

Fourier Transforms and Diffusions in the Upper Halfplane

Def: Let f be an absolutely integrable function on $(-\infty, \infty)$ (i.e. let $f \in L^1(\mathbb{R})$).

The Fourier transform of f is the function given by

$$\mathcal{F}(f)(\xi) = \hat{f}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\xi x} dx \quad (-\infty < \xi < \infty).$$

Ex 1 Compute the Fourier transform of the function

$$\chi_{(c,d)}(x) = \begin{cases} 1 & \text{if } x \in (c,d), \\ 0 & \text{o.w.} \end{cases}$$

Solution:

$$\begin{aligned} \hat{\chi}_{(c,d)}(\xi) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \chi_{(c,d)}(x) e^{-i\xi x} dx \\ &= \frac{1}{\sqrt{2\pi}} \int_c^d e^{-i\xi x} dx \\ &= \frac{1}{\sqrt{2\pi}} \left(\frac{e^{-i\xi x}}{-i\xi} \Big|_{x=c}^{x=d} \right) \quad (\text{provided } \xi \neq 0) \end{aligned}$$

Formula B.

$$\hat{\chi}_{(c,d)}(\xi) = \frac{e^{-i\xi c} - e^{-i\xi d}}{i\xi \sqrt{2\pi}} \quad (\xi \neq 0)$$

Notes: ① $\hat{\chi}_{(c,d)}(0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \chi_{(c,d)}(x) e^{-i0x} dx = \frac{d-c}{\sqrt{2\pi}} \stackrel{\text{l'H\^opital}}{=} \lim_{\xi \rightarrow 0} \hat{\chi}_{(c,d)}(\xi).$

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② The special case $c = -b$ and $d = b$ in Ex 1 yields

Formula A.

$$\hat{\chi}_{(-b,b)}(\xi) = \frac{e^{i\xi b} - e^{-i\xi b}}{i\xi \sqrt{2\pi}} = \frac{2i \sin(b\xi)}{i\xi \sqrt{2\pi}} = \sqrt{\frac{2}{\pi}} \frac{\sin(b\xi)}{\xi}$$

Unnormalised
sinc function

$$\xi \rightarrow \frac{\pi \xi}{b} \Rightarrow \frac{\sin(\xi)}{\xi}$$

(used in
digital signal processing
and information theory)

$\mathcal{F}(f)(\xi) = \hat{f}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\xi x} dx$

Properties of the Fourier Transform:

- (Linearity) ① $\widehat{(c_1 f_1 + c_2 f_2)}(\xi) = c_1 \hat{f}_1(\xi) + c_2 \hat{f}_2(\xi)$ (Valid if f_1, f_2 are absolutely integrable and $c_1, c_2 \in \mathbb{R}$.)
- (Transform of a derivative) ② $\widehat{f'}(\xi) = i\xi \hat{f}(\xi)$ (Valid if f is differentiable and f and f' are ~~absolutely integrable~~ absolutely integrable.)
- (Derivative of a transform) ③ $\frac{d}{d\xi} \mathcal{F}(f)(\xi) = -i \mathcal{F}(x f(x))(\xi)$ (Valid if f and $x \mapsto x f(x)$ are ~~absolutely integrable~~ absolutely integrable.)
- (Convolutions) ④ $\widehat{f * g}(\xi) = \sqrt{2\pi} \hat{f}(\xi) \hat{g}(\xi)$ (Valid if f, g are absolutely integrable)
- (Inversion Theorem) ⑤ $f(x) = \lim_{M \rightarrow \infty} \frac{1}{\sqrt{2\pi}} \int_{-M}^M \hat{f}(\xi) e^{i\xi x} d\xi$ (Valid if f is continuous, piecewise smooth, and absolutely integrable.)

See also the shifting theorems (exercises #4, 6 for Fourier transforms).

