

1. (40 pts.) (a) Show that the full Fourier series of $f(x) = x^3 - x$ on $[-1, 1]$ is $\sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi x)}{(\pi n)^3}$.
- 5 (b) On the same coordinate axes, sketch the graph of f and the sum of the first three nonzero terms of the full Fourier series of f on $[-1, 1]$.
- 5 (c) Discuss the L^2 -convergence, or lack thereof, for the full Fourier series of f on $[-1, 1]$.
- 5 (d) Discuss the pointwise convergence, or lack thereof, for the full Fourier series of f on $[-1, 1]$.
- 7 (e) Discuss the uniform convergence, or lack thereof, for the full Fourier series of f on $[-1, 1]$.
- 5 (f) Use the results above to help compute the sum of the series $\sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)^3}$.
- 5 (g) Use the results above to help compute the sum of the series $\sum_{n=1}^{\infty} \frac{1}{n^6}$.

(a) $f(x) \sim a_0 + \sum_{n=1}^{\infty} [a_n \cos(n\pi x) + b_n \sin(n\pi x)]$ where $a_0 = \frac{1}{2} \int_{-1}^1 f(x) dx = 0$, 1 pt. to here

$a_n = \int_{-1}^1 \underbrace{f(x)}_{\text{odd}} \underbrace{\cos(n\pi x)}_{\text{even}} dx = 0$, $b_n = \int_{-1}^1 \underbrace{f(x)}_{\text{odd}} \underbrace{\sin(n\pi x)}_{\text{odd}} dx = 2 \int_0^1 (x^3 - x) \sin(n\pi x) dx =$ 3 pts. to here.

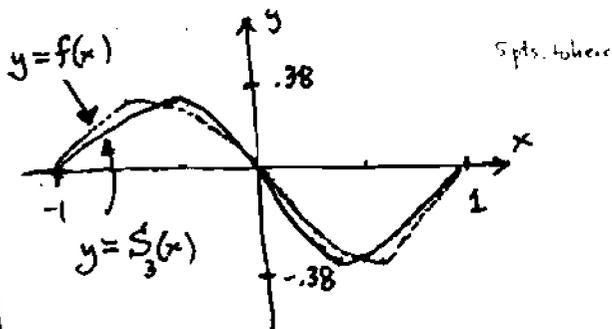
where $2 \int_0^1 (x^3 - x) \left(\frac{-\cos(n\pi x)}{n\pi} \right) \Big|_0^1 + \frac{2}{n\pi} \int_0^1 \underbrace{(3x^2 - 1)}_u \underbrace{\cos(n\pi x)}_{dv} dx = \frac{2}{n\pi} (3x^2 - 1) \frac{\sin(n\pi x)}{n\pi} \Big|_0^1 - \frac{12}{(n\pi)^2} \int_0^1 \underbrace{x}_{u} \underbrace{\sin(n\pi x)}_{dv} dx$ 5 pts. to here.

$= \frac{-12}{(n\pi)^2} x \left(\frac{-\cos(n\pi x)}{n\pi} \right) \Big|_0^1 + \frac{12}{(n\pi)^3} \int_0^1 \cos(n\pi x) dx = \frac{12 \cos(n\pi)}{(n\pi)^3} = \frac{12(-1)^n}{(n\pi)^3}$ 7 pts. to here

Thus $f(x) \sim \sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi x)}{(\pi n)^3}$

(b) $S_3(x) = \sum_{n=1}^3 \frac{12(-1)^n \sin(n\pi x)}{(\pi n)^3}$

$\Rightarrow S_3(x) = \frac{12}{\pi^3} \left(-\sin(\pi x) + \frac{1}{8} \sin(2\pi x) - \frac{1}{27} \sin(3\pi x) \right)$



