

1.(33 pts.) Use Fourier transform methods to solve

$$u_t - u_{xx} + tu = 0 \quad \text{for } -\infty < x < \infty, 0 < t < \infty,$$

$$u(x, 0) = f(x) \quad \text{for } -\infty < x < \infty.$$

$$\mathcal{F}(u_t - u_{xx} + tu)(\xi) = \mathcal{F}(0)(\xi)$$

$$\frac{\partial}{\partial t} \mathcal{F}(u)(\xi) + \xi^2 \mathcal{F}(u)(\xi) + t \mathcal{F}(u)(\xi) = 0$$

Integrating factor: $e^{\int (\xi^2 + t) dt} = e^{\xi^2 t + \frac{t^2}{2}}$

$$e^{\xi^2 t + \frac{t^2}{2}} \frac{\partial}{\partial t} \mathcal{F}(u)(\xi) + (\xi^2 + t) e^{\xi^2 t + \frac{t^2}{2}} \mathcal{F}(u)(\xi) = 0$$

$$\frac{\partial}{\partial t} \left(e^{\xi^2 t + \frac{t^2}{2}} \mathcal{F}(u)(\xi) \right) = 0$$

$$e^{\xi^2 t + \frac{t^2}{2}} \mathcal{F}(u)(\xi) = A(\xi)$$

$$\mathcal{F}(u)(\xi) = A(\xi) e^{-\xi^2 t - \frac{t^2}{2}}$$

$$\hat{f}(\xi) = \mathcal{F}(u(\cdot, 0))(\xi) = \mathcal{F}(u)(\xi) \Big|_{t=0} = A(\xi)$$

$$\therefore \mathcal{F}(u)(\xi) = \hat{f}(\xi) e^{-\xi^2 t - \frac{t^2}{2}}$$

Take $a = \frac{1}{4t}$ in formula I in the table of Fourier transform to get

$$\mathcal{F}\left(\frac{e^{-\frac{(\cdot)^2}{4t}}}{\sqrt{2t}}\right)(\xi) = e^{-\xi^2 t}$$

$$\therefore \mathcal{F}(u)(\xi) = \mathcal{F}(f)(\xi) \cdot \mathcal{F}\left(\frac{e^{-\frac{(\cdot)^2}{4t}}}{\sqrt{2t}}\right)(\xi) \cdot e^{-\frac{t^2}{2}}$$

Using the property that

$$\mathcal{F}(g * h)(\xi) = \sqrt{2\pi} \mathcal{F}(g)(\xi) \mathcal{F}(h)(\xi)$$

gives

$$\mathcal{F}(u)(\xi) = \frac{1}{\sqrt{2\pi}} \mathcal{F}\left(f * \frac{e^{-\frac{(\cdot)^2}{4t}}}{\sqrt{2t}}\right)(\xi) \cdot e^{-\frac{t^2}{2}}$$

$$= \mathcal{F}\left(f * \frac{e^{-\frac{(\cdot)^2}{4t} - \frac{t^2}{2}}}{\sqrt{4\pi t}}\right)(\xi)$$

The inversion theorem then implies

$$u(x, t) = \left(f * \frac{e^{-\frac{(\cdot)^2}{4t} - \frac{t^2}{2}}}{\sqrt{4\pi t}}\right)(x)$$

$$= \int_{-\infty}^{\infty} \frac{e^{-\frac{(x-y)^2}{4t} - \frac{t^2}{2}}}{\sqrt{4\pi t}} f(y) dy$$

$$= \frac{e^{-\frac{t^2}{2}}}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4t}} f(y) dy$$

2.(33 pts.) (a) State (without providing a proof) the weak maximum/minimum principle for solutions to the diffusion equation.

(b) Show that solutions to the wave equation need not satisfy a maximum principle.

(c) Use the weak maximum/minimum principle to show that there is at most one solution to the Poisson equation with Dirichlet boundary conditions:

$$\begin{aligned}u_t - ku_{xx} &= f(x,t) \quad \text{for } 0 < x < L, 0 < t \leq T, \\u(x,0) &= \phi(x) \quad \text{for } 0 \leq x \leq L, \\u(0,t) &= g(t) \quad \text{and } u(L,t) = h(t) \quad \text{for } 0 \leq t \leq T.\end{aligned}$$

(a) Let $u = u(x,t)$ be a solution of $u_t - ku_{xx} = 0$ in $R = \{(x,t) : 0 < x < L, 0 < t \leq T\}$ and let u be continuous on $\bar{R} = \{(x,t) : 0 \leq x \leq L, 0 \leq t \leq T\}$. Then the maximum and minimum values of u on \bar{R} are attained on $\bar{R} \setminus R$, i.e. either initially ($t=0$) or on the lateral walls ($x=0$ or $x=L$).

