hyperbolic type

10 pts. (b)
$$\left(\frac{\partial^2}{\partial x^2} - 3\frac{\partial^2}{\partial x \partial y} - 4\frac{\partial^2}{\partial y^2}\right)u = 0$$

2 pts. to here.
$$\left(\frac{\partial}{\partial x} + \frac{\partial}{\partial y}\right)\left(\frac{\partial}{\partial x} - 4\frac{\partial}{\partial y}\right)u = 0$$

$$\therefore 5\frac{\partial}{\partial \eta}\left(5\frac{\partial}{\partial \xi}\right)u = 0$$

$$\Rightarrow \frac{\partial}{\partial \eta}\left(\frac{\partial u}{\partial \xi}\right) = 0$$
 6 pts. to here.

$$\Rightarrow \frac{\partial u}{\partial \xi} = \int 0 d\eta + c_1(\xi)$$

$$\Rightarrow u = \int c_1(\xi) d\xi + c_2(\eta)$$

$$\therefore u = f(\xi) + g(\eta)$$
 8 pts. to here.

$$\therefore u(x,y) = f(x-y) + g(4x+y)$$
 where f and g are C^2 -functions of a single real variable. 10 pts. to here.

10 pts. (c)
$$u_y(x,y) = -f'(x-y) + g'(4x+y)$$
 so

(*)
$$-3x^2 = u_y(x,0) = -f'(x) + g'(4x)$$
 for all real x . 2 pts. to here

(†)
$$x^3 = u(x,0) = f(x) + g(4x)$$
 " " " "

(**)
$$\Rightarrow 3x^2 = f'(x) + 4g'(4x)$$
 " " " " (OVER)
4 pts. to here.

4 pts. to here.
Let $\xi = \beta x - \alpha y = x - y$
 $\eta = -(\delta x - \gamma y) = -(4x - y) = 4x + y$
$$\frac{\partial}{\partial x} = \frac{\partial \xi}{\partial x} \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial x} \frac{\partial}{\partial \eta} = \frac{\partial}{\partial \xi} + 4\frac{\partial}{\partial \eta}$$

$$\frac{\partial}{\partial y} = \frac{\partial \xi}{\partial y} \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial y} \frac{\partial}{\partial \eta} = -\frac{\partial}{\partial \xi} + \frac{\partial}{\partial \eta}$$

$$\therefore \frac{\partial}{\partial x} + \frac{\partial}{\partial y} = 5\frac{\partial}{\partial \eta}$$

and
$$\frac{\partial}{\partial x} - 4\frac{\partial}{\partial y} = 5\frac{\partial}{\partial \xi}$$

Adding (*) and (**) yields $0 = 5g'(x)$ for all real $x \Rightarrow g(x) = c$ 6 pts. to here
for all real x . Substituting in (†) gives $x^3 = f(x) + c$. 8 pts. to here.

$$\therefore u(x, y) = f(x-y) + g(x+y) = (x-y)^3 - c + c$$

$$\boxed{u(x, y) = (x-y)^3}$$

10 pts. to here.

3.(25 pts.) A homogeneous solid material occupying $D = \{(x, y, z) \in \mathbb{R}^3 : 4 \leq x^2 + y^2 + z^2 \leq 100\}$ is completely insulated and its initial temperature at position (x, y, z) in D is $200/\sqrt{x^2 + y^2 + z^2}$.

(a) Write (without proof or derivation) the partial differential equation and initial/boundary conditions that completely govern the temperature $u(x, y, z, t)$ at position (x, y, z) in D and time $t \geq 0$.

(b) Use Gauss' divergence theorem to help show that the heat energy $H(t) = \iiint_D c\rho u(x, y, z, t) dV$ of the material in D at time t is a constant function of time. Here c and ρ denote the (constant) specific heat and mass density, respectively, of the material in D .

Bonus (10 pts.). Compute the (constant) steady-state temperature that the material in D reaches after a long time.

15 pts. (a)
$$\begin{cases} u_t - k\nabla^2 u = 0 & \text{if } 4 < x^2 + y^2 + z^2 < 100 \text{ and } 0 < t < \infty, & (5) \\ \frac{\partial u}{\partial n} = 0 & \text{if } t \geq 0 \text{ and } x^2 + y^2 + z^2 = 4 \text{ or } x^2 + y^2 + z^2 = 100, & (5) \\ u(x, y, z, 0) = \frac{200}{\sqrt{x^2 + y^2 + z^2}} & \text{if } 4 \leq x^2 + y^2 + z^2 \leq 100. & (5) \end{cases}$$

10 pts. (b)
$$\begin{aligned} \frac{dH}{dt} &= \iiint_D \frac{\partial}{\partial t} (c\rho u(x, y, z, t)) dV = \iiint_D c\rho \frac{\partial u}{\partial t} dV = \iiint_D c\rho k \nabla^2 u dV = \iiint_D c\rho k \nabla \cdot (\nabla u) dV \\ &= c\rho k \iint_{\partial D} \nabla u \cdot \vec{n} dS = c\rho k \iint_{\partial D} \frac{\partial u}{\partial n} dV = 0. \end{aligned}$$
 Therefore $H(t) = \text{constant}$ for all $t \geq 0$.
5 pts. to here.
10 pts. to here.

