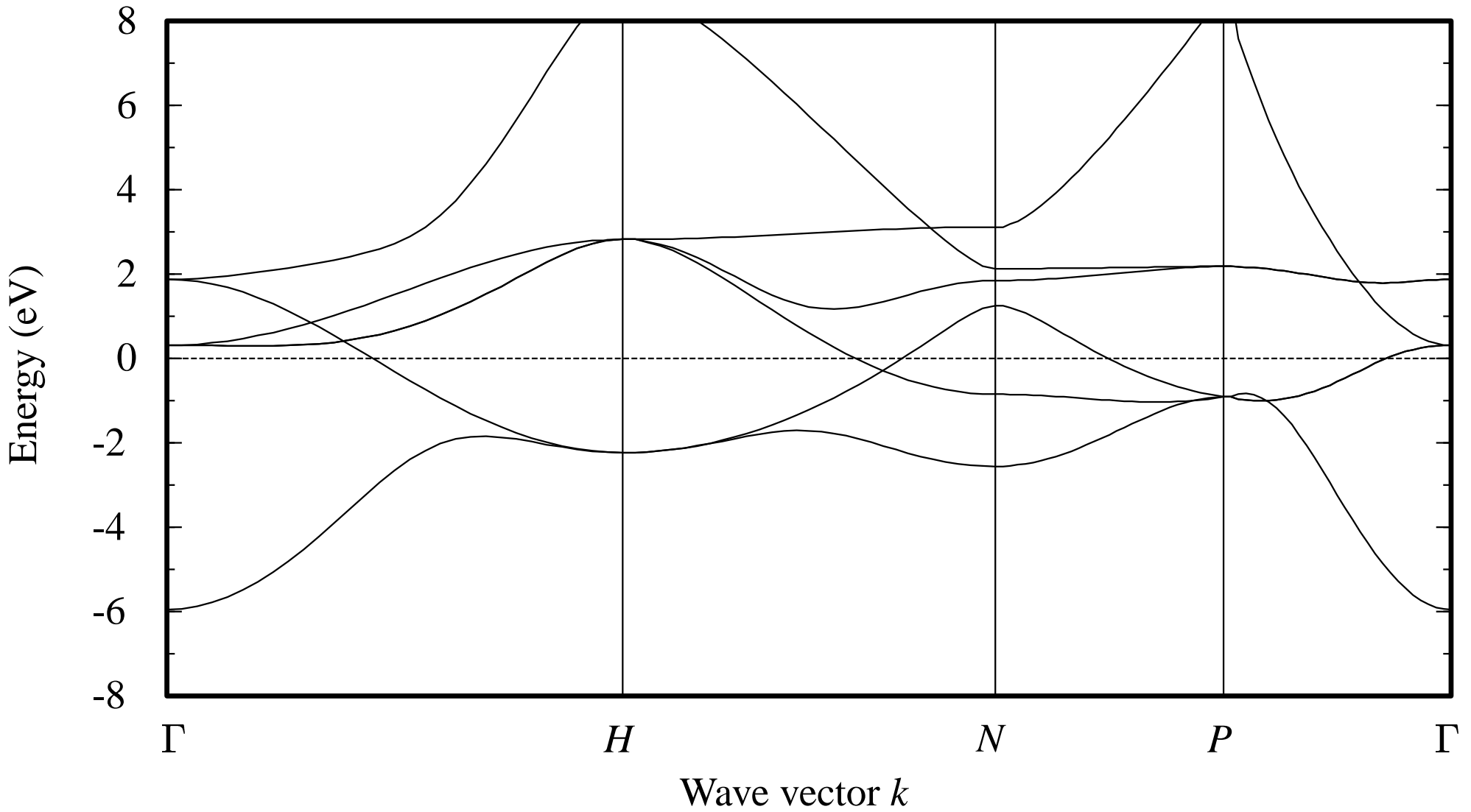


Band Structure Calculations



Question: How could it ever be true that electrons in a metal think they are moving freely in an empty box?

Pseudopotentials give conceptual answer.

