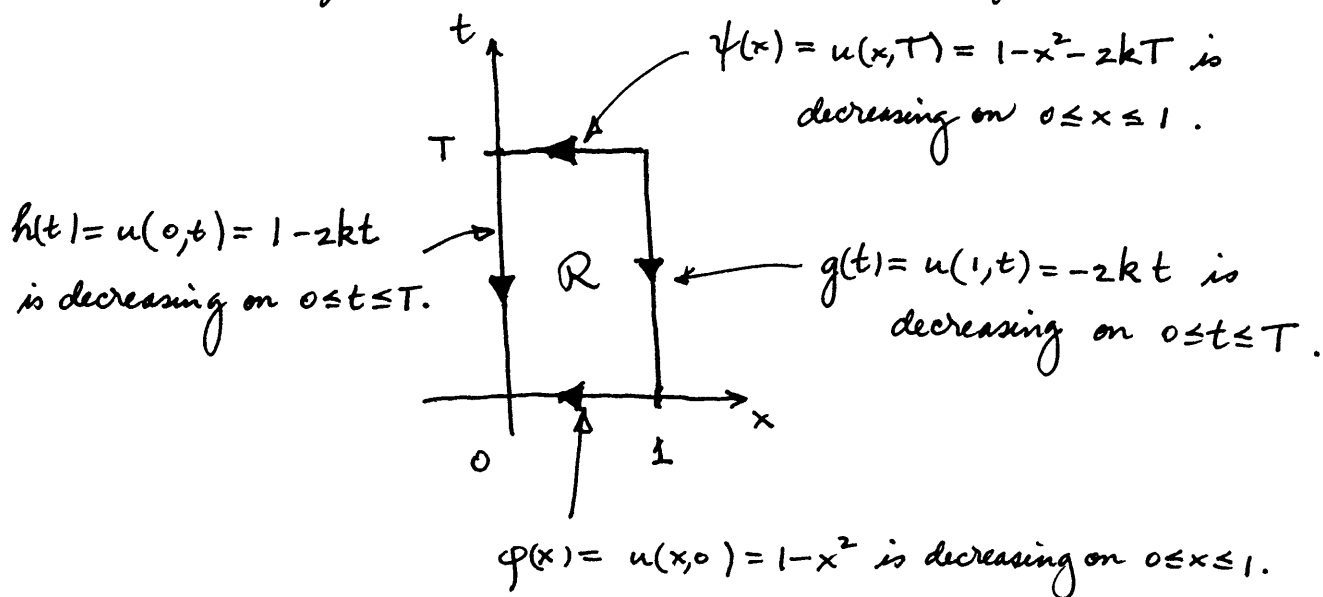


Sec. 2.3, pp. 44-45.

#1. Consider the solution $u(x,t) = 1 - x^2 - 2kt$ of the diffusion equation $u_t - ku_{xx} = 0$. Find the locations of its maximum and minimum in the closed rectangle $\bar{R} = \{(x,t) : 0 \leq x \leq 1, 0 \leq t \leq T\}$.

Since $u_t = -2k \neq 0$, the maximum and minimum of u on \bar{R} cannot occur at an interior point of \bar{R} . The values of u on $\partial\bar{R}$ behave as follows (arrows show the directions of increasing u).



Therefore $\max_{(x,t) \in \bar{R}} u(x,t) = u(0,0) = 1$

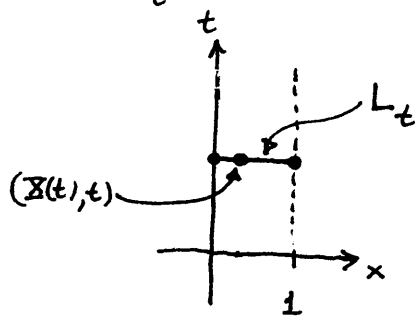
and $\min_{(x,t) \in \bar{R}} u(x,t) = u(1,T) = -2kT$.

(Note that the maximum and minimum values of u occur on the initial and side walls of \bar{R} , in accordance with the maximum/minimum principle.)