Property ①: $\widehat{(c_1 f_1 + c_2 f_2)}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} [c_1 f_1(x) + c_2 f_2(x)] e^{-i\xi x} dx$

$$= c_1 \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f_1(x) e^{-i\xi x} dx + c_2 \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f_2(x) e^{-i\xi x} dx$$

$$= c_1 \hat{f}_1(\xi) + c_2 \hat{f}_2(\xi).$$

Property ②: $\widehat{f'}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f'(x) e^{-i\xi x} dx$

$$= \lim_{\substack{M \rightarrow \infty \\ N \rightarrow -\infty}} \frac{1}{\sqrt{2\pi}} \int_N^M f'(x) e^{-i\xi x} dx$$

Parts: $v = e^{-i\xi x}$ $dv = f'(x) dx$ $dU = -i\xi e^{-i\xi x} dx$ $U = f(x)$

(+)

$$= \lim_{\substack{M \rightarrow \infty \\ N \rightarrow -\infty}} \left\{ \frac{f(M) e^{-i\xi M}}{\sqrt{2\pi}} - \frac{f(N) e^{-i\xi N}}{\sqrt{2\pi}} + \frac{1}{\sqrt{2\pi}} \int_{-N}^M i\xi f(x) e^{-i\xi x} dx \right\}$$

Since f' is absolutely integrable on $(-\infty, \infty)$, to each $\epsilon > 0$ there corresponds

$\therefore M_0 = M_0(\epsilon) > 0$ such that

$$\left| f(M_2) - f(M_1) \right| = \left| \int_{M_1}^{M_2} f'(x) dx \right| \leq \int_{M_1}^{M_2} |f'(x)| dx < \epsilon$$

for all $M_1, M_2 \geq M_0$. It follows that $\lim_{M \rightarrow \infty} f(M)$ exists.

A similar argument shows that $\lim_{N \rightarrow +\infty} f(-N)$ exists. However,

f is absolutely integrable on $(-\infty, \infty)$ so $\lim_{M \rightarrow \infty} f(M) = 0 = \lim_{N \rightarrow +\infty} f(-N)$.

Returning to (†), we see that

$$\lim_{M \rightarrow \infty} \frac{\overbrace{f(M)}^{\text{bounded}} e^{-i\xi M}}{\sqrt{2\pi}} = 0 \quad \text{and} \quad \lim_{N \rightarrow +\infty} \frac{\overbrace{f(-N)}^{\text{bounded}} e^{+i\xi N}}{\sqrt{2\pi}} = 0.$$

Consequently, $\widehat{f}'(\xi) = i\xi \widehat{f}(\xi)$.

Property ③: $\frac{d}{d\xi} \widehat{f}(f)(\xi) = \frac{d}{d\xi} \left(\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\xi x} dx \right)$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial}{\partial \xi} \left(f(x) e^{-i\xi x} \right) dx$$

(Theorem 2, p. 390, Appendix A.3)

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} -ix f(x) e^{-i\xi x} dx$$

$$= -i \widehat{f}(xf(x))(\xi).$$

Ex 2 Compute the Fourier transform of the function $f(x) = e^{-ax^2}$ where a is a positive constant.

Solution: $f(x) = e^{-ax^2}$
 $f'(x) = -2ax e^{-ax^2}$

$\therefore f'(x) + 2axf(x) = 0 \quad (-\infty < x < \infty)$

We take the Fourier transform of both sides of this identity and use linearity:

$$\mathcal{F}(f')(\xi) + 2a\mathcal{F}(xf(x))(\xi) = 0$$

Applying properties (2) and (3) gives

$$i\xi \hat{f}(\xi) + 2ai \frac{d}{d\xi} \hat{f}(\xi) = 0$$

$$\Rightarrow \int \frac{d\hat{f}(\xi)}{\hat{f}(\xi)} = \int \frac{-\xi}{2a} d\xi$$