10 pts. Bonus: Let $U_0 = \lim_{t \rightarrow \infty} u(x, y, z, t)$ be the steady-state temperature of the material in D .

Then $H(\infty) = \lim_{t \rightarrow \infty} H(t) = \lim_{t \rightarrow \infty} \iiint_D c\rho u(x, y, z, t) dV = \iiint_D c\rho U_0 dV = c\rho U_0 \cdot \text{vol}(D)$.
1 pt. to here.

But $\text{vol}(D) = \frac{4}{3}\pi(R_2^3 - R_1^3) = \frac{4}{3}\pi(10^3 - 2^3) = \frac{4 \cdot 992\pi}{3}$ and $H(0) = \iiint_D c\rho u(x, y, z, 0) dV$.
4 pts. to here.

$$= \iiint_D c\rho \frac{200}{\sqrt{x^2 + y^2 + z^2}} dV = \int_0^{2\pi} \int_0^\pi \int_2^{10} c\rho \frac{200}{r} r^2 \sin\phi dr d\phi d\theta = \int_0^{2\pi} \int_0^\pi c\rho 100r^2 \Big|_{r=2}^{r=10} \sin\phi d\phi d\theta$$

$$= 9600c\rho(-\cos\phi) \Big|_{\phi=0}^{\phi=\pi} \cdot 2\pi = 38,400c\rho\pi.$$
 Thus $38,400c\rho\pi = c\rho U_0 \frac{4}{3} \cdot 992\pi$
7 pts. to here.

$$\Rightarrow U_0 = \frac{3}{4 \cdot 992} \cdot 38,400 = \frac{900}{31} \approx 29.$$
 10 pts. to here.

4. (25 pts.) Use Fourier transform methods to find a solution to

$$u_{xx} + u_{yy} \stackrel{\textcircled{1}}{=} 0 \quad \text{for } -\infty < x < \infty, 0 < y < \infty,$$

which satisfies

$$u(x, 0) \stackrel{\textcircled{2}}{=} \begin{cases} 1 & \text{if } |x| < 1, \\ 0 & \text{otherwise,} \end{cases}$$

and

$$\lim_{y \rightarrow \infty} u(x, y) \stackrel{\textcircled{3}}{=} 0 \quad \text{for each } x \text{ in } (-\infty, \infty).$$

We take the Fourier transform of $\textcircled{1}$ with respect to the variable x :

$$\mathcal{F}(u_{xx} + u_{yy})(\xi) = \mathcal{F}(0)(\xi)$$

$$(i\xi)^2 \mathcal{F}(u)(\xi) + \frac{\partial^2 \mathcal{F}(u)(\xi)}{\partial y^2} = 0$$

$$\frac{\partial^2 \mathcal{F}(u)(\xi)}{\partial y^2} - \xi^2 \mathcal{F}(u)(\xi) = 0.$$

The general solution to this second-order linear homogeneous ODE in the variable y , with ξ as a parameter, is $\mathcal{F}(u)(\xi) = c_1(\xi)e^{\xi y} + c_2(\xi)e^{-\xi y}$. We apply $\textcircled{3}$ to get

$$\lim_{y \rightarrow \infty} (c_1(\xi)e^{\xi y} + c_2(\xi)e^{-\xi y}) = \lim_{y \rightarrow \infty} \mathcal{F}(u)(\xi) = \lim_{y \rightarrow \infty} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u(x, y) e^{-i\xi x} dx = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \left(\lim_{y \rightarrow \infty} u(x, y) \right) e^{-i\xi x} dx$$

$$= 0. \quad \text{Consider first the case when } \xi > 0. \text{ If } c_1(\xi) \neq 0 \text{ then } \lim_{y \rightarrow \infty} (c_1(\xi)e^{\xi y} + c_2(\xi)e^{-\xi y}) = \lim_{y \rightarrow \infty} (c_1(\xi)e^{\xi y}) = \pm \infty, \text{ depending on the algebraic sign of } c_1(\xi). \text{ Since this}$$

contradicts the fact that the limit must be 0, it follows that $c_1(\xi) = 0$ if $\xi > 0$.