Sec. 2.3, pp. 44-45.

3 (cont.) points (x, t) such that $0 < x < 1$ and $0 < t < \infty$.

(b) Let $\bar{x}(t)$ denote the number in $[0, 1]$ such that $\mu(t) = u(\bar{x}(t), t)$, i.e. the maximum of $u(x, t)$ on the horizontal line segment $L_t = \{(x, t) : 0 \leq x \leq 1\}$ occurs at the point $(\bar{x}(t), t)$.



(If the maximum of $u(x, t)$ on L_t occurs at more than one point, then take $(\bar{x}(t), t)$ as the left-most point on L_t where the maximum of $u(x, t)$ occurs.)

By the chain rule applied to $\mu(t) = u(\bar{x}(t), t)$, we have

$$\begin{aligned} \frac{d\mu}{dt} &= \frac{\partial u}{\partial \bar{x}} \cdot \frac{d\bar{x}}{dt} + \frac{\partial u}{\partial t} \cdot \frac{dt}{dt} \\ &= \frac{\partial u}{\partial \bar{x}} \cdot \frac{d\bar{x}}{dt} + \frac{\partial^2 u}{\partial x^2} \end{aligned}$$

since u is a solution to $u_t - u_{xx} = 0$. But $\frac{\partial u}{\partial x}(\bar{x}(t), t) = 0$ and $\frac{\partial^2 u}{\partial x^2}(\bar{x}(t), t) \leq 0$ since the maximum of u on L_t occurs at $(\bar{x}(t), t)$. Consequently

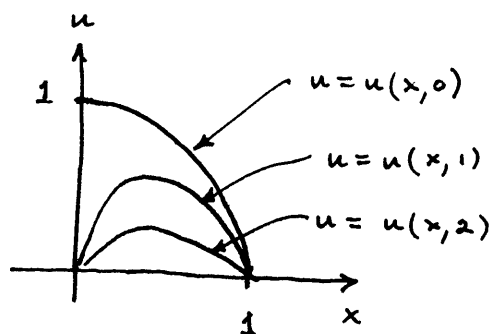
$$\frac{d\mu}{dt} = 0 \cdot \frac{d\bar{x}}{dt} + \frac{\partial^2 u}{\partial x^2} \leq 0;$$

i.e. μ is a nonincreasing function of t .

Sec 2.3, pp. 44-45.

3 (cont.)

(c)



#5. The purpose of this exercise is to show that the conclusion of the maximum principle need not hold for the equation $u_t = x u_{xx}$, which has a variable coefficient.

(a) Verify that $u(x, t) = -2xt - x^2$ is a solution. Find the location of its maximum in the rectangle $\bar{R} = \{(x, t) : -2 \leq x \leq 2, 0 \leq t \leq 1\}$.

(b) Where precisely does our proof of the maximum principle break down for this equation?

(a) If $u(x, t) = -2xt - x^2$ then

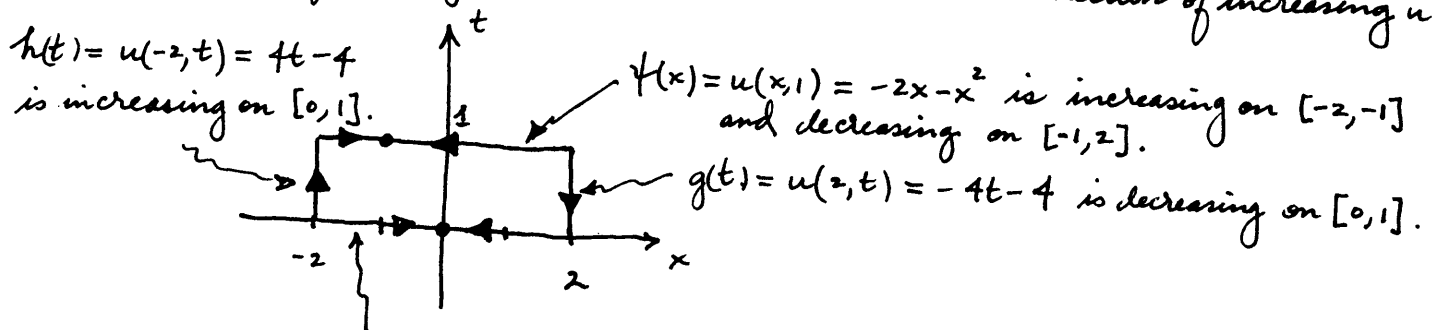
$$u_t - x u_{xx} = -2x - x(-2) \stackrel{v}{=} 0.$$