$$\ln \hat{f}(\xi) = -\frac{\xi^2}{4a} + C$$

$$\hat{f}(\xi) = A e^{-\frac{\xi^2}{4a}} \quad (A = e^C)$$

To find A, we evaluate at $\xi = 0$:

$$A = \hat{f}(0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i0x} dx = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-ax^2} dx \quad \left\{ \begin{array}{l} \text{Let } p = x\sqrt{a} \\ dp = \sqrt{a} dx \end{array} \right.$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-p^2} \frac{dp}{\sqrt{a}} = \frac{1}{\sqrt{2\pi a}} \sqrt{\pi} = \frac{1}{\sqrt{2a}}$$

$$\therefore \mathcal{F}(e^{-ax^2})(\xi) = \frac{e^{-\frac{\xi^2}{4a}}}{\sqrt{2a}} \quad \text{Formula I.}$$

Note: $\mathcal{F}(e^{-\frac{x^2}{2}})(\xi) = e^{-\frac{\xi^2}{2}}$. Therefore $f(x) = e^{-\frac{x^2}{2}}$ is a fixed point of the Fourier transform operator; i.e. $\mathcal{F}f = f$.
 Take $a = \frac{1}{2}$ in Formula I to get //

Convolutions

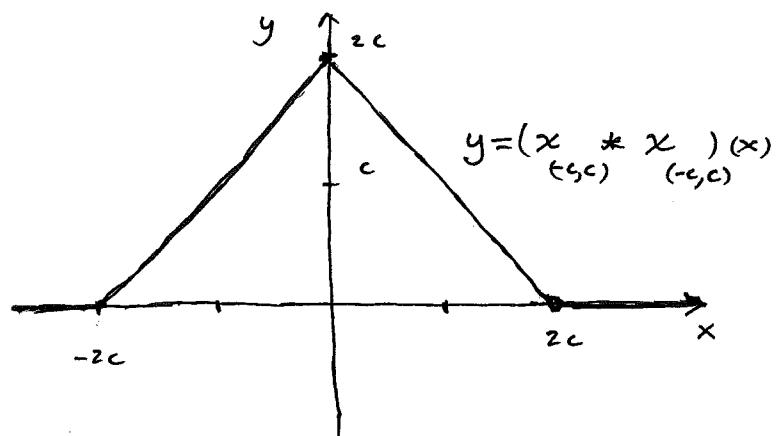
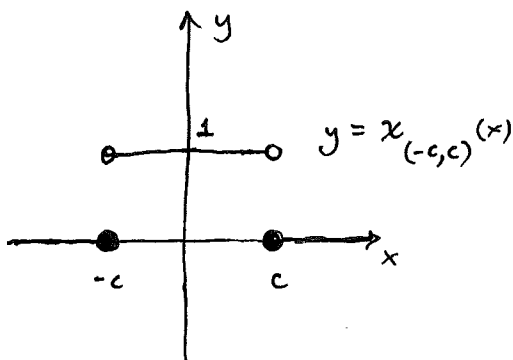
Def: Let f and g be absolutely integrable on $(-\infty, \infty)$. Their convolution product (denoted $f * g$) is the function given by

$$(f * g)(x) = \int_{-\infty}^{\infty} f(x-y)g(y) dy \quad (-\infty < x < \infty).$$

Ex 3: Compute the convolution product of $\chi_{(-c,c)}$ with itself.

$$\begin{aligned} \text{Soln: } (\chi_{(-c,c)} * \chi_{(-c,c)})(x) &= \int_{-\infty}^{\infty} \chi_{(-c,c)}(x-y) \chi_{(-c,c)}(y) dy \\ &= \int_{-\infty}^{\infty} \chi_{(x-c, x+c)}(y) \chi_{(-c,c)}(y) dy \\ &= \int_{-\infty}^{\infty} \chi_{(x-c, x+c) \cap (-c,c)}(y) dy \\ &= \text{length of } \{ (x-c, x+c) \cap (-c,c) \} \\ &= \begin{cases} 2c - |x| & \text{if } |x| \leq 2c, \\ 0 & \text{o.w.} \end{cases} \end{aligned}$$

$$\begin{aligned} -c < x-y < c \\ \updownarrow \\ -c < y-x < c \\ \updownarrow \\ x-c < y < x+c \end{aligned}$$



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Note: Ex 3 illustrates the basic principle that convolution is a

"smoothing" operation. ^{Moreover} $f * g$ inherits the smoothness properties of either factor.