A similar argument shows that $c_2(\xi) = 0$ if $\xi < 0$. Consequently,

$$(*) \quad \mathcal{F}(u)(\xi) = \begin{cases} c_2(\xi)e^{-\xi y} & \text{if } \xi > 0 \\ c_1(\xi)e^{\xi y} & \text{if } \xi < 0 \end{cases} = c(\xi)e^{-|\xi|y}.$$

Evaluating $(*)$ at $y=0$ and applying $\textcircled{2}$ gives

$$(**) \quad c(\xi) = \mathcal{F}(u)(\xi) \Big|_{y=0} = \mathcal{F}(u(\cdot, 0))(\xi) = \mathcal{F}(x_{(-1,1)})(\xi). \quad (\text{OVER})$$

4 pts.
to here.

8 pts.
here.

12 pts.
to here.

15 pts.
to here.

By entry C in the table of Fourier transforms accompanying this exam,

$$(***) \quad \mathcal{F}\left(\frac{1}{(\cdot)^2+y^2}\right)(\xi) = \sqrt{\frac{\pi}{2}} \cdot \frac{e^{-|\xi|y}}{y} \quad \text{so} \quad \mathcal{F}\left(\sqrt{\frac{2}{\pi}} \cdot \frac{y}{(\cdot)^2+y^2}\right)(\xi) = e^{-|\xi|y}.$$

17 pts.
to here.

Substituting from (**) and (***) into (*) gives

$$\begin{aligned} \mathcal{F}(u)(\xi) &= \mathcal{F}\left(\chi_{(-1,1)}\right)(\xi) \mathcal{F}\left(\sqrt{\frac{2}{\pi}} \cdot \frac{y}{(\cdot)^2+y^2}\right)(\xi) \\ &= \frac{1}{\sqrt{2\pi}} \mathcal{F}\left(\chi_{(-1,1)} * \sqrt{\frac{2}{\pi}} \cdot \frac{y}{(\cdot)^2+y^2}\right)(\xi) \\ &= \mathcal{F}\left(\frac{1}{\pi} \chi_{(-1,1)} * \frac{y}{(\cdot)^2+y^2}\right)(\xi). \end{aligned}$$

20 pts.
to here.

Applying the inversion theorem,

$$u(x,y) = \left(\frac{1}{\pi} \chi_{(-1,1)} * \frac{y}{(\cdot)^2+y^2}\right)(x)$$

21 pts.
to here.

$$= \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{y}{(x-z)^2+y^2} \chi_{(-1,1)}(z) dz$$

$$= \frac{1}{\pi} \int_{-1}^1 \frac{y}{(x-z)^2+y^2} dz \quad \leftarrow \begin{cases} \text{Let } s = z-x. \\ \text{Then } ds = dz. \end{cases}$$

23 pts.
to here

$$= \frac{1}{\pi} \int_{-1-x}^{1-x} \frac{y}{s^2+y^2} ds$$

$$= \frac{1}{\pi} \operatorname{Arctan}\left(\frac{s}{y}\right) \Big|_{s=-1-x}^{1-x}$$

$$= \boxed{\frac{1}{\pi} \operatorname{Arctan}\left(\frac{1-x}{y}\right) + \frac{1}{\pi} \operatorname{Arctan}\left(\frac{1+x}{y}\right)}.$$

25 pts.
to here.

②-③ $u_x(0,y,z) = 0 = u_x(1,y,z)$ if $0 \leq y \leq 1, 0 \leq z \leq 1$

⑥ $u_y(x,1,z) = 0$ if $0 \leq x \leq 1, 0 \leq z \leq 1$

④-⑤ $u_z(x,y,0) = 0 = u_z(x,y,1)$ if $0 \leq x \leq 1, 0 \leq y \leq 1$

5. (25 pts.) Solve $\nabla^2 u = 0$ in the cube $0 < x < 1, 0 < y < 1, 0 < z < 1$ given that u satisfies the inhomogeneous **Dirichlet** condition $u(x,0,z) = 4 \sin^2(\pi x) \sin^2(\pi z)$ if $0 \leq x \leq 1, 0 \leq z \leq 1$, and u satisfies homogeneous **Neumann** boundary conditions on the other five faces. (Hint: You may find the identity $2 \sin^2(\theta) = 1 - \cos(2\theta)$ useful.)

1 pt. to here.

We seek nontrivial solutions of the form $u(x,y,z) = X(x)Y(y)Z(z)$ to the homogeneous portion of this problem: ①-②-③-④-⑤-⑥. Substituting into ① leads to

$$X''(x)Y(y)Z(z) + X(x)Y''(y)Z(z) + X(x)Y(y)Z''(z) = 0 \Rightarrow -\frac{X''(x)}{X(x)} = \frac{Y''(y)}{Y(y)} + \frac{Z''(z)}{Z(z)} = \lambda,$$

and then to $-\frac{Z''(z)}{Z(z)} = \frac{Y''(y)}{Y(y)} - \lambda = \mu$. Substituting in ②, ③, ④, ⑤, and ⑥ and

using the assumption that $X(x)Y(y)Z(z)$ is not identically 0 on the unit cube yields

$$X'(0) = 0 = X'(1), \quad Z'(0) = 0 = Z'(1), \quad \text{and} \quad Y'(1) = 0. \quad \text{Thus}$$

$$(*) \begin{cases} X''(x) + \lambda X(x) = 0, & X'(0) = 0 = X'(1), & 4 \text{ pts. to here.} \\ Z''(z) + \mu Z(z) = 0, & Z'(0) = 0 = Z'(1), & 7 \text{ pts. to here.} \\ Y''(y) - (\lambda + \mu)Y(y) = 0, & Y'(1) = 0. & 9 \text{ pts. to here.} \end{cases}$$