(e) The full Fourier series of f is the Fourier series of f with respect to the orthogonal set of functions $\mathcal{E} = \{1, \cos(\pi x), \sin(\pi x), \cos(2\pi x), \sin(2\pi x), \dots\}$ on the interval $(-1, 1)$. These functions are the complete set of eigenfunctions for the problem

$\mathcal{E}''(x) + \lambda \mathcal{E}(x) = 0, \mathcal{E}(1) = \mathcal{E}(-1), \mathcal{E}'(1) = \mathcal{E}'(-1)$

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Note that $f(x) = x^3 - x$ satisfies the (periodic) boundary conditions:

3 pts. to here. $f(1) = 0 = f(-1)$ and $f'(1) = 2 = f'(-1)$. Furthermore $f(x) = x^3 - x$, $f'(x) = 3x^2 - 1$,
 5 pts. to here. and $f''(x) = 6x$ are continuous on the closed interval $-1 \leq x \leq 1$. Consequently,
 7 pts. to here. Theorem 2 guarantees that the full Fourier series of f converges uniformly to f
 on $[-1, 1]$.

(c)-(d). Because uniform convergence ^{to f} on $[-1, 1]$ implies both L^2 -convergence
 and pointwise convergence ^{to f} on $[-1, 1]$, it follows that the full Fourier series
 of f converges to f , both in the mean-square sense and the pointwise sense,
 on $[-1, 1]$.

(f) $\left(\frac{1}{2}\right)^3 - \frac{1}{2} = f\left(\frac{1}{2}\right) = \sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi/2)}{(n\pi)^3}$ ^{by part (d)} 3 pts. to here.

$$\Rightarrow -\frac{3}{8} = \frac{12}{\pi^3} \sum_{k=0}^{\infty} \frac{(-1)^{2k+1} (-1)^k}{(2k+1)^3}$$

n	$\sin(n\pi/2)$
1	1
2	0
3	-1
4	0
\vdots	\vdots
$2k$	0
$2k+1$	$(-1)^k$

$$\Rightarrow \left(-\frac{\pi^3}{12}\right) \left(-\frac{3}{8}\right) = \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)^3}$$

$$\boxed{\frac{\pi^3}{32} = \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)^3}}$$

5 pts. to here.

(g) By Parseval's identity, $\sum_{n=1}^{\infty} |c_n|^2 \int_a^b |\sum_{n=1}^{\infty} c_n \phi_n(x)|^2 dx = \int_a^b |f(x)|^2 dx$ 2 pts. to here.

when $\{\phi_1, \phi_2, \phi_3, \dots\}$ is a complete set of orthogonal functions on $[a, b]$.

Applying this to our case, $\sum_{n=1}^{\infty} |b_n|^2 \int_{-1}^1 \sin^2(n\pi x) dx = \int_{-1}^1 |x^3 - x|^2 dx$

(since $a_0 = a_1 = \dots = 0$). Thus $\sum_{n=1}^{\infty} \left| \frac{12(-1)^n}{(n\pi)^3} \right|^2 \cdot 1 = 2 \int_0^1 (x-x^3)^2 dx = \frac{16}{105}$

so $\sum_{n=1}^{\infty} \frac{1}{n^6} = \frac{16}{105} \cdot \frac{\pi^6}{3} = \frac{\pi^6}{945}$ 5 pts. to here.

Please recall that the Laplacian of u in polar coordinates is $\nabla^2 u = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2}$ while in

spherical coordinates it is $\nabla^2 u = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial u}{\partial r} \right) + \frac{1}{r^2 \sin(\phi)} \frac{\partial}{\partial \phi} \left(\sin(\phi) \frac{\partial u}{\partial \phi} \right) + \frac{1}{r^2 \sin^2(\phi)} \frac{\partial^2 u}{\partial \theta^2}$.

2.(30 pts.) Let $0 < a < b < \infty$. Solve $\nabla^2 u = 1$ in $\mathcal{R} : a^2 < x^2 + y^2 + z^2 < b^2$ subject to $u = 0$ on $\partial \mathcal{R}$.