Note that $u_x = -2t - 2x$ and $u_t = -2x$. Therefore if the maximum of u occurs at an interior point of \bar{R} , then $u_t = u_x = 0$ there. But $u_x = u_t = 0$ implies $x = t = 0$, and $(0, 0)$ is not an interior point of \bar{R} . Therefore the maximum of u on \bar{R} must occur at a boundary point.

On the boundary of \bar{R} , the function $u(x, t)$ has the

Sec. 2.3, pp. 44-45.

5 (cont.) following behaviors (arrows indicate the direction of increasing u).



$\varphi(x) = u(x, 0) = -x^2$ is increasing on $[-2, 0]$ and decreasing on $[0, 2]$.

$$u(0, 0) = 0 - 0 = 0, \quad u(-1, 1) = 2 - 1 = 1$$

Thus the maximum of u on the rectangle occurs at $(-1, 1)$, a point on the back wall $t = 1$. Thus the maximum of u on \bar{R} does not occur at a point on the initial or side walls of \bar{R} , in opposition to the conclusion of the maximum principle.

(b) The analogue of the diffusion inequality (2), p. 42, does not hold for this equation on the rectangle R :

$$v_t - x v_{xx} = (u + \epsilon x^2)_t - x (u + \epsilon x^2)_{xx} = u_t - x (u_{xx} + 2\epsilon)$$

$$= \underbrace{u_t - x u_{xx}}_0 - 2x\epsilon$$

$$= \underbrace{-2x\epsilon}$$

This need not be negative for all $-2 \leq x \leq 2$.

Sec. 2.3, pp. 44-45.

#8. Consider the diffusion equation $u_t - ku_{xx} = 0$ on $0 < x < l$ with the Robin boundary conditions $u_x(0,t) - a_0 u(0,t) = 0$ and $u_x(l,t) + a_l u(l,t) = 0$. If $a_0 > 0$ and $a_l > 0$, use the energy method to show that the endpoints contribute to the decrease of $\int_0^l u^2(x,t) dx$. (This is interpreted to mean that part of the "energy" is lost at the boundary, so we call the boundary conditions "radiating" or "dissipative".)

Let $E(t) = \int_0^l u^2(x,t) dx$. Then

$$\frac{dE}{dt} = \frac{d}{dt} \int_0^l u^2(x,t) dx$$

$$= \int_0^l \frac{\partial}{\partial t} \{u^2(x,t)\} dx$$

$$= \int_0^l 2u(x,t) u_t(x,t) dx$$

$$= \int_0^l \overbrace{2u(x,t)}^v \overbrace{ku_{xx}(x,t)}^{dv} dx$$

(Because $u_t - ku_{xx} = 0$.)

$$= \overbrace{2ku(x,t)}^v \overbrace{u_x(x,t)}^v \Big|_{x=0}^{x=l} - \int_0^l \overbrace{2ku_x(x,t)}^{v'} \overbrace{u_x(x,t)}^v dx$$

$$= 2k(u(l,t)u_x(l,t) - u(0,t)u_x(0,t)) - 2k \int_0^l u_x^2(x,t) dx$$

$$= \underbrace{2k(-a_l u^2(l,t) - a_0 u^2(0,t))}_{\text{the contribution of the endpoints to the loss in } E(t)} - 2k \int_0^l u_x^2(x,t) dx$$

the contribution of the endpoints to the loss in $E(t)$.