

1.(25 pts.) (a) Let $u = u(x, t)$ be a classical solution to the diffusion equation $u_t - u_{xx} = 0$ in the strip $0 < x < 1, 0 < t < \infty$, such that $u(0, t) = 0 \stackrel{\text{②}}{=} u_x(1, t)$ if $t \geq 0$. Show that the energy of the solution u ,

given by $E(t) = \int_0^1 u^2(x, t) dx$, is a nonincreasing function on the interval $0 \leq t < \infty$.

$$\begin{aligned} \frac{dE}{dt} &= \frac{d}{dt} \int_0^1 u^2(x, t) dx = \int_0^1 \frac{\partial}{\partial t} [u^2(x, t)] dx = \int_0^1 2u(x, t) u_t(x, t) dx = \int_0^1 \overbrace{2u(x, t)}^{\text{by ①}} \overbrace{u_{xx}(x, t)}^{dV} dx \\ &= 2u(x, t) u_x(x, t) \Big|_{x=0}^1 - \int_0^1 2u_x(x, t) u_x(x, t) dx = 2u(1, t) u_x(1, t) - 2u(0, t) u_x(0, t) - 2 \int_0^1 u_x^2(x, t) dx \\ &= -2 \int_0^1 u_x^2(x, t) dx \leq 0. \end{aligned}$$

8 pts.
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14 pts.
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15 pts.
to here.

Therefore $E = E(t)$ is a nonincreasing function on the interval $0 \leq t < \infty$.

(b) Use part (a) to help show that there is at most one classical solution to the following problem.

$$(*) \quad \begin{cases} u_t - u_{xx} = f(x, t) & \text{if } 0 < x < 1, 0 < t < \infty, \\ u(0, t) = g(t) \text{ and } u_x(1, t) = h(t) & \text{if } 0 \leq t < \infty, \\ u(x, 0) = \phi(x) & \text{if } 0 \leq x \leq 1. \end{cases}$$

(Here f, g, h , and ϕ are fixed continuous functions.)

2 pts.
to here.

Suppose $u_1 = u_1(x, t)$ and $u_2 = u_2(x, t)$ are classical solutions to $(*)$. Then $v(x, t) = u_1(x, t) - u_2(x, t)$ is a classical solution to

$$(***) \quad \begin{cases} v_t - v_{xx} = 0 & \text{if } 0 < x < 1, 0 < t < \infty, \\ v(0, t) = 0 = v_x(1, t) & \text{if } 0 \leq t < \infty, \\ v(x, 0) = 0 & \text{if } 0 \leq x \leq 1. \end{cases}$$

4 pts.
to here.

If $E(t) = \int_0^1 v^2(x, t) dx$ ($t \geq 0$) is the energy of the solution v , then E is a non-increasing function on the interval $[0, \infty)$. But $0 \leq E(t) \leq E(0) = \int_0^1 v^2(x, 0) dx = 0$ for all $t \geq 0$ so by the vanishing theorem, the nonnegative continuous integrand of E must be identically zero for each $t \geq 0$ and $0 \leq x \leq 1$. I.e. $v(x, t) \equiv 0$ so $u_1 = u_2$.

6 pt.
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8 pts.
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10 pts. to here.

2.(25 pts.) Solve the heat equation in the upper half-plane

$$u_t - u_{xx} = 0 \text{ if } -\infty < x < \infty, 0 < t < \infty,$$

subject to the initial condition

$$u(x, 0) = x^2 \text{ if } -\infty < x < \infty.$$

Note: You may find the following facts useful.

$$\int_{-\infty}^{\infty} pe^{-p^2} dp = 0 \quad \text{and} \quad \int_{-\infty}^{\infty} p^2 e^{-p^2} dp = \frac{\sqrt{\pi}}{2}.$$

Although $\varphi(x) = x^2$ is not bounded on $-\infty < x < \infty$, we use the formula for solutions to the diffusion I.V.P. in the upper half plane:

$$u(x, t) = \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4t}} \varphi(y) dy \quad (t > 0).$$

4 pts.
to here.

Since $k=1$ and $\varphi(x) = x^2$ this becomes

$$\begin{aligned} u(x, t) &= \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4t}} y^2 dy \quad \dots \quad \text{Let } p = \frac{y-x}{\sqrt{4t}}. \text{ Then } dp = \frac{dy}{\sqrt{4t}}. \\ &= \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} (x + p\sqrt{4t})^2 dp \\ &= \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} (x^2 + 2px\sqrt{4t} + 4tp^2) dp \\ &= \frac{x^2}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-p^2} dp + \frac{2x\sqrt{4t}}{\sqrt{\pi}} \int_{-\infty}^{\infty} pe^{-p^2} dp + \frac{4t}{\sqrt{\pi}} \int_{-\infty}^{\infty} p^2 e^{-p^2} dp \end{aligned}$$

$$\boxed{u(x, t) = x^2 + 2t}$$

23 pts.
to here.