3 pts. to here. Assume $u = u(r)$, independent of the spherical angles ϕ and θ . (This is reasonable to assume because the PDE, region, and boundary conditions are invariant with respect to rotations.) Then $\frac{\partial u}{\partial \phi} = \frac{\partial^2 u}{\partial \theta^2} = 0$ so $\nabla^2 u = 1$ becomes $\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial u}{\partial r} \right) = 1$ 9 pts. to here

$$\Rightarrow \frac{\partial}{\partial r} \left(r^2 \frac{\partial u}{\partial r} \right) = r^2 \Rightarrow r^2 \frac{\partial u}{\partial r} = \frac{r^3}{3} + c_1 \Rightarrow \frac{\partial u}{\partial r} = \frac{r}{3} + \frac{c_1}{r^2} \Rightarrow$$

15 pts. to here

21 pts. to here. $u = \frac{r^2}{6} - \frac{c_1}{r} + c_2$. Applying the boundary conditions $u|_{r=a} = 0 = u|_{r=b}$ we have

$$\left. \begin{aligned} 0 &= \frac{a^2}{6} - \frac{c_1}{a} + c_2 \\ 0 &= \frac{b^2}{6} - \frac{c_1}{b} + c_2 \end{aligned} \right\} \text{ subtract: } 0 = \frac{b^2 - a^2}{6} - c_1 \left(\frac{1}{b} - \frac{1}{a} \right) \Rightarrow -c_1 \left(\frac{b-a}{ab} \right) = \frac{b^2 - a^2}{6}$$

$$\therefore c_1 = -\frac{(b+a)ab}{6}$$

27 pts. to here.

$$\therefore c_2 = \frac{c_1}{a} - \frac{a^2}{6} = -\frac{(b+a)b}{6} - \frac{a^2}{6} \quad 29 \text{ pts. to here}$$

Thus $u(r, \theta, \phi) = \frac{r^2}{6} + \frac{(b+a)ab}{6r} - \frac{(b+a)b}{6} - \frac{a^2}{6}$ 30 pts. to here

This simplifies to $u(r, \theta, \phi) = \frac{(r-a)(r-b)(r+a+b)}{6r}$.

Note: The solution above is unique. This is a direct consequence of the maximum/minimum principle for harmonic functions.

3. (30 pts.) Solve $\nabla^2 u = 0$ in the cube $0 < x < 1$, $0 < y < 1$, $0 < z < 1$ given that u satisfies the inhomogeneous Dirichlet condition $u(x, 0, z) = 4 \sin^2(\pi x) \sin^2(\pi z)$ if $0 \leq x \leq 1$, $0 \leq z \leq 1$, and u satisfies homogeneous Neumann boundary conditions on the other five faces. Hint: You may find the identity $2 \sin^2(\theta) = 1 - \cos(2\theta)$ useful.

2 pts. to here.

$u(x, y, z) = X(x)Y(y)Z(z)$ (nontrivial solution to the homogeneous portion of this problem) leads

to $X''(x)Y(y)Z(z) + X(x)Y''(y)Z(z) + X(x)Y(y)Z''(z) = 0$ (when substituted in the PDE)

$\Rightarrow -\frac{X'(x)}{X(x)} = \frac{Y'(y)}{Y(y)} + \frac{Z'(z)}{Z(z)} = \lambda$ and $-\frac{Z'(z)}{Z(z)} = \frac{Y'(y)}{Y(y)} - \lambda = \mu$. Thus

(*) $\left\{ \begin{array}{l} X''(x) + \lambda X(x) = 0, \quad X'(0) = 0 = X'(1). \\ Z''(z) + \mu Z(z) = 0, \quad Z'(0) = 0 = Z'(1). \\ Y''(y) - (\lambda + \mu) Y(y) = 0, \quad Y'(1) = 0. \end{array} \right.$ 5 pts. to here
8 pts. to here
10 pts. to here.