The eigenvalue problems in the first two lines of (*) have solutions

$$\lambda_l = (l\pi)^2, \quad X_l(x) = \cos(l\pi x) \quad (l = 0, 1, 2, \dots) \quad 12 \text{ pts. to here.}$$

$$\mu_m = (m\pi)^2, \quad Z_m(z) = \cos(m\pi z) \quad (m = 0, 1, 2, \dots) \quad 15 \text{ pts. to here.}$$

Substituting $\lambda + \mu = \lambda_l + \mu_m = \pi^2(l^2 + m^2)$ in the DE in the third line of (*) and applying the B.C. yields $Y_{l,m}(y) = \cosh(\pi(y-1)\sqrt{l^2 + m^2})$, up to a constant factor.

17 pts. to here.

Thus the superposition principle implies that

$$u(x,y,z) = \sum_{l=0}^{\infty} \sum_{m=0}^{\infty} A_{l,m} \cos(l\pi x) \cos(m\pi z) \cosh(\pi(y-1)\sqrt{l^2 + m^2})$$

19 pts. to here.

is a formal solution to ①-②-③-④-⑤-⑥ for any choice of constants $A_{l,m}$.

We need to choose these constants so that ⑦ is satisfied:

(OVER)

$$\sum_{l=0}^{\infty} \sum_{m=0}^{\infty} A_{l,m} \cosh(-\pi\sqrt{l^2+m^2}) \cos(l\pi x) \cos(m\pi z) = u(x, 0, z)$$

$$= 4 \sin^2(\pi x) \sin^2(\pi z)$$

$$= [1 - \cos(2\pi x)] [1 - \cos(2\pi z)]$$

$$= 1 - \cos(2\pi x) - \cos(2\pi z) + \cos(2\pi x) \cos(2\pi z)$$

must hold for all $0 \leq x \leq 1$ and $0 \leq z \leq 1$. Consequently, equating "like" coefficients produces

$$A_{0,0} = 1$$

$$A_{2,0} \cosh(-2\pi) = -1$$

$$A_{0,2} \cosh(-2\pi) = -1$$

$$A_{2,2} \cosh(-2\pi\sqrt{2}) = 1$$

$$A_{0,0} = 1$$

$$\Rightarrow A_{2,0} = \frac{1}{\cosh(2\pi)} = A_{0,2}$$

$$A_{2,2} = \frac{1}{\cosh(2\pi\sqrt{2})}$$

and all other $A_{l,m} = 0$. That is,

$$u(x, y, z) = 1 - \frac{\cos(2\pi x) \cosh(2\pi(y-1))}{\cosh(2\pi)} - \frac{\cos(2\pi z) \cosh(2\pi(y-1))}{\cosh(2\pi)} + \frac{\cos(2\pi x) \cos(2\pi z) \cosh(2\pi(y-1))}{\cosh(2\pi\sqrt{2})}$$

solves ①-②-③-④-⑤-⑥-⑦.

6. (25 pts.) (a) Show that the full Fourier series of $f(x) = x^3 - x$ on $[-1, 1]$ is $\sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi x)}{(\pi n)^3}$.

7 (b) Show that the full Fourier series of f converges uniformly to f on $[-1, 1]$.

5 (c) Use the results above to help compute the sum of the series $\sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)^3}$.

5 (d) Use Parseval's identity and the results above to help compute the sum of the series $\sum_{n=1}^{\infty} \frac{1}{n^6}$.

(a) The full Fourier series of f is $a_0 + \sum_{n=1}^{\infty} [a_n \cos(n\pi x) + b_n \sin(n\pi x)]$ where

1 pt. to here. $a_0 = \frac{\langle f, 1 \rangle}{\langle 1, 1 \rangle} = \frac{1}{2} \int_{-1}^1 \underbrace{(x^3 - x)}_{\text{odd}} dx = 0,$

2 pts. to here. $a_n = \frac{\langle f, \cos(n\pi \cdot) \rangle}{\langle \cos(n\pi \cdot), \cos(n\pi \cdot) \rangle} = \int_{-1}^1 \underbrace{(x^3 - x)}_{\text{odd}} \underbrace{\cos(n\pi x)}_{\text{even}} dx = 0$ for $n=1, 2, 3, \dots,$