E.g. if f is differentiable (with f, f', g absolutely integrable) then $(f * g)'(x) = (f' * g)(x)$.

Elementary Properties of the Convolution Product

1. $f * g = g * f$

2. $(c_1 f_1 + c_2 f_2) * g = c_1 (f_1 * g) + c_2 (f_2 * g)$

3. If f and g are absolutely integrable then so is $f * g$:

$$\int_{-\infty}^{\infty} |(f * g)(x)| dx \leq \int_{-\infty}^{\infty} |f(x)| dx \cdot \int_{-\infty}^{\infty} |g(y)| dy$$

Property (4) of Fourier Transforms: $\widehat{f * g}(\xi) = \sqrt{2\pi} \hat{f}(\xi) \hat{g}(\xi)$

(4): $\widehat{f * g}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} (f * g)(x) e^{-i\xi x} dx$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} f(x-y) g(y) dy \right) e^{-i\xi x} dx$$

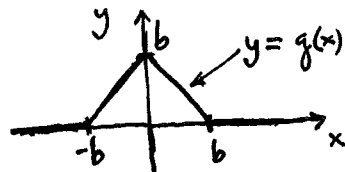
$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} f(x-y) e^{-i\xi(x-y)} dx \right) g(y) e^{-i\xi y} dy$$

Let $z = x - y$.
Then $dz = dx$.

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} f(z) e^{-i\xi z} dz \right) g(y) e^{-i\xi y} dy$$

$$= \frac{1}{\sqrt{2\pi}} \left(\int_{-\infty}^{\infty} f(z) e^{-i\xi z} dz \right) \int_{-\infty}^{\infty} g(y) e^{-i\xi y} dy$$

$$= \sqrt{2\pi} \hat{f}(\xi) \hat{g}(\xi).$$



Ex 4] Compute the Fourier transform of $g(x) = \begin{cases} b - |x| & \text{if } |x| \leq b, \\ 0 & \text{if } |x| > b. \end{cases}$

Solution: Recall from Ex 3 with $c = \frac{b}{2}$ that $g(x) = \chi_{(-\frac{b}{2}, \frac{b}{2})} * \chi_{(-\frac{b}{2}, \frac{b}{2})}(x)$.

Verification of Property 3 of Convolutions:

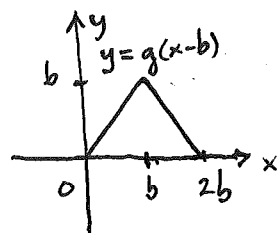
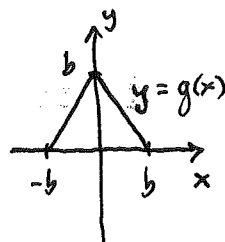
$$\begin{aligned}\int_{-\infty}^{\infty} |(f * g)(x)| dx &= \int_{-\infty}^{\infty} \left| \int_{-\infty}^{\infty} f(x-y)g(y) dy \right| dx \\ &\leq \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} |f(x-y)g(y)| dy dx \\ &= \int_{-\infty}^{\infty} |g(y)| \left(\int_{-\infty}^{\infty} |f(x-y)| dx \right) dy \\ &= \int_{-\infty}^{\infty} |g(y)| \left[\int_{-\infty}^{\infty} |f(z)| dz \right] dy \\ &= \left(\int_{-\infty}^{\infty} |f(z)| dz \right) \int_{-\infty}^{\infty} |g(y)| dy\end{aligned}$$