[The Neumann boundary conditions $u_x(0, y, z) = u_x(1, y, z) = 0$ for all $0 \leq y \leq 1, 0 \leq z \leq 1$, $u_z(x, y, 0) = u_z(x, y, 1) = 0$ for all $0 \leq x \leq 1, 0 \leq y \leq 1$, and $u_y(x, 1, z) = 0$ for all $0 \leq x \leq 1, 0 \leq z \leq 1$ yield the boundary conditions $X'(0) = 0 = X'(1)$, $Z'(0) = Z'(1) = 0$, and $Y'(1) = 0$ above.]

The eigenvalue problems in the first two lines of (*) have solutions:

$\lambda_\ell = (\ell\pi)^2$, $X_\ell(x) = \cos(\ell\pi x)$ ($\ell = 0, 1, 2, \dots$), 14 pts. to here.

$\mu_m = (m\pi)^2$, $Z_m(z) = \cos(m\pi z)$ ($m = 0, 1, 2, \dots$). 18 pts. to here

Substituting $\lambda + \mu = \lambda_\ell + \mu_m = \pi^2(\ell^2 + m^2)$ in the third line of (*) and solving yields $Y_{\ell, m}(y) = \cosh(\pi(y-1)\sqrt{\ell^2 + m^2})$ (up to a constant factor). 20 pts. to here

Thus, the superposition principle implies that

$u(x, y, z) = \sum_{\ell=0}^{\infty} \sum_{m=0}^{\infty} A_{\ell, m} \cos(\ell\pi x) \cos(m\pi z) \cosh(\pi(y-1)\sqrt{\ell^2 + m^2})$ 22 pts. to here

is a formal solution to the homogeneous portion of this problem for any

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choice of constants $A_{l,m}$ ($l, m = 0, 1, 2, \dots$). Therefore, to satisfy the inhomogeneous Dirichlet boundary condition, we require that

$$\begin{aligned} \sum_{l=0}^{\infty} \sum_{m=0}^{\infty} A_{l,m} \cosh(-\pi\sqrt{l^2+m^2}) \cos(l\pi x) \cos(m\pi z) &= u(x, 0, z) = 4 \sin^2(\pi x) \sin^2(\pi z) \\ &= [1 - \cos(2\pi x)][1 - \cos(2\pi z)] \\ &= 1 - \cos(2\pi x) - \cos(2\pi z) + \cos(2\pi x) \cos(2\pi z) \end{aligned}$$

for all $0 \leq x \leq 1$ and $0 \leq z \leq 1$. Consequently, equating like coefficients produces

$$\left. \begin{aligned} A_{0,0} &= 1 \\ A_{2,0} \cosh(-2\pi) &= -1 \\ A_{0,2} \cosh(-2\pi) &= -1 \\ A_{2,2} \cosh(-2\pi\sqrt{2}) &= 1 \end{aligned} \right\} \Rightarrow \begin{aligned} A_{0,0} &= 1 \\ A_{2,0} &= \frac{-1}{\cosh(2\pi)} = A_{0,2} \\ A_{2,2} &= \frac{1}{\cosh(2\pi\sqrt{2})} \end{aligned}$$

2.6 pts. to here.

and all other $A_{l,m} = 0$. That is,

$$\begin{aligned} u(x, y, z) &= 1 - \frac{\cos(2\pi x) \cosh(2\pi(y-1))}{\cosh(2\pi)} - \frac{\cos(2\pi z) \cosh(2\pi(y-1))}{\cosh(2\pi)} \\ &\quad + \frac{\cos(2\pi x) \cos(2\pi z) \cosh(2\pi(y-1)\sqrt{2})}{\cosh(2\pi\sqrt{2})} \end{aligned}$$

3.0 pts. to here.

solves the problem.