3 pts. to here. $b_n = \frac{\langle f, \sin(n\pi \cdot) \rangle}{\langle \sin(n\pi \cdot), \sin(n\pi \cdot) \rangle} = \int_{-1}^1 \underbrace{(x^3 - x)}_{\text{even}} \underbrace{\sin(n\pi x)}_{\text{odd}} dx = 2 \int_0^1 \underbrace{x^3}_{u} \underbrace{\sin(n\pi x)}_{dv} dx =$

4 pts. to here $\left[2(x^3 - x) \left(-\frac{\cos(n\pi x)}{n\pi} \right) \right]_0^1 + \frac{2}{n\pi} \int_0^1 \underbrace{(3x^2 - 1)}_u \underbrace{(\cos(n\pi x))}_{dv} dx = \frac{2}{n\pi} \left[\left(3x^2 - 1 \right) \frac{\sin(n\pi x)}{n\pi} \right]_0^1 - \int_0^1 6x \frac{\sin(n\pi x)}{n\pi} dx$

5 pts. to here $= -\frac{12}{(n\pi)^2} \int_0^1 \underbrace{x}_u \underbrace{\sin(n\pi x)}_{dv} dx = -\frac{12}{(n\pi)^2} \left[x \left(-\frac{\cos(n\pi x)}{n\pi} \right) \right]_0^1 + \frac{12}{(n\pi)^2} \int_0^1 \frac{-\cos(n\pi x)}{n\pi} dx = \frac{12(-1)^n}{(n\pi)^3}$

6 pts. to here $= \frac{12(-1)^n}{(n\pi)^3}$ 7 pts. to here

$\therefore f(x) \sim \sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi x)}{(\pi n)^3}$ 8 pts. to here

(b) The set of orthogonal functions for the full Fourier series of f on $[-1, 1]$ is

$\Phi = \{ 1, \cos(\pi x), \sin(\pi x), \cos(2\pi x), \sin(2\pi x), \dots \}$, and this is the complete set of eigenfunctions for the problem $-\mathcal{L}''(x) = \lambda \mathcal{L}(x)$ with periodic boundary

conditions $\mathcal{L}(1) = \mathcal{L}(-1)$ and $\mathcal{L}'(1) = \mathcal{L}'(-1)$. The function $f(x) = x^3 - x$ is

symmetric (or hermitian) B.C.'s.

(OVER)

1 pt. to here

continuous on $[-1, 1]$ and has derivatives $f'(x) = 3x^2 - 1$ and $f''(x) = 6x$

3 pts. to here.

which are also continuous on $[-1, 1]$. Note that f satisfies the periodic

5 pts. to here.

B.C.'s on $[-1, 1]$: $f(1) = 0 = f(-1)$ and $f'(1) = 2 = f'(-1)$. Therefore

7 pts. to here.

Theorem 2 (on the "Convergence Theorems" attached to this exam) implies that

the Fourier series of f converges to f uniformly on $[-1, 1]$.

1 pt. to here.

(c) $x^3 - x = f(x) = \sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi x)}{(n\pi)^3}$ for all $-1 \leq x \leq 1$. Set

2 pts. to here

$x = 1/2$ in this identity to obtain

$$-\frac{3}{8} = \left(\frac{1}{2}\right)^3 - \frac{1}{2} = \sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi/2)}{(n\pi)^3}$$

$$= \sum_{k=0}^{\infty} \frac{12(-1)^{2k+1} \sin((2k+1)\pi/2)}{((2k+1)\pi)^3}$$

n	$\sin(n\pi/2)$
1	1
2	0
3	-1
4	0
\vdots	
$2k$	0
$2k+1$	$(-1)^k$

$$\Rightarrow -\frac{3}{8} = - \sum_{k=0}^{\infty} \frac{12(-1)^k}{[(2k+1)\pi]^3}$$

5 pts. to here.

$$\Rightarrow \boxed{\frac{\pi^3}{32}} = \left(-\frac{\pi^3}{12}\right) \left(-\frac{3}{8}\right) = \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)^3}$$

2 pts. to here.

(d) By Parseval's identity, $|a_0|^2 \int_a^b 1^2 dx + \sum_{n=1}^{\infty} (|a_n|^2 \int_a^b \cos^2(n\pi x) dx + |b_n|^2 \int_a^b \sin^2(n\pi x) dx) = \int_a^b |f(x)|^2 dx$. Applying this to our case produces $\sum_{n=1}^{\infty} \left| \frac{12(-1)^n}{(n\pi)^3} \right|^2 \int_{-1}^1 \sin^2(n\pi x) dx = \int_{-1}^1 |x^3 - x|^2 dx$

But $\int_{-1}^1 \sin^2(n\pi x) dx = \int_{-1}^1 \left(\frac{1}{2} - \frac{1}{2} \cos(2n\pi x)\right) dx = 1$ and $\int_{-1}^1 |x^3 - x|^2 dx = 2 \int_0^1 (x^3 - x)^2 dx$

$$= 2 \int_0^1 (x^2 - 2x^4 + x^6) dx = 2 \left(\frac{x^3}{3} - \frac{2x^5}{5} + \frac{x^7}{7} \right) \Big|_0^1 = 2 \left(\frac{1}{3} - \frac{2}{5} + \frac{1}{7} \right) = \frac{16}{105}$$

$$\frac{(12)^2}{\pi^6} \sum_{n=1}^{\infty} \frac{1}{n^6} = \frac{16}{105} \Rightarrow \sum_{n=1}^{\infty} \frac{1}{n^6} = \frac{16}{105} \cdot \frac{\pi^6}{16 \cdot 9} = \boxed{\frac{\pi^6}{945}}$$

5 pts. to here.