Since $\varphi(x) = x^2$ is not bounded on $-\infty < x < \infty$, we have "abused" the formula for solutions to the diffusion I.V.P. in the upper halfplane and must check our answer

$$u_t - u_{xx} = 2 - 2 = 0 \quad \checkmark$$

$$u(x, 0) = x^2 \quad \checkmark$$

25 pts.
to here.

3.(25 pts.) Use the Fourier transform method to derive a formula for the solution to the following initial value problem.

$$\begin{aligned} u_t - u_{xx} + tu &= 0 \quad \text{if } -\infty < x < \infty, 0 < t < \infty, \\ u(x, 0) &\stackrel{(2)}{=} f(x) \quad \text{if } -\infty < x < \infty. \end{aligned}$$

Taking the Fourier transform of ① with respect to x yields $\mathcal{F}(u_t - u_{xx} + tu)(\xi) = \mathcal{F}(f)(\xi)$

so $\mathcal{F}(u_t)(\xi) - \mathcal{F}(u_{xx})(\xi) + t\mathcal{F}(u)(\xi) = 0$ by linearity. But $\mathcal{F}(f')(\xi) = i\xi\mathcal{F}(f)(\xi)$

and $\mathcal{F}(u_t)(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial u}{\partial t}(x, t) e^{-i\xi x} dx = \frac{\partial}{\partial t} \left(\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u(x, t) e^{-i\xi x} dx \right) = \frac{\partial}{\partial t} \mathcal{F}(u)(\xi)$ so

$$\frac{\partial}{\partial t} \mathcal{F}(u)(\xi) - (i\xi)^2 \mathcal{F}(u)(\xi) + t\mathcal{F}(u)(\xi) = 0$$

$$(†) \quad \frac{\partial}{\partial t} \mathcal{F}(u)(\xi) + (t + \xi^2) \mathcal{F}(u)(\xi) = 0.$$

An integrating factor for this linear first-order ODE in the ^{independent} variable t (and parameter ξ) is

$$\mu(t) = e^{\int (t + \xi^2) dt} = e^{\frac{t^2}{2} + \xi^2 t}.$$

Multiplying (†) by the integrating factor, we obtain an "exact" expression in the left member:

$$\frac{\partial}{\partial t} \left(e^{\frac{t^2}{2} + \xi^2 t} \mathcal{F}(u)(\xi) \right) = e^{\frac{t^2}{2} + \xi^2 t} \frac{\partial}{\partial t} \mathcal{F}(u)(\xi) + (t + \xi^2) e^{\frac{t^2}{2} + \xi^2 t} \mathcal{F}(u)(\xi) = 0.$$

Integrating with respect to t (holding ξ fixed) yields

$$e^{\frac{t^2}{2} + \xi^2 t} \mathcal{F}(u)(\xi) = c(\xi) \Rightarrow \mathcal{F}(u)(\xi) = c(\xi) e^{-\frac{t^2}{2} - \xi^2 t}$$

Applying the initial condition ② yields $\mathcal{F}(f)(\xi) = \mathcal{F}(u(\cdot, 0))(\xi) = c(\xi)$ so

(*) $\mathcal{F}(u)(\xi) = \mathcal{F}(f)(\xi) e^{-\frac{t^2}{2}} \cdot e^{-\xi^2 t}$. Applying formula I in the table of Fourier

transforms: $\mathcal{F}(e^{-ax^2})(\xi) = \frac{1}{\sqrt{2a}} e^{-\frac{|x|}{4a}}$ with $\frac{1}{4a} = t$, we find

$\mathcal{F}\left(\frac{1}{\sqrt{2t}} e^{-\frac{x^2}{4t}}\right)(\xi) = \mathcal{F}\left(\sqrt{2a} e^{-ax^2}\right)(\xi) = e^{-\frac{\xi^2}{4a}} = e^{-\xi^2 t}$. Substituting

into (*) we have $\mathcal{F}(u)(\xi) = e^{-\frac{t^2}{2}} \mathcal{F}(f)(\xi) \mathcal{F}\left(\frac{1}{\sqrt{2t}} e^{-\frac{x^2}{4t}}\right)(\xi)$. (OVER)