Bonus (30 pts.). Solve $u_t - u_{xx} \stackrel{\textcircled{1}}{=} 0$ in $-1 < x < 1$, $0 < t < \infty$, subject to the boundary conditions $u(1,t) \stackrel{\textcircled{2}}{=} u(-1,t)$ and $u_x(1,t) \stackrel{\textcircled{3}}{=} u_x(-1,t)$ if $t \geq 0$, and the initial condition $u(x,0) \stackrel{\textcircled{4}}{=} x^3 - x$ if $-1 \leq x \leq 1$. Furthermore, show that this solution is unique. Hint: You may find the results of problem 1 useful.

2 pts. to here. We seek nontrivial solutions of $\textcircled{1}-\textcircled{2}-\textcircled{3}$ of the form $u(x,t) = X(x)T(t)$. Substituting into $\textcircled{1}$ yields $T'(t)X(x) - T(t)X''(x) = 0 \Rightarrow \frac{T'(t)}{T(t)} = \frac{X''(x)}{X(x)} = -\lambda$. Substituting in $\textcircled{2}-\textcircled{3}$ yields $X(1)T(t) = X(-1)T(t)$ and $X'(1)T(t) = X'(-1)T(t)$ for all $t \geq 0$. Using the fact that $u(x,t) = X(x)T(t)$ is not identically zero yields

6 pts. to here.

$$(*) \begin{cases} X''(x) + \lambda X(x) = 0, & X(1) = X(-1), \quad X'(1) = X'(-1), \\ T'(t) + \lambda T(t) = 0. \end{cases}$$

7 pts. to here.

pts. to here.

11 pts. to here.

13 pts. to here.

15 pts. to here.

The eigenvalue problem in the first line of $(*)$ has solutions $\lambda = \lambda_n = (n\pi)^2$ and $X_n(x) = a_n \cos(n\pi x) + b_n \sin(n\pi x)$ ($n=0,1,2,3,\dots$ and a_n, b_n arbitrary constants).

Setting $\lambda = \lambda_n = (n\pi)^2$ in the second line of $(*)$ yields $T_n(t) = e^{-(n\pi)^2 t}$ (up to a constant factor). Thus $u(x,t) = a_0 + \sum_{n=1}^{\infty} [a_n \cos(n\pi x) + b_n \sin(n\pi x)] e^{-(n\pi)^2 t}$

is a formal solution to $\textcircled{1}-\textcircled{2}-\textcircled{3}$ for any choice of constants $a_0, a_1, b_1, a_2, b_2, \dots$. We want to choose these constants so the initial condition $\textcircled{4}$ is met. Using the full Fourier series for $f(x) = x^3 - x$ in $[-1,1]$ in problem 1 yields

17 pts. to here.

$$a_0 + \sum_{n=1}^{\infty} [a_n \cos(n\pi x) + b_n \sin(n\pi x)] = u(x,0) = x^3 - x = \sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi x)}{(\pi n)^3}$$

19 pts. to here.

for all $-1 \leq x \leq 1$. Therefore choose $a_0 = a_1 = a_2 = \dots = 0$ and $b_n = \frac{12(-1)^n}{(\pi n)^3}$ for $n=1,2,3,\dots$

That is,

21 pts. to here.

$$u(x,t) = \sum_{n=1}^{\infty} \frac{12(-1)^n \sin(n\pi x) e^{-(n\pi)^2 t}}{(\pi n)^3}$$

To show that this solution is unique, suppose that $v = v(x,t)$ is any other

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22 pts. to here. solution of ①-②-③-④ and consider $w(x,t) = u(x,t) - v(x,t)$. Then w satisfies

①' $w_t - w_{xx} = 0$ if $-1 < x < 1, 0 < t < \infty$,

24 pts. to here.

②'-③' $w(1,t) = w(-1,t)$ and $w_x(1,t) = w_x(-1,t)$ if $t \geq 0$,

④' $w(x,0) = 0$ if $-1 \leq x \leq 1$.