(b) $u(x,t) = \sin(\pi t) \sin(\pi x)$ solves the wave equation $u_{tt} - u_{xx} = 0$ in $R = \{(x,t) : 0 < x < 1, 0 < t < 1\}$ and is continuous on $\bar{R} = \{(x,t) : 0 \leq x \leq 1, 0 \leq t \leq 1\}$. Yet $u = 0$ on $\bar{R} \setminus R$ while the maximum value of u on \bar{R} is $u(\frac{1}{2}, \frac{1}{2}) = 1$.

(c) Suppose $u = u_1(x,t)$ and $u = u_2(x,t)$ solve the problem in part (c) and are continuous functions on $\bar{R} = \{(x,t) : 0 \leq x \leq L, 0 \leq t \leq T\}$. Then $w(x,t) = u_1(x,t) - u_2(x,t)$ solves $u_t - ku_{xx} = 0$ in $R = \{(x,t) : 0 < x < L, 0 < t \leq T\}$ and is continuous on \bar{R} . Therefore the maximum and minimum of w on \bar{R} are attained on $\bar{R} \setminus R$. However $w(x,0) = 0$ for $0 \leq x \leq L$ and $w(0,t) = 0 = w(L,t)$ for $0 \leq t \leq T$. Therefore $w = 0$ on $\bar{R} \setminus R$, and it follows that $w(x,t) = 0$ for all (x,t) in \bar{R} ; i.e. $u_1(x,t) = u_2(x,t)$ for all (x,t) in \bar{R} .

3.(33 pts.) Use energy methods to show that there is at most one solution to the Poisson equation with Neumann boundary conditions:

$$\begin{aligned} u_t - ku_{xx} &= f(x,t) \quad \text{for } 0 < x < L, 0 < t < \infty, \\ u(x,0) &= \phi(x) \quad \text{for } 0 \leq x \leq L, \\ u_x(0,t) &= g(t) \quad \text{and } u_x(L,t) = h(t) \quad \text{for } 0 \leq t < \infty. \end{aligned}$$

Suppose $u = u_1(x,t)$ and $u = u_2(x,t)$ solve the above problem. Then $w(x,t) = u_1(x,t) - u_2(x,t)$ solves

① $u_t - ku_{xx} = 0$ for $0 < x < L, 0 < t < \infty,$

② $u(x,0) = 0$ for $0 \leq x \leq L,$

③-④ $u_x(0,t) = 0 = u_x(L,t)$ for $0 \leq t < \infty.$

Consider the energy function of w : $E(t) = \int_0^L w^2(x,t) dx$. Then $\frac{dE}{dt} =$

$$\int_0^L \frac{\partial}{\partial t} (w^2(x,t)) dx = \int_0^L 2w(x,t)w_t(x,t) dx \stackrel{\textcircled{1}}{=} 2k \int_0^L \overbrace{w(x,t)}^u \overbrace{w_{xx}(x,t)}^{dv} dx =$$

$$2k \cdot w(x,t)w_x(x,t) \Big|_{x=0}^L - 2k \int_0^L w_x^2(x,t) dx \stackrel{\textcircled{3-4}}{=} 0 - 2k \int_0^L w_x^2(x,t) dx \leq 0.$$

Therefore E is a nonincreasing function of t on $[0, \infty)$. Therefore, for $0 < t < \infty,$

$0 \leq E(t) \leq E(0) = \int_0^L w^2(x,0) dx \stackrel{\textcircled{2}}{=} \int_0^L 0 dx = 0.$ It follows that

$\int_0^L w^2(x,t) dx = E(t) = 0$ for all $0 \leq t < \infty.$ Consequently, the vanishing

theorem implies $w^2(x,t) = 0$ for all $-\infty < x < \infty$ for each $t \geq 0,$ and

hence $w(x,t) = 0$ for all $-\infty < x < \infty$ and $0 \leq t < \infty.$ That is, $u_1(x,t) = u_2(x,t)$

in the upper halfplane.