But $\mathcal{F}(f * g)(\xi) = \sqrt{2\pi} \mathcal{F}(f)(\xi) \mathcal{F}(g)(\xi)$ so the last expression is equivalent to

$$\text{to here. } \mathcal{F}(u)(\xi) = \frac{e^{-t^2/2}}{\sqrt{2\pi}} \mathcal{F}\left(f * e^{-\frac{(\cdot)^2}{4t}}\right)(\xi) = \mathcal{F}\left(f * \frac{1}{\sqrt{4\pi t}} e^{-\frac{t^2}{2} - \frac{(\cdot)^2}{4t}}\right)(\xi).$$

By the uniqueness theorem for Fourier transforms, it follows that

$$\text{to here. } u(x, t) = \left(f * \frac{1}{\sqrt{4\pi t}} e^{-\frac{t^2}{2} - \frac{(x-y)^2}{4t}} \right)(x)$$

or equivalently

$$u(x, t) = \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-\frac{t^2}{2} - \frac{(x-y)^2}{4t}} f(y) dy$$

$$\boxed{u(x, t) = \frac{e^{-t^2/2}}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4t}} f(y) dy}$$

for $t > 0$ and $-\infty < x < \infty$.

4.(25 pts.) Solve the diffusion problem in the strip

$$u_t - u_{xx} \stackrel{(1)}{=} 0 \text{ if } 0 < x < 1, 0 < t < \infty,$$

subject to the mixed boundary conditions (Dirichlet at the left end, Neumann at the right end)

$$u(0, t) \stackrel{(2)}{=} 0 \stackrel{(3)}{=} u_x(1, t) \text{ if } 0 \leq t < \infty,$$

and the initial condition

$$u(x, 0) \stackrel{(4)}{=} 3 \sin\left(\frac{\pi x}{2}\right) - \sin\left(\frac{3\pi x}{2}\right) \text{ if } 0 \leq x \leq 1.$$

(Note: The solution to this problem is unique by problem 1 on this exam.)

We use separation of variables. We seek nontrivial solutions of ①-②-③ of the form

$$u(x, t) = \Xi(x)T(t). \quad \text{Substituting into ① and rearranging we have } -\frac{T'(t)}{T(t)} = -\frac{\Xi''(x)}{\Xi(x)} = \lambda$$

for some constant λ and all $0 < x < 1, 0 < t < \infty$. Substituting into ② and ③ yields

$$\Xi(0)T(t) = 0 \text{ and } \Xi'(1)T(t) = 0 \text{ for all } t \geq 0. \quad \text{Since we require that } T(t) \neq 0$$

it follows that $\Xi(0) = 0 = \Xi'(1)$. Therefore we are led to the coupled system of ODE's

$$\begin{cases} \boxed{\Xi''(x) + \lambda \Xi(x) = 0 \text{ if } 0 < x < 1 \text{ with } \Xi(0) = 0 = \Xi'(1)}, \\ \text{and} \\ T'(t) + \lambda T(t) = 0 \text{ if } t > 0. \end{cases} \quad \text{Eigenvalue Problem}$$

We assume that the eigenvalues λ above are real.

Case $\lambda = 0$: Then $\Xi''(x) = 0$ has general solution $\Xi(x) = c_1 x + c_2$. $0 = \Xi(0) = c_2$

and $0 = \Xi'(1) = c_1$, so there are no nontrivial solutions in this case; i.e. $\lambda = 0$ is not an eigenvalue.

Case $\lambda < 0$ (say $\lambda = -\mu^2$ where $\mu > 0$): Then $\Xi''(x) - \mu^2 \Xi(x) = 0$ has general solution $\Xi(x) = c_1 \cosh(\mu x) + c_2 \sinh(\mu x)$. Then $0 = \Xi(0) = c_1$, and $0 = \Xi'(1) = \mu c_2 \cosh(\mu)$ imply $c_1 = c_2 = 0$. I.e. there are no nontrivial solutions in this case and hence there are no negative eigenvalues.

Case $\lambda > 0$ (say $\lambda = \mu^2$ where $\mu > 0$): Then $\Xi''(x) + \mu^2 \Xi(x) = 0$ has general solution $\Xi(x) = c_1 \cos(\mu x) + c_2 \sin(\mu x)$. Also $0 = \Xi(0) = c_1$, and $0 = \Xi'(1) = \mu c_2 \cos(\mu)$

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13 pts.
to here.

imply either $c_1 = c_2 = 0$ or $c_1 = 0$ and $\cos(\mu) = 0$. Therefore in order to have a nontrivial solution in this case, we must have $\mu_n = (n + \frac{1}{2})\pi$ for some integer $n = 0, 1, 2, \dots$. Hence the eigenvalues are