7.(25 pts.). Solve

$$u, -u_x \stackrel{\textcircled{1}}{=} 0 \text{ in } -1 < x < 1, 0 < t < \infty,$$

subject to the boundary conditions

$$u(1,t) \stackrel{\textcircled{2}}{=} u(-1,t) \text{ and } u_x(1,t) \stackrel{\textcircled{3}}{=} u_x(-1,t) \text{ if } t \geq 0,$$

and the initial condition

$$u(x,0) \stackrel{\textcircled{4}}{=} x^3 - x \text{ if } -1 \leq x \leq 1.$$

(Hint: You may find the results of problem 6 useful.)

Bonus (10 pts.). Show that the solution to the problem above is unique.

We seek nontrivial solutions of $\textcircled{1}-\textcircled{2}-\textcircled{3}$ of the form $u(x,t) = \Sigma(x)T(t)$. ^{2 pts. to here.} Substituting into $\textcircled{1}$ yields $T'(t)\Sigma(x) - T(t)\Sigma''(x) = 0 \Rightarrow \frac{T'(t)}{T(t)} = \frac{\Sigma''(x)}{\Sigma(x)} = -\lambda$. Substituting into $\textcircled{2}-\textcircled{3}$ yields $\Sigma(1)T(t) = \Sigma(-1)T(t)$ and $\Sigma'(1)T(t) = \Sigma'(-1)T(t)$ for all $t \geq 0$. But $u(x,t) = \Sigma(x)T(t)$ is not identically zero so

$$(*) \begin{cases} \Sigma''(x) + \lambda\Sigma(x) = 0, & \Sigma(1) = \Sigma(-1), \Sigma'(1) = \Sigma'(-1), \\ T'(t) + \lambda T(t) = 0. \end{cases}$$

8 pts. to here.
10 pts. to here.

^{12 pts. to here.} The eigenvalue problem in the first line of $(*)$ has solutions $\lambda = \lambda_n = (n\pi)^2$ and $\Sigma_n(x) = a_n \cos(n\pi x) + b_n \sin(n\pi x)$ ($n=0,1,2,\dots$ and a_n, b_n arbitrary constants). ^{14 pts. to here.} Setting $\lambda = \lambda_n = (n\pi)^2$ in the second line of $(*)$ and solving yields $T_n(t) = e^{-(n\pi)^2 t}$ (up to a constant factor). ^{16 pts. to here.} Thus

$$u(x,t) = a_0 + \sum_{n=1}^{\infty} [a_n \cos(n\pi x) + b_n \sin(n\pi x)] e^{-(n\pi)^2 t}$$

^{18 pts. to here.} is a formal solution to $\textcircled{1}-\textcircled{2}-\textcircled{3}$ for any choice of constants $a_0, a_1, b_1, a_2, b_2, \dots$. We want to choose these constants so $\textcircled{4}$ is met. Using the full Fourier series for $f(x) = x^3 - x$ in $[-1,1]$ from problem 6 yields

$$a_0 + \sum_{n=1}^{\infty} [a_n \cos(n\pi x) + b_n \sin(n\pi x)] = u(x,0) = x^3 - x = \sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi x)}{(n\pi)^3} \text{ for } -1 \leq x \leq 1$$

^{23 pts. to here.} Therefore choose $b_n = \frac{12(-1)^n}{(\pi n)^3}$ for $n=1,2,3,\dots$ and all other coefficients 0. Therefore a solution

of $\textcircled{1}-\textcircled{2}-\textcircled{3}-\textcircled{4}$ is

$$u(x,t) = \sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi x) e^{-(n\pi)^2 t}}{(\pi n)^3}$$

(OVER for bonus)

^{25 pts. to here.}

Bonus

To show that this solution is unique, suppose that $v = v(x, t)$ is any other solution of ①-②-③-④ and consider $w(x, t) = u(x, t) - v(x, t)$. Then w satisfies

1 pt. to here.

①' $w_t - w_{xx} = 0$ if $-1 < x < 1, 0 < t < \infty$,

3 pts. to here.

②'-③' $w(1, t) = w(-1, t)$ and $w_x(1, t) = w_x(-1, t)$ if $t \geq 0$,

④' $w(x, 0) = 0$ if $-1 \leq x \leq 1$.