26 pts. to here

Consider the energy $E(t) = \int_{-1}^1 w^2(x,t) dx$ of this solution at $t \geq 0$.

Show $\frac{dE}{dt} = \int_{-1}^1 \frac{\partial}{\partial t} [w^2(x,t)] dx = \int_{-1}^1 2w(x,t) w_t(x,t) dx = \int_{-1}^1 \underbrace{2w(x,t)}_U \underbrace{w_{xx}(x,t)}_{dV} dx$

$= 2w(x,t) w_x(x,t) \Big|_{x=-1}^1 - \int_{-1}^1 2w_x^2(x,t) dx$

$= \underbrace{2w(1,t)w_x(1,t) - 2w(-1,t)w_x(-1,t)}_{0 \text{ by } \textcircled{2}'-\textcircled{3}'} - 2 \int_{-1}^1 w_x^2(x,t) dx$

28 pts. to here.

$= -2 \int_{-1}^1 w_x^2(x,t) dx \leq 0$. (I.e. $E = E(t)$ is a decreasing function.)

Therefore $0 \leq E(t) \leq E(0) = \int_{-1}^1 w^2(x,0) dx \stackrel{\text{by } \textcircled{4}'}{=} \int_{-1}^1 0 dx = 0$ for $t \geq 0$. By the

vanishing theorem, the nonnegative continuous integrand of $E(t)$ must be identically zero for all $-1 \leq x \leq 1$ and each $t \geq 0$. That is,

30 pts. to here.

$u(x,t) - v(x,t) = w(x,t) = 0$ for all $-1 \leq x \leq 1, 0 \leq t < \infty$,

and hence the solution to ①-②-③-④ is unique.

Convergence Theorems

$$X'' + \lambda X = 0 \text{ in } (a, b) \text{ with any symmetric BC.} \quad (1)$$

Now let $f(x)$ be any function defined on $a \leq x \leq b$. Consider the Fourier series for the problem (1) with any given boundary conditions that are symmetric. We now state a convergence theorem for each of the three modes of convergence. They are partly proved in the next section.

Theorem 2. Uniform Convergence The Fourier series $\sum A_n X_n(x)$ converges to $f(x)$ uniformly on $[a, b]$ provided that

- (i) $f(x)$, $f'(x)$, and $f''(x)$ exist and are continuous for $a \leq x \leq b$ and
- (ii) $f(x)$ satisfies the given boundary conditions.

Theorem 3. L^2 Convergence The Fourier series converges to $f(x)$ in the mean-square sense in (a, b) provided only that $f(x)$ is any function for which

$$\int_a^b |f(x)|^2 dx \text{ is finite.} \quad (8)$$

Theorem 4. Pointwise Convergence of Classical Fourier Series

- (i) The classical Fourier series (full or sine or cosine) converges to $f(x)$ pointwise on (a, b) , provided that $f(x)$ is a continuous function on $a \leq x \leq b$ and $f'(x)$ is piecewise continuous on $a \leq x \leq b$.
- (ii) More generally, if $f(x)$ itself is only piecewise continuous on $a \leq x \leq b$ and $f'(x)$ is also piecewise continuous on $a \leq x \leq b$, then the classical Fourier series converges at every point x ($-\infty < x < \infty$). The sum is

$$\sum_n A_n X_n(x) = \frac{1}{2} [f(x+) + f(x-)] \quad \text{for all } a < x < b. \quad (9)$$

The sum is $\frac{1}{2} [f_{\text{ext}}(x+) + f_{\text{ext}}(x-)]$ for all $-\infty < x < \infty$, where $f_{\text{ext}}(x)$ is the extended function (periodic, odd periodic, or even periodic).

Theorem 4^∞ . If $f(x)$ is a function of period $2l$ on the line for which $f(x)$ and $f'(x)$ are piecewise continuous, then the classical full Fourier series converges to $\frac{1}{2} [f(x+) + f(x-)]$ for $-\infty < x < \infty$.