Bonus.(33 pts.) (a) Use Fourier transform methods to show that, under appropriate hypotheses, a solution to the Poisson initial value problem:

$$u_t - ku_{xx} = f(x,t) \quad \text{for } -\infty < x < \infty, 0 < t < \infty,$$

$$u(x,0) = \phi(x) \quad \text{for } -\infty < x < \infty,$$

is given by

$$u(x,t) = \int_{-\infty}^{\infty} \frac{e^{-\frac{(x-y)^2}{4kt}}}{\sqrt{4k\pi t}} \phi(y) dy + \int_0^t \int_{-\infty}^{\infty} \frac{e^{-\frac{(x-y)^2}{4k(t-\tau)}}}{\sqrt{4k\pi(t-\tau)}} f(y,\tau) dy d\tau.$$

(b) Use the formula in part (a) to solve the Poisson initial value problem with $\phi(x) \equiv 0$ and

$$f(x,t) = \begin{cases} x & \text{if } |x| < t, \\ 0 & \text{otherwise.} \end{cases}$$

(a) $\mathcal{F}(u_t - ku_{xx})(z) = \mathcal{F}(f(x,t))(z) \equiv F(z,t)$

$\frac{\partial}{\partial t} \mathcal{F}(u)(z) + k z^2 \mathcal{F}(u)(z) = F(z,t)$. An integrating factor is $e^{\int k z^2 dt} = e^{k z^2 t}$

$e^{k z^2 t} \frac{\partial}{\partial t} \mathcal{F}(u)(z) + k z^2 e^{k z^2 t} \mathcal{F}(u)(z) = e^{k z^2 t} F(z,t) \rightarrow e^{k z^2 t} \mathcal{F}(u)(z) = A(z) + \int_0^t e^{k z^2 \tau} F(z,\tau) d\tau$

$\frac{\partial}{\partial t} (e^{k z^2 t} \mathcal{F}(u)(z)) = e^{k z^2 t} F(z,t)$

$\therefore \mathcal{F}(u)(z) = A(z) e^{-k z^2 t} + e^{-k z^2 t} \int_0^t e^{k z^2 \tau} F(z,\tau) d\tau$

$= A(z) e^{-k z^2 t} + \int_0^t e^{-k z^2 (t-\tau)} F(z,\tau) d\tau$

$\mathcal{F}(\phi)(z) = \mathcal{F}(u(\cdot, 0))(z) = \mathcal{F}(u)(z) \Big|_{t=0} = A(z)$

$\therefore \mathcal{F}(u)(z) = \mathcal{F}(\phi)(z) e^{-k z^2 t} + \int_0^t e^{-k z^2 (t-\tau)} F(z,\tau) d\tau$

From formula I in the table of Fourier transforms, $\mathcal{F}\left(\frac{e^{-\frac{(x)^2}{4kt}}}{\sqrt{2kt}}\right)(z) = e^{-k z^2 t}$

(OVER)

$$\therefore \mathcal{F}_t(u)(z) = \mathcal{F}_t(\varphi)(z) \mathcal{F}_t\left(\frac{e^{-\frac{(\cdot)^2}{4kt}}}{\sqrt{2kt}}\right)(z) + \int_0^t \mathcal{F}_t\left(\frac{e^{-\frac{(\cdot)^2}{4k(t-\tau)}}}{\sqrt{2k(t-\tau)}}\right)(z) \mathcal{F}_t(f(\cdot, \tau)) d\tau$$

Using the convolution formula: $\mathcal{F}_t(g * h)(z) = \sqrt{2\pi} \mathcal{F}_t(g)(z) \mathcal{F}_t(h)(z)$ yields

$$\mathcal{F}_t(u)(z) = \frac{1}{\sqrt{2\pi}} \mathcal{F}_t\left(\varphi * \frac{e^{-\frac{(\cdot)^2}{4kt}}}{\sqrt{2kt}}\right)(z) + \frac{1}{\sqrt{2\pi}} \int_0^t \mathcal{F}_t\left(f(\cdot, \tau) * \frac{e^{-\frac{(\cdot)^2}{4k(t-\tau)}}}{\sqrt{2k(t-\tau)}}\right)(z) d\tau$$

$$= \mathcal{F}_t\left(\varphi * \frac{e^{-\frac{(\cdot)^2}{4kt}}}{\sqrt{4k\pi t}} + \int_0^t f(\cdot, \tau) * \frac{e^{-\frac{(\cdot)^2}{4k\pi(t-\tau)}}}{\sqrt{4k\pi(t-\tau)}} d\tau\right)(z)$$