6 pts. to here.

Consider the energy $E(t) = \int_{-1}^1 w^2(x, t) dx$ of this solution at t in $[0, \infty)$. Then

$$\frac{dE}{dt} = \int_{-1}^1 \frac{\partial}{\partial t} [w^2(x, t)] dx = \int_{-1}^1 2w(x, t) w_t(x, t) dx = \int_{-1}^1 \underbrace{2w(x, t)}_U \underbrace{w_{xx}(x, t)}_{dV} dx$$

$$= \underbrace{2w(1, t)w_x(1, t) - 2w(-1, t)w_x(-1, t)}_{0 \text{ by } \textcircled{2'} - \textcircled{3'}} - 2 \int_{-1}^1 w_x^2(x, t) dx$$

$$= -2 \int_{-1}^1 w_x^2(x, t) dx$$

$$\leq 0.$$

8 pts. to here.

Therefore $E = E(t)$ is a decreasing function on $[0, \infty)$, so for $t \geq 0$,

$$0 \leq E(t) \leq E(0) = \int_{-1}^1 w^2(0, x) dx = \int_{-1}^1 0 dx = 0 \text{ (by } \textcircled{4'}) .$$

By the vanishing theorem, the nonnegative continuous integrand of

$E(t)$ must be identically zero on $-1 \leq x \leq 1$ for each $t \geq 0$. That is,

$$0 = w(x, t) = u(x, t) - v(x, t) \text{ for } -1 \leq x \leq 1 \text{ and } 0 \leq t < \infty.$$

10 pts. to here.

Hence the solution to ①-②-③-④ is unique.

8. (25 pts.) (a) Solve $\nabla^2 u = 0$ inside the unit disk, subject to the boundary condition

$$u(1; \theta) = \begin{cases} 1 & \text{if } 0 < \theta < \pi, \\ -1 & \text{if } -\pi < \theta < 0, \\ 0 & \text{if } \theta = 0 \text{ or } \pi. \end{cases}$$

(b) What is the value of the solution to part (a) at the center of the disk? Support your answer.

4 pts.
to here.

(a) Poisson's formula, $u(r; \theta) = \frac{a^2 - r^2}{2\pi} \int_{-\pi}^{\pi} \frac{h(\varphi) d\varphi}{a^2 - 2ar \cos(\theta - \varphi) + r^2}$, applied

to this case yields

10 pts.
to here.

$$u(r; \theta) = \frac{1-r^2}{2\pi} \left[\int_{-\pi}^0 \frac{-1 d\varphi}{1-2r \cos(\varphi - \theta) + r^2} + \int_0^{\pi} \frac{1 d\varphi}{1-2r \cos(\varphi - \theta) + r^2} \right].$$

12 pts.
to here.

Making use of the integration fact $\int \frac{d\psi}{A+B \cos(\psi)} = \frac{2}{\sqrt{A^2-B^2}} \text{Arctan} \left(\tan\left(\frac{\psi}{2}\right) \sqrt{\frac{A-1}{A+1}} \right)$

if $A > |B|$ (OVER for details), we have

16 pts.
to here.

$$u(r; \theta) = \frac{1-r^2}{2\pi} \left[\frac{2}{\sqrt{(1+r^2)^2 - 4r^2}} \text{Arctan} \left(\tan\left(\frac{\varphi - \theta}{2}\right) \sqrt{\frac{1+r^2+2r}{1+r^2-2r}} \right) \right]_0^{\pi} - \frac{2}{\sqrt{(1+r^2)^2 - 4r^2}} \text{Arctan} \left(\tan\left(\frac{\varphi - \theta}{2}\right) \sqrt{\frac{1+r^2+2r}{1+r^2-2r}} \right) \Big|_{-\pi}^0 \right]$$

$$= \frac{1}{\pi} \left[\text{Arctan} \left(\tan\left(\frac{\pi - \theta}{2}\right) \left(\frac{1+r}{1-r}\right) \right) - \text{Arctan} \left(\tan\left(\frac{\theta}{2}\right) \left(\frac{1+r}{1-r}\right) \right) - \text{Arctan} \left(\tan\left(\frac{\theta}{2}\right) \left(\frac{1+r}{1-r}\right) \right) + \text{Arctan} \left(\tan\left(\frac{\pi - \theta}{2}\right) \left(\frac{1+r}{1-r}\right) \right) \right]$$

$$= \boxed{\frac{2}{\pi} \text{Arctan} \left(\cot\left(\frac{\theta}{2}\right) \left(\frac{1+r}{1-r}\right) \right)}.$$

20 pts.
to here.