$$\text{Therefore } \hat{g}(\xi) = \widehat{x_{(-\frac{b}{2}, \frac{b}{2})} * x_{(-\frac{b}{2}, \frac{b}{2})}}(\xi)$$

$$= \sqrt{2\pi} \hat{x}_{(-\frac{b}{2}, \frac{b}{2})}(\xi) \hat{x}_{(-\frac{b}{2}, \frac{b}{2})}(\xi)$$

$$= \sqrt{2\pi} \left(\sqrt{\frac{2}{\pi}} \frac{\sin(b\xi/2)}{\xi} \right) \left(\sqrt{\frac{2}{\pi}} \frac{\sin(b\xi/2)}{\xi} \right) \quad (\text{by Ex 1})$$

$$= \boxed{2\sqrt{\frac{2}{\pi}} \frac{\sin^2(b\xi/2)}{\xi^2}}$$



Notes: $g(x-b) = \begin{cases} x & \text{if } 0 < x \leq b, \\ 2b-x & \text{if } b < x < 2b, \\ 0 & \text{o.w.,} \end{cases}$

so, by Ex 4 and shifting on the x-axis (exercise #4),

$$\mathcal{F}\{g(x-b)\}(\xi) = e^{-i\xi b} \mathcal{F}\{g\}(\xi)$$

$$= 2\sqrt{\frac{2}{\pi}} e^{-i\xi b} \left(\frac{e^{i\xi b/2} - e^{-i\xi b/2}}{2i} \right)^2 \cdot \frac{1}{\xi^2}$$

$$= 2\sqrt{\frac{2}{\pi}} e^{-i\xi b} \left(\frac{e^{i\xi b} - 2 + e^{-i\xi b}}{-4} \right) \cdot \frac{1}{\xi^2}$$

$$= \frac{-1 + 2e^{-i\xi b} - e^{-2i\xi b}}{\xi^2 \sqrt{2\pi}}$$

Formula D.

omit if
pressed
for time.

Property ⑤: For a proof of the inversion theorem see: R. Goldberg, Fourier Transforms, Cambridge U. Press, 1962, pp. 10-13.

or M. Pinsky, PDE's and BVP's with Applications (2nd ed.), McGraw-Hill, 1991, pp. 253-4.

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Solution of the Diffusion Equation in the Upper Half plane (via Fourier transforms)

Consider the I.V.P.

$$(*) \quad \begin{cases} u_t - ku_{xx} = f(x,t) & \text{for } -\infty < x < \infty, 0 < t < \infty, \\ u(x,0) = \varphi(x) & \text{for } -\infty < x < \infty. \end{cases}$$

We will use Fourier transform techniques to derive the solution

$$(**) \quad u(x,t) = \int_{-\infty}^{\infty} \frac{e^{-\frac{(x-y)^2}{4kt}}}{\sqrt{4k\pi t}} \varphi(y) dy + \int_0^t \int_0^{\infty} \frac{1}{\sqrt{4k\pi(t-\tau)}} e^{-\frac{(x-y)^2}{4k(t-\tau)}} f(y,\tau) dy d\tau.$$

The derivation will be "formal", i.e. we won't worry about the validity of interchanging integration and differentiation, convergence of improper integrals, etc.

STEP 1. Convert the PDE to an ODE ^{problem} by taking the Fourier transform.