- ➡ Restrict attention to single unit cell.
- ➡ Let $|\vec{k}\rangle$ denote plane waves $e^{i\vec{k}\cdot\vec{r}}$.
- ➡ Let $|\psi_c\rangle$ denote core states.

$$|\vec{k}_{ps}\rangle = |\vec{k}\rangle - \sum_c |\psi_c\rangle \langle \psi_c | \vec{k}\rangle, \quad (\text{L1})$$

$$\hat{U}|\vec{k}_{ps}\rangle = \hat{U}|\vec{k}\rangle - \sum_c \hat{U} \langle \psi_c | \vec{k}\rangle |\psi_c\rangle. \quad (\text{L2})$$

$$\begin{aligned} (\hat{\mathcal{H}} - \mathcal{E})|\vec{k}_{ps}\rangle &= \left(\frac{\hat{P}^2}{2m} + \hat{U} - \mathcal{E} \right) |\vec{k}_{ps}\rangle \\ &= ? \end{aligned} \quad (\text{L3})$$

Pseudopotentials

$$? = \left(\frac{\hat{P}^2}{2m} + \hat{U}_{\text{ps}} - \mathcal{E} \right) |\vec{k}\rangle = (\hat{\mathcal{H}}_{\text{ps}} - \mathcal{E}) |\vec{k}\rangle, \quad (\text{L6})$$

$$\hat{U}_{\text{ps}} = \hat{U} - \sum_c (\mathcal{E}_c - \mathcal{E}) |\psi_c\rangle \langle \psi_c|. \quad (\text{L7})$$

$$(\hat{\mathcal{H}} - \mathcal{E}) |\vec{k}_{\text{ps}}\rangle = (\hat{\mathcal{H}}_{\text{ps}} - \mathcal{E}) |\vec{k}\rangle. \quad (\text{L8})$$

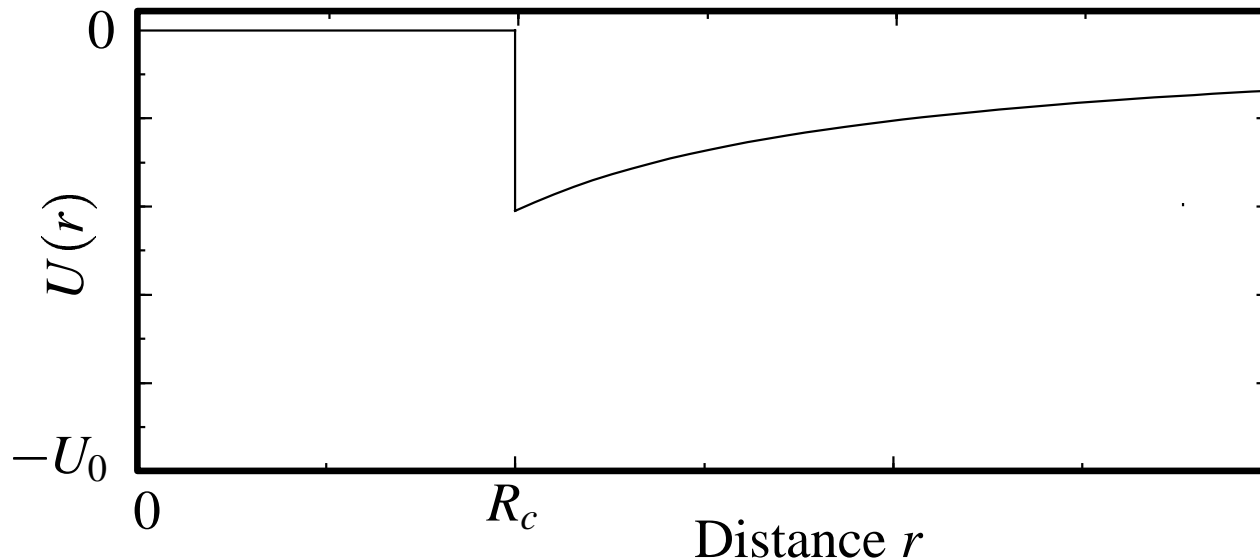


Figure 1: The Ashcroft empty core pseudopotential is zero up to a critical radius R_c , and it equals a screened Coulomb potential $-U_0 \exp[-r/d]/r$ thereafter.