(b) By the mean value property for harmonic functions,

5 pts.

$$u(0; \theta) = \frac{1}{2\pi} \int_{-\pi}^{\pi} h(\varphi) d\varphi = \frac{1}{2\pi} \left[\int_{-\pi}^0 -1 d\varphi + \int_0^{\pi} 1 d\varphi \right] = \boxed{0}.$$

A Brief Table of Fourier Transforms

$f(x)$	$\hat{f}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\xi x} dx$
A. $\begin{cases} 1 & \text{if } -b < x < b, \\ 0 & \text{otherwise.} \end{cases}$	$\sqrt{\frac{2}{\pi}} \frac{\sin(b\xi)}{\xi}$
B. $\begin{cases} 1 & \text{if } c < x < d, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{e^{-ic\xi} - e^{-id\xi}}{i\xi\sqrt{2\pi}}$
C. $\frac{1}{x^2 + a^2} \quad (a > 0)$	$\sqrt{\frac{\pi}{2}} \frac{e^{-a \xi }}{a}$
D. $\begin{cases} x & \text{if } 0 < x \leq b, \\ 2b - x & \text{if } b < x < 2b, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{-1 + 2e^{-ib\xi} - e^{-2ib\xi}}{\xi^2\sqrt{2\pi}}$
E. $\begin{cases} e^{-ax} & \text{if } x > 0, \\ 0 & \text{otherwise.} \end{cases}$ (a > 0)	$\frac{1}{(a + i\xi)\sqrt{2\pi}}$
F. $\begin{cases} e^{ax} & \text{if } b < x < c, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{e^{(a-i\xi)c} - e^{(a-i\xi)b}}{(a - i\xi)\sqrt{2\pi}}$
G. $\begin{cases} e^{iax} & \text{if } -b < x < b, \\ 0 & \text{otherwise.} \end{cases}$	$\sqrt{\frac{2}{\pi}} \frac{\sin(b(\xi-a))}{\xi - a}$
H. $\begin{cases} e^{iax} & \text{if } c < x < d, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{i}{\sqrt{2\pi}} \frac{e^{ic(a-\xi)} - e^{id(a-\xi)}}{a - \xi}$
I. $e^{-ax^2} \quad (a > 0)$	$\frac{1}{\sqrt{2a}} e^{-\xi^2/(4a)}$
J. $\frac{\sin(ax)}{x} \quad (a > 0)$	$\begin{cases} 0 & \text{if } \xi \geq a, \\ \sqrt{\pi/2} & \text{if } \xi < a. \end{cases}$

Convergence Theorems

$$X'' + \lambda X = 0 \text{ in } (a, b) \text{ with any symmetric BC.} \quad (1)$$

Now let $f(x)$ be any function defined on $a \leq x \leq b$. Consider the Fourier series for the problem (1) with any given boundary conditions that are symmetric. We now state a convergence theorem for each of the three modes of convergence. They are partly proved in the next section.

Theorem 2. Uniform Convergence The Fourier series $\sum A_n X_n(x)$ converges to $f(x)$ uniformly on $[a, b]$ provided that

- (i) $f(x)$, $f'(x)$, and $f''(x)$ exist and are continuous for $a \leq x \leq b$ and
- (ii) $f(x)$ satisfies the given boundary conditions.

Theorem 3. L^2 Convergence The Fourier series converges to $f(x)$ in the mean-square sense in (a, b) provided only that $f(x)$ is any function for which

$$\int_a^b |f(x)|^2 dx \text{ is finite.} \quad (8)$$

Theorem 4. Pointwise Convergence of Classical Fourier Series

- (i) The classical Fourier series (full or sine or cosine) converges to $f(x)$ pointwise on (a, b) , provided that $f(x)$ is a continuous function on $a \leq x \leq b$ and $f'(x)$ is piecewise continuous on $a \leq x \leq b$.
- (ii) More generally, if $f(x)$ itself is only piecewise continuous on $a \leq x \leq b$ and $f'(x)$ is also piecewise continuous on $a \leq x \leq b$, then the classical Fourier series converges at every point x ($-\infty < x < \infty$). The sum is

$$\sum_n A_n X_n(x) = \frac{1}{2} [f(x+) + f(x-)] \quad \text{for all } a < x < b. \quad (9)$$

The sum is $\frac{1}{2} [f_{\text{ext}}(x+) + f_{\text{ext}}(x-)]$ for all $-\infty < x < \infty$, where $f_{\text{ext}}(x)$ is the extended function (periodic, odd periodic, or even periodic).

Theorem 4 $^\infty$. If $f(x)$ is a function of period $2l$ on the line for which $f(x)$ and $f'(x)$ are piecewise continuous, then the classical full Fourier series converges to $\frac{1}{2} [f(x+) + f(x-)]$ for $-\infty < x < \infty$.

Math 325
Final Exam
Summer 2008

n : 20

μ : 119.8

σ : 51.5

Distribution of Scores:

174 - 200	2
146 - 173	6
120 - 145	3
100 - 119	3
0 - 99	6

Distribution of Letter Grades

A	2
B	7
C	5
D	6
F	0