We begin by taking the Fourier transform of the p.d.e. in (*) with respect to the variable x .

$$\mathcal{F}(u_t(\cdot, t) - ku_{xx}(\cdot, t))(\xi) = \mathcal{F}(f(\cdot, t))(\xi) = \hat{f}(\xi, t)$$

Using properties ① and ② of the Fourier transform,

$$\mathcal{F}(u_t)(\xi) - k\mathcal{F}(u_{xx})(\xi) = \hat{f}(\xi, t)$$

$$\mathcal{F}(u_t)(\xi) - k(i\xi)^2 \mathcal{F}(u)(\xi) = \hat{f}(\xi, t)$$

$$\text{But } \mathcal{F}(u_t)(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial u(x,t)}{\partial t} e^{-i\xi x} dx = \frac{\partial}{\partial t} \left(\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u(x,t) e^{-i\xi x} dx \right) = \frac{\partial}{\partial t} \mathcal{F}(u)(\xi)$$

$$\text{so } (+) \quad \frac{\partial}{\partial t} \mathcal{F}(u)(\xi) + k\xi^2 \mathcal{F}(u)(\xi) = \hat{f}(\xi, t).$$

(Observe that via the Fourier transform, we have converted the p.d.e.

in (*) in the variables x and t into an o.d.e. in the variable t ,

STEP 2. Solve the ODE problem

with parameter ξ .) An integrating factor for the first-order linear equation (†) is

$$e^{\int k\xi^2 dt} = e^{k\xi^2 t + C}$$

Multiplying (†) by this integrating factor yields

$$e^{k\xi^2 t} \frac{\partial \mathcal{F}(u)(\xi)}{\partial t} + k\xi^2 e^{k\xi^2 t} \mathcal{F}(u)(\xi) = e^{k\xi^2 t} \hat{f}(\xi, t)$$

$$\frac{\partial}{\partial t} \left(e^{k\xi^2 t} \mathcal{F}(u)(\xi) \right) = e^{k\xi^2 t} \hat{f}(\xi, t)$$

Consequently,
$$e^{k\xi^2 t} \mathcal{F}(u)(\xi) = \int_0^t \hat{f}(\xi, \tau) e^{k\xi^2 \tau} d\tau + A(\xi) \quad \text{or}$$

$$\mathcal{F}(u)(\xi) = \int_0^t \hat{f}(\xi, \tau) e^{-k\xi^2(t-\tau)} d\tau + A(\xi) e^{-k\xi^2 t}$$

In order to determine the factor $A(\xi)$ we use the initial condition

in (*):

$$A(\xi) = \left[A(\xi) e^{-k\xi^2 t} + \int_0^t \hat{f}(\xi, \tau) e^{-k\xi^2(t-\tau)} d\tau \right] \Big|_{t=0} = \mathcal{F}(u(\cdot, t))(\xi) \Big|_{t=0} = \mathcal{F}(u(\cdot, 0))(\xi) = \mathcal{F}(\varphi)(\xi)$$

Therefore

$$(†) \quad \mathcal{F}(u)(\xi) = \mathcal{F}(\varphi)(\xi) e^{-k\xi^2 t} + \int_0^t \hat{f}(\xi, \tau) e^{-k\xi^2(t-\tau)} d\tau \equiv U_1(\xi, t) + U_2(\xi, t)$$

STEP 3.

Convert the solution of the ODE problem into a solution of the PDE.

(I.e. compute the inverse Fourier transform of the ODE's solution.)

Using formula I in the table of Fourier transforms,

$$\mathcal{F}(e^{-a(\cdot)^2})(\xi) = \frac{e^{-\xi^2/4a}}{\sqrt{2a}},$$

with $a = \frac{1}{4kt}$ yields

$$e^{-kt\xi^2} = \mathcal{F}\left(\sqrt{\frac{1}{2kt}} e^{-\frac{(\cdot)^2}{4kt}}\right)(\xi).$$

Thus $U_1(\xi, t) = \mathcal{F}(\varphi)(\xi) \mathcal{F}\left(\frac{1}{\sqrt{2kt}} e^{-\frac{(\cdot)^2}{4kt}}\right)(\xi)$. Using property (4) gives

$$\begin{aligned} \text{(III)} \quad U_1(\xi, t) &= \frac{1}{\sqrt{2\pi}} \mathcal{F}\left(\varphi * \frac{1}{\sqrt{2kt}} e^{-\frac{(\cdot)^2}{4kt}}\right)(\xi) \\ &= \mathcal{F}\left(\varphi * \frac{1}{\sqrt{4k\pi t}} e^{-\frac{(\cdot)^2}{4kt}}\right)(\xi). \end{aligned}$$