By the inversion theorem,

$$u(x, t) = \left(\varphi * \frac{e^{-\frac{(\cdot)^2}{4kt}}}{\sqrt{4k\pi t}}\right)(x) + \int_0^t \left(f(\cdot, \tau) * \frac{e^{-\frac{(\cdot)^2}{4k\pi(t-\tau)}}}{\sqrt{4k\pi(t-\tau)}}\right)(x) d\tau$$

$$= \int_{-\infty}^{\infty} \frac{e^{-\frac{(x-y)^2}{4kt}}}{\sqrt{4k\pi t}} \varphi(y) dy + \int_0^t \int_{-\infty}^{\infty} \frac{e^{-\frac{(x-y)^2}{4k\pi(t-\tau)}}}{\sqrt{4k\pi(t-\tau)}} f(y, \tau) dy d\tau$$

$$(b) \quad u(x, t) = \int_0^t \int_{-\tau}^{\tau} \frac{e^{-\frac{(x-y)^2}{4k(t-\tau)}}}{\sqrt{4k\pi(t-\tau)}} y dy d\tau$$

$$\text{let } p = \frac{y-x}{\sqrt{4k(t-\tau)}}$$

$$\text{then } dp = \frac{dy}{\sqrt{4k(t-\tau)}}$$

$$= \int_0^t \frac{1}{\sqrt{\pi}} \int_{\frac{-\tau-x}{\sqrt{4k(t-\tau)}}}^{\frac{\tau-x}{\sqrt{4k(t-\tau)}}} e^{-p^2} (\sqrt{4k(t-\tau)} p + x) dp d\tau$$

(Cont.)

$$u(x,t) = \int_0^t \frac{\sqrt{4k(t-\tau)}}{(-2\sqrt{\pi})} \int_{\frac{-\tau-x}{\sqrt{4k(t-\tau)}}}^{\frac{\tau-x}{\sqrt{4k(t-\tau)}}} e^{-p^2} dp d\tau + \int_0^t \frac{x}{\sqrt{\pi}} \int_{\frac{-\tau-x}{\sqrt{4k(t-\tau)}}}^{\frac{\tau-x}{\sqrt{4k(t-\tau)}}} e^{-p^2} dp d\tau$$

$$= \int_0^t \frac{\sqrt{4k(t-\tau)}}{2\sqrt{\pi}} \left[-e^{-\frac{(\tau-x)^2}{4k(t-\tau)}} + e^{-\frac{(\tau+x)^2}{4k(t-\tau)}} \right] d\tau + \frac{x}{2} \int_0^t \left[\operatorname{Erf} \left(\frac{\tau-x}{\sqrt{4k(t-\tau)}} \right) + \operatorname{Erf} \left(\frac{\tau+x}{\sqrt{4k(t-\tau)}} \right) \right] d\tau$$

A Brief Table of Fourier Transforms

	f(x)	$\hat{f}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\xi x} dx$
A.	$\begin{cases} 1 & \text{if } -b < x < b, \\ 0 & \text{otherwise.} \end{cases}$	$\sqrt{\frac{2}{\pi}} \frac{\sin(b\xi)}{\xi}$
B.	$\begin{cases} 1 & \text{if } c < x < d, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{e^{-ic\xi} - e^{-id\xi}}{i\xi\sqrt{2\pi}}$
C.	$\frac{1}{x^2 + a^2} \quad (a > 0)$	$\sqrt{\frac{\pi}{2}} \frac{e^{-a \xi }}{a}$
D.	$\begin{cases} x & \text{if } 0 < x \leq b, \\ 2b - x & \text{if } b < x < 2b, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{-1 + 2e^{-ib\xi} - e^{-2ib\xi}}{\xi^2\sqrt{2\pi}}$
E.	$\begin{cases} e^{-ax} & \text{if } x > 0, \\ 0 & \text{otherwise.} \end{cases}$ <p style="text-align: center;">(a > 0)</p>	$\frac{1}{(a + i\xi)\sqrt{2\pi}}$
F.	$\begin{cases} e^{ax} & \text{if } b < x < c, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{e^{(a-i\xi)c} - e^{(a-i\xi)b}}{(a - i\xi)\sqrt{2\pi}}$
G.	$\begin{cases} e^{iax} & \text{if } -b < x < b, \\ 0 & \text{otherwise.} \end{cases}$	$\sqrt{\frac{2}{\pi}} \frac{\sin(b(\xi-a))}{\xi - a}$
H.	$\begin{cases} e^{iax} & \text{if } c < x < d, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{i}{\sqrt{2\pi}} \frac{e^{ic(a-\xi)} - e^{id(a-\xi)}}{a - \xi}$
I.	$e^{-ax^2} \quad (a > 0)$	$\frac{1}{\sqrt{2a}} e^{-\xi^2/(4a)}$
J.	$\frac{\sin(ax)}{x} \quad (a > 0)$	$\begin{cases} 0 & \text{if } \xi \geq a, \\ \sqrt{\pi/2} & \text{if } \xi < a. \end{cases}$