15 pts.
to here.

$$\lambda_n = \mu_n^2 = (n + \frac{1}{2})\pi^2 \quad (n = 0, 1, 2, \dots)$$

and the corresponding eigenfunctions are (up to a constant factor)

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$$X_n(x) = \sin(\mu_n x) = \sin((n + \frac{1}{2})\pi x) \quad (n = 0, 1, 2, \dots)$$

The corresponding solutions to the t-equation,

$$T_n'(t) + \lambda_n T_n(t) = 0 \quad \text{or equivalently} \quad T_n'(t) + (n + \frac{1}{2})\pi^2 T_n(t) = 0,$$

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are (up to a constant factor) $T_n(t) = e^{-(n + \frac{1}{2})\pi^2 t}$. Therefore

$$u_n(x, t) = X_n(x) T_n(t) = e^{-(n + \frac{1}{2})\pi^2 t} \sin((n + \frac{1}{2})\pi x) \quad (n = 0, 1, 2, \dots)$$

are nontrivial solutions to ①-②-③. By the superposition principle

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$$u(x, t) = \sum_{n=0}^N b_n e^{-(n + \frac{1}{2})\pi^2 t} \sin((n + \frac{1}{2})\pi x)$$

also solves ①-②-③ for any integer $N \geq 0$ and any choice of constants b_0, b_1, \dots, b_N .

In order to satisfy ④ as well, we must have

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$$3 \sin\left(\frac{\pi x}{2}\right) - \sin\left(\frac{3\pi x}{2}\right) \stackrel{④}{=} u(x, 0) = \sum_{n=0}^N b_n \sin\left((n + \frac{1}{2})\pi x\right) \quad \text{for all } 0 \leq x \leq 1.$$

By inspection, we may take $N = 1$ and $b_0 = 3, b_1 = -1$ to satisfy ④. Thus

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pts. to here.

$$\boxed{u(x, t) = 3e^{-\frac{\pi^2}{4}t} \sin\left(\frac{\pi x}{2}\right) - e^{-\frac{9\pi^2}{4}t} \sin\left(\frac{3\pi x}{2}\right)}$$

solves ①-②-③-④.

A Brief Table of Fourier Transforms

$$f(x) \quad \hat{f}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ix\xi} dx$$

A.	$\begin{cases} 1 & \text{if } -b < x < b, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{2}{\pi} \frac{\sin(b\xi)}{\xi}$
B.	$\begin{cases} 1 & \text{if } c < x < d, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{e^{-ic\xi} - e^{-id\xi}}{i\xi\sqrt{2\pi}}$
C.	$\frac{1}{x^2 + a^2} \quad (a > 0)$	$\frac{\pi}{2} \frac{e^{-a \xi }}{a}$
D.	$\begin{cases} x & \text{if } 0 < x \leq b, \\ 2b - x & \text{if } b < x < 2b, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{-1 + 2e^{-ib\xi} - e^{-2ib\xi}}{\xi^2\sqrt{2\pi}}$
E.	$\begin{cases} e^{-ax} & \text{if } x > 0, \\ 0 & \text{otherwise.} \end{cases} \quad (a > 0)$	$\frac{1}{(a + i\xi)\sqrt{2\pi}}$
F.	$\begin{cases} e^{ax} & \text{if } b < x < c, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{e^{(a-i\xi)c} - e^{(a-i\xi)b}}{(a - i\xi)\sqrt{2\pi}}$
G.	$\begin{cases} e^{iax} & \text{if } -b < x < b, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{2}{\pi} \frac{\sin(b(\xi-a))}{\xi - a}$
H.	$\begin{cases} e^{iax} & \text{if } c < x < d, \\ 0 & \text{otherwise.} \end{cases}$	$\frac{i}{\sqrt{2\pi}} \frac{e^{ic(a-\xi)} - e^{id(a-\xi)}}{a - \xi}$
I.	$e^{-ax^2} \quad (a > 0)$	$\frac{1}{\sqrt{2a}} e^{-\xi^2/(4a)}$
J.	$\frac{\sin(ax)}{x} \quad (a > 0)$	$\begin{cases} 0 & \text{if } \xi \geq a, \\ \sqrt{\pi/2} & \text{if } \xi < a. \end{cases}$

Math 325
Exam II
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mean: 78.1

standard deviation: 17.6

number of scores: 16

<u>Distribution of Scores:</u>		<u>frequency</u>
87 - 100	A	6
73 - 86	B	4
60 - 72	C(graduate), B(undergraduate)	4
50 - 59	C	1
0 - 49	F(graduate), D(undergraduate)	1