Using formula I in the table of Fourier transforms with $a = \frac{1}{4k(t-\tau)}$ gives

$$e^{-k\xi^2(t-\tau)} = \mathcal{F}\left(\sqrt{\frac{1}{2k(t-\tau)}} e^{-\frac{(\cdot)^2}{4k(t-\tau)}}\right)(\xi).$$

Therefore

$$\begin{aligned} U_2(\xi, t) &= \int_0^t \hat{f}(\xi, \tau) e^{-k\xi^2(t-\tau)} d\tau \\ &= \int_0^t \mathcal{F}(f(\cdot, \tau))(\xi) \mathcal{F}\left(\frac{1}{\sqrt{2k(t-\tau)}} e^{-\frac{(\cdot)^2}{4k(t-\tau)}}\right)(\xi) d\tau. \end{aligned}$$

Property (7) and interchanging the order of integration then yields

$$\begin{aligned} \text{(IV)} \quad U_2(\xi, t) &= \int_0^t \frac{1}{\sqrt{2\pi}} \mathcal{F}\left(f(\cdot, \tau) * \frac{1}{\sqrt{2k(t-\tau)}} e^{-\frac{(\cdot)^2}{4k(t-\tau)}}\right)(\xi) d\tau \\ &= \mathcal{F}\left(\int_0^t f(\cdot, \tau) * \frac{1}{\sqrt{4k\pi(t-\tau)}} e^{-\frac{(\cdot)^2}{4k(t-\tau)}} d\tau\right)(\xi). \end{aligned}$$

Consequently, substituting from (III) and (IV) into (II) leads to

$$\begin{aligned} \mathcal{F}(u(\cdot, t))(\xi) &= \mathcal{F}\left(\varphi * \frac{1}{\sqrt{4k\pi t}} e^{-\frac{(\cdot)^2}{4kt}}\right)(\xi) + \mathcal{F}\left(\int_0^t f(\cdot, \tau) * \frac{1}{\sqrt{4k\pi(t-\tau)}} e^{-\frac{(\cdot)^2}{4k(t-\tau)}} d\tau\right)(\xi) \\ &= \mathcal{F}\left(\varphi * \frac{1}{\sqrt{4k\pi t}} e^{-\frac{(\cdot)^2}{4kt}} + \int_0^t f(\cdot, \tau) * \frac{1}{\sqrt{4k\pi(t-\tau)}} e^{-\frac{(\cdot)^2}{4k(t-\tau)}} d\tau\right)(\xi) \end{aligned}$$

The inversion theorem - i.e. property (5) - implies two continuous functions whose Fourier transforms agree for every $\xi \in \mathbb{R}$ are actually identical functions of $x \in \mathbb{R}$. (This result is known as the uniqueness theorem for Fourier transforms.) Therefore

$$u(x, t) = \left(\varphi * \frac{1}{\sqrt{4k\pi t}} e^{-\frac{(\cdot)^2}{4kt}} + \int_0^t f(\cdot, \tau) * \frac{1}{\sqrt{4k\pi(t-\tau)}} e^{-\frac{(\cdot)^2}{4k(t-\tau)}} d\tau \right)(x)$$

for all $x \in \mathbb{R}$ (and $t > 0$). Writing out the convolutions in terms of integrals gives (**):

$$u(x, t) = \int_{-\infty}^{\infty} \frac{1}{\sqrt{4k\pi t}} e^{-\frac{(x-y)^2}{4kt}} \varphi(y) dy + \int_0^t \int_{-\infty}^{\infty} \frac{1}{\sqrt{4k\pi(t-\tau)}} e^{-\frac{(x-y)^2}{4k(t-\tau)}} f(y, \tau) dy d\tau.$$