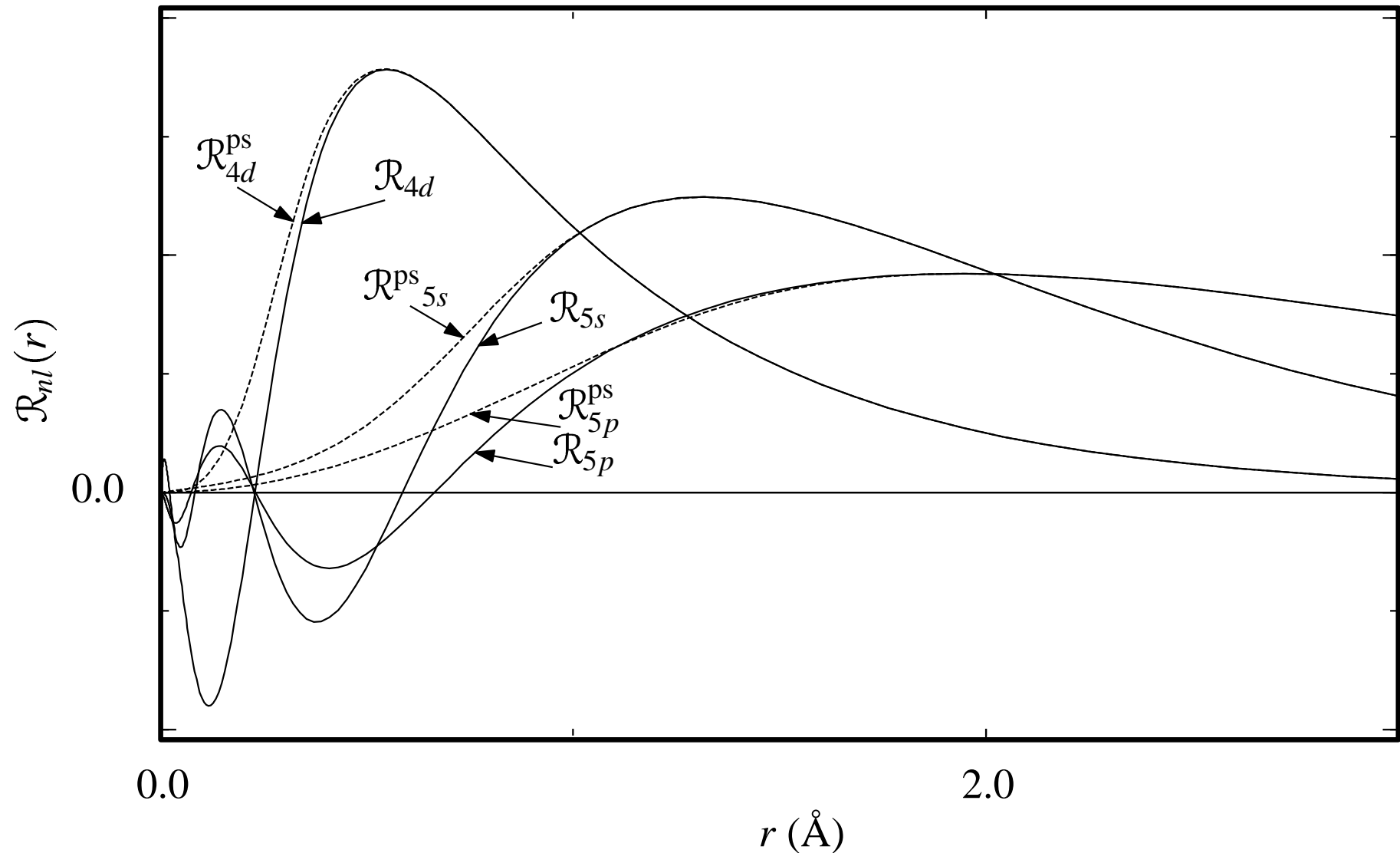


Figure 2: Real and pseudo wave functions for the 5s, 5p, and 4d levels of silver. The 5p level is not much occupied in the ground state of silver, but it can be included in the pseudopotential nevertheless.

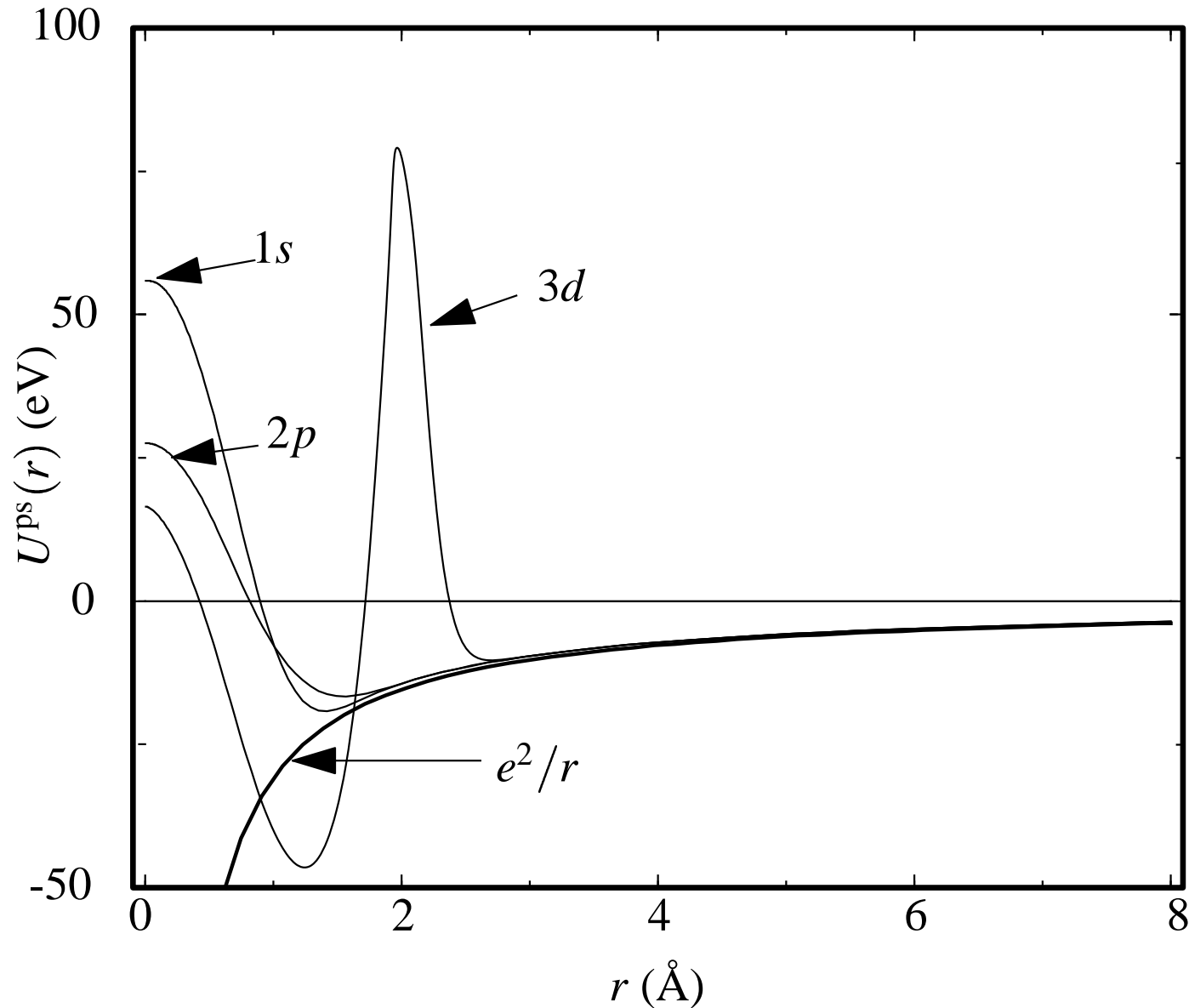


Figure 3: Pseudopotentials for the $5s$, $5p$, and $4d$ states of silver.

Electron density $n(\vec{r}) = \sum |\psi_i(\vec{r})|^2$ is spherically symmetrical in vicinity of nucleus

Equation for radial functions is

$$-\frac{\hbar^2}{2m} \left[\frac{1}{r} \frac{\partial^2}{\partial r^2} r - \frac{l(l+1)}{r^2} \right] \mathcal{R}_{nl} + \left[\int \frac{e^2 n(r')}{|\vec{r} - \vec{r}'|} d\vec{r}' - \frac{e^2 Z}{r} + \frac{\delta \mathcal{E}_{xc}}{\delta n} - \mathcal{E}_{nl} \right] \mathcal{R}_{nl}(r) = 0. \quad (\text{L9})$$

$$U_l^{\text{ps}}(r) = \frac{\hbar^2}{2m} \left[\frac{1}{r \mathcal{R}_{nl}^{\text{ps}}} \frac{\partial^2 r \mathcal{R}_{nl}^{\text{ps}}}{\partial r^2} - \frac{l(l+1)}{r^2} \right] - \left[\int \frac{e^2 n^{\text{ps}}(r')}{|\vec{r} - \vec{r}'|} d\vec{r}' + \frac{\delta \mathcal{E}_{xc}}{\delta n^{\text{ps}}} - \mathcal{E}_{nl} \right]. \quad (\text{L10})$$

$$\psi(\vec{r}) = \sum_{lm} Y_{lm}(\theta, \phi) \psi_{lm}(r); \quad \psi_{lm}(r) = \int d\theta d\phi \sin\theta Y_{lm}^*(\theta, \phi) \psi(\vec{r}), \quad (\text{L11})$$

$$U^{\text{ps}} = \frac{4\pi Z e^2}{q^2 + \kappa^2}. \quad (\text{L12})$$

Result from later work....

$$\frac{1}{\Omega} U^{\text{ps}}(q=0) = -\frac{2}{3} \mathcal{E}_F. \quad (\text{L13})$$

$$\psi_{\vec{k}}(\vec{r}) = \frac{1}{\sqrt{N}} \sum_{\vec{R}, l} b_l e^{i\vec{k} \cdot \vec{R}} a_l^{\text{at}}(\vec{r} - \vec{R}), \quad (\text{L14})$$

$$\langle \psi | \hat{\mathcal{H}} - \mathcal{E} | \psi \rangle. \quad (\text{L15})$$

$$\langle \psi | \mathcal{E} | \psi \rangle = \mathcal{E} \sum_{\vec{R}, \vec{R}'} \int d\vec{r} a^{\text{at}}(\vec{r} - \vec{R}) a^{\text{at}}(\vec{r} - \vec{R}') \frac{e^{i\vec{k} \cdot (\vec{R} - \vec{R}')}}{N} b^2 \quad (\text{L16})$$

$$= b^2 \mathcal{E} \left(1 + \sum_{\vec{\delta}} \alpha e^{i\vec{k} \cdot \vec{\delta}} \right) \quad (\text{L17})$$

and

$$\alpha = \int d\vec{r} a^{\text{at}}(\vec{r}) a^{\text{at}}(\vec{r} + \vec{\delta}). \quad (\text{L18})$$

$$\langle \psi | \hat{\mathcal{H}} | \psi \rangle = \sum_{\vec{R}, \vec{R}'} \int d\vec{r} a^{\text{at}}(\vec{r} - \vec{R}) \left[-\frac{\hbar^2}{2m} \nabla^2 + U(\vec{r}) \right] a^{\text{at}}(\vec{r} - \vec{R}') \frac{e^{i\vec{k} \cdot (\vec{R} - \vec{R}')}}{N} b^2 \quad (\text{L19})$$

$$= \sum_{\vec{R}\vec{R}'} \int d\vec{r} a^{\text{at}}(\vec{r} - \vec{R}) \left\{ \begin{array}{l} [-\frac{\hbar^2}{2m} \nabla^2 + U^{\text{at}}(\vec{r} - \vec{R}')] a^{\text{at}}(\vec{r} - \vec{R}') \\ + [U(\vec{r}) - U^{\text{at}}(\vec{r} - \vec{R}')] a^{\text{at}}(\vec{r} - \vec{R}') \end{array} \right\} \frac{e^{i\vec{k} \cdot (\vec{R} - \vec{R}')}}{N} b^2 \quad (\text{L20})$$

$$= \int d\vec{r} \sum_{\vec{R}\vec{R}'} \mathcal{E}^{\text{at}} \frac{a^{\text{at}}(\vec{r} - \vec{R}) a^{\text{at}}(\vec{r} - \vec{R}')}{N} e^{i\vec{k} \cdot (\vec{R} - \vec{R}')} b^2$$

$$+ \int d\vec{r} \sum_{\vec{R}\vec{R}'} a^{\text{at}}(\vec{r} - \vec{R}) [U(\vec{r}) - U^{\text{at}}(\vec{r} - \vec{R}')] a^{\text{at}}(\vec{r} - \vec{R}') \frac{e^{i\vec{k} \cdot (\vec{R} - \vec{R}')}}{N} b^2. \quad (\text{L21})$$

$$\mathcal{E}(1 + \sum_{\vec{\delta}} \alpha e^{i\vec{k} \cdot \vec{\delta}}) = \mathcal{E}^{\text{at}} \sum_{\vec{\delta}} \alpha e^{i\vec{k} \cdot \vec{\delta}} + U + (\mathfrak{t} - \alpha \mathcal{E}^{\text{at}}) \sum_{\vec{\delta}} e^{i\vec{k} \cdot \vec{\delta}}, \quad (\text{L22})$$

where

$$U = \mathcal{E}^{\text{at}} + \int d\vec{r} a^{\text{at}}(\vec{r}) [U(\vec{r}) - U^{\text{at}}(\vec{r})] a^{\text{at}}(\vec{r}) \quad (\text{L23})$$

and

$$\mathfrak{t} = \alpha \mathcal{E}^{\text{at}} + \int d\vec{r} a^{\text{at}}(\vec{r}) [U(\vec{r}) - U^{\text{at}}(\vec{r} + \vec{\delta})] a^{\text{at}}(\vec{r} + \vec{\delta}). \quad (\text{L24})$$

$$\mathcal{E} = U + t \sum_{\vec{\delta}} e^{i\vec{k} \cdot \vec{\delta}}. \quad (\text{L25})$$

$$\vec{K} = l_1 \vec{b}_1 + l_2 \vec{b}_2 + l_3 \vec{b}_3 \quad (\text{L26})$$

$$\sum_{i=1}^N \lambda_i^r \hat{e}_i (\hat{e}_i \cdot \vec{a}_1). \quad (\text{L27})$$

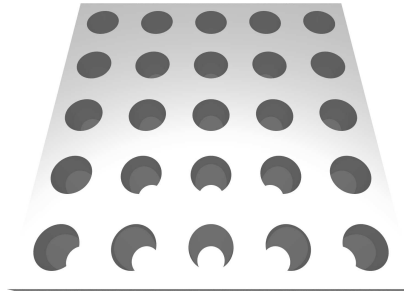
$$\hat{\mathcal{H}} = \frac{\hat{P}^2}{2m} + U(\hat{R}). \quad (\text{L28})$$

$$\hbar^2 K_{\max}^2 / 2m = \mathcal{E}_{\max}$$

$$\hat{\mathcal{H}}\psi = \sum_{\vec{K}'} \left\{ \left[\mathcal{E}_{\vec{k}+\vec{K}'}^0 - \mathcal{E}_{\max} \right] \delta_{\vec{K}\vec{K}'} + U_{\vec{K}-\vec{K}'} \right\} \psi_{n\vec{k}}(\vec{K}'), \quad (\text{L29})$$

$$1 + \hat{\mathcal{H}} dt / \hbar. \quad (\text{L30})$$

$$\psi_{n+1} = (1 + \hat{\mathcal{H}}dt/\hbar)\psi_n \Rightarrow \frac{\psi_{n+1} - \psi_n}{dt} = \frac{1}{\hbar}\hat{\mathcal{H}}\psi_n, \quad (\text{L31})$$



$$\phi_{\varepsilon\vec{k}} = e^{i\vec{k}\cdot\vec{r}}$$

$$-\frac{1}{2m}\hbar^2\nabla^2\phi_{\varepsilon\vec{k}} + U(r)\phi_{\varepsilon\vec{k}} = \varepsilon\phi_{\varepsilon\vec{k}}$$

$$\psi_{\varepsilon} = Y_{lm}\mathcal{R}_{l\varepsilon}(r), \quad (\text{L32})$$

$$\frac{-\hbar^2}{2mr^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} \mathcal{R}_{l\varepsilon}(r) + [U(r) + \frac{\hbar^2 l(l+1)}{2mr^2}] \mathcal{R}_{l\varepsilon}(r) = \varepsilon \mathcal{R}_{l\varepsilon}(r). \quad (\text{L33})$$

$$\phi_{\varepsilon\vec{k}} = \sum_{l=0}^{\infty} \sum_{m=-l}^l A_{lm} Y_{lm}(\hat{r}) \mathcal{R}_{l\varepsilon}(r), \quad (\text{L34})$$

$$e^{i\vec{k}\cdot\vec{r}} = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^l i^l j_l(kr) Y_{lm}^*(\hat{k}) Y_{lm}(\hat{r}). \quad (\text{L35})$$

$$\phi_{\varepsilon\vec{k}} = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{i^l j_l(kR_h) Y_{lm}^*(\hat{k})}{\mathcal{R}_{l\varepsilon}(R_h)} Y_{lm}(\hat{r}) \mathcal{R}_{l\varepsilon}(r). \quad (\text{L36})$$

$$\psi_{\vec{k}} = \sum_{\vec{K}} b_{\vec{k}+\vec{K}} \phi_{\varepsilon, \vec{k}+\vec{K}}. \quad (\text{L37})$$

$$\langle \psi | \hat{\mathcal{H}} - \varepsilon | \psi \rangle, \quad (\text{L38})$$

$$0 = \sum_{\vec{K}} \langle \phi_{\varepsilon\vec{q}} | \hat{\mathcal{H}} - \varepsilon | \phi_{\varepsilon\vec{q}+\vec{K}} \rangle b_{\vec{q}+\vec{K}}, \quad (\text{L39})$$

where

$$\langle \phi_{\varepsilon\vec{q}} | \hat{\mathcal{H}} - \varepsilon | \phi_{\varepsilon\vec{q}'} \rangle = \left(\frac{\hbar^2 \vec{q} \cdot \vec{q}'}{2m} - \varepsilon \right) \Omega \delta_{\vec{q}, \vec{q}'} + \mathcal{U}_{\vec{q}, \vec{q}'} \quad (\text{L40})$$

$$\mathcal{U}_{\vec{q},\vec{q}'} = 4\pi R_h^2 \left\{ \begin{array}{l} - \left(\frac{\hbar^2 \vec{q} \cdot \vec{q}'}{2m} - \varepsilon \right) \frac{j_1(|\vec{q} - \vec{q}'| R_h)}{|\vec{q} - \vec{q}'|} \\ + \sum_{l=0}^{\infty} \frac{\hbar^2}{2m} (2l+1) P_l(\hat{q} \cdot \hat{q}') j_l(q R_h) j_l(q' R_h) \frac{\mathcal{R}'_{l\varepsilon}(R_h)}{\mathcal{R}_{l\varepsilon}(R_h)} \end{array} \right\}. \quad (\text{L41})$$

$$\chi_{lm}(\mathcal{E}, r) = \begin{cases} i^l Y_l^m(\hat{r}) [\mathcal{R}_{l\mathcal{E}}(r) + (\frac{r}{R_h})^l p_{l\mathcal{E}}] & \text{for } r < R_h \\ i^l Y_l^m(\hat{r}) (\frac{R_h}{r})^{l+1} & \text{for } r > R_h. \end{cases} \quad (\text{L42})$$

$$\begin{aligned} \mathcal{R}_{l\mathcal{E}}(R_h) + p_{l\mathcal{E}} &= 1 \\ \mathcal{R}'_{l\mathcal{E}}(R_h) + l \frac{p_{l\mathcal{E}}}{R_h} &= -\frac{l+1}{R_h}. \end{aligned} \quad (\text{L43})$$

$$p_{l\mathcal{E}} = \frac{-(l+1)/R_h - \left. \frac{\partial \mathcal{R}_{l\mathcal{E}}}{\partial r} \right|_{R_h}}{l/R_h - \left. \frac{\partial \mathcal{R}_{l\mathcal{E}}}{\partial r} \right|_{R_h}} = \frac{\mathcal{D}_l(\mathcal{E}) + l + 1}{\mathcal{D}_l(\mathcal{E}) - l} \quad (\text{L44a})$$

with

$$\mathcal{D}_l(\mathcal{E}) = \frac{R_h}{\mathcal{R}_{l\mathcal{E}}(R_h)} \left. \frac{\partial \mathcal{R}_{l\mathcal{E}}}{\partial r} \right|_{R_h}. \quad (\text{L44b})$$

$$\psi(\vec{r}) = \sum_{lm} B_{lm}^{\vec{k}} \sum_{\vec{R}} e^{i\vec{k} \cdot \vec{R}} \chi_{lm}(\mathcal{E}, \vec{r} - \vec{R}) \quad (\text{L45})$$

$$0 = \sum_{lm} B_{lm}^{n\vec{k}} \left[p_{l\varepsilon} i^l Y_l^m(\hat{r}) \left(\frac{r}{R_h}\right)^l + \sum_{\vec{R} \neq 0} i^l Y_l^m(\widehat{\vec{r} - \vec{R}}) \left(\frac{R_h}{|\vec{r} - \vec{R}|}\right)^{l+1} e^{i\vec{k} \cdot \vec{R}} \right]. \quad (\text{L46})$$

$$\sum_{\vec{R} \neq 0} i^l Y_l^m(\widehat{\vec{r} - \vec{R}}) \left(\frac{R_h}{|\vec{r} - \vec{R}|}\right)^{l+1} e^{i\vec{k} \cdot \vec{R}} = - \sum_{l'm'} \frac{\mathcal{S}_{ll'mm'}^{\vec{k}}}{2(2l'+1)} \left(\frac{r}{R_h}\right)^{l'} i^{l'} Y_{l'}^{m'}(\hat{r}). \quad (\text{L47})$$

$$0 = \sum_{lm} B_{lm}^{n\vec{k}} \left[p_{l\varepsilon} i^l Y_l^m(\hat{r}) \left(\frac{r}{R_h}\right)^l - \sum_{l'm'} \frac{\mathcal{S}_{ll'mm'}^{\vec{k}}}{2(2l'+1)} \left(\frac{r}{R_h}\right)^{l'} i^{l'} Y_{l'}^{m'}(\hat{r}) \right] \quad (\text{L48})$$

$$\Rightarrow 0 = \sum_{lml'm'} B_{lm}^{n\vec{k}} \left[p_{l\varepsilon} \delta_{ll'} \delta_{mm'} - \frac{\mathcal{S}_{ll'mm'}^{\vec{k}}}{2(2l'+1)} \right] \left(\frac{r}{R_h}\right)^{l'} i^{l'} Y_{l'}^{m'}(\hat{r}). \quad (\text{L49})$$

$$0 = \sum_{lm} \left[2(2l+1) p_{l\varepsilon} \delta_{ll'} \delta_{mm'} - \mathcal{S}_{ll'mm'}^{\vec{k}} \right] B_{lm}^{n\vec{k}}. \quad (\text{L50})$$

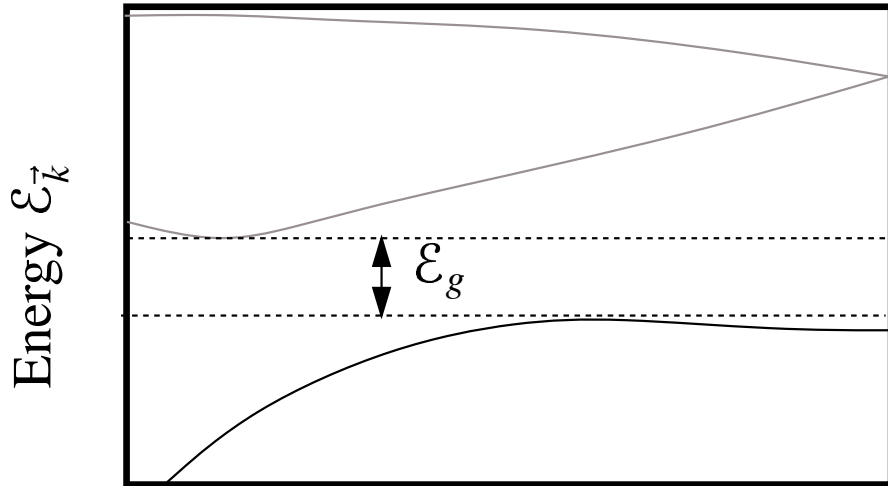
Definition of Metals, Insulators, and Semiconductors

Wilson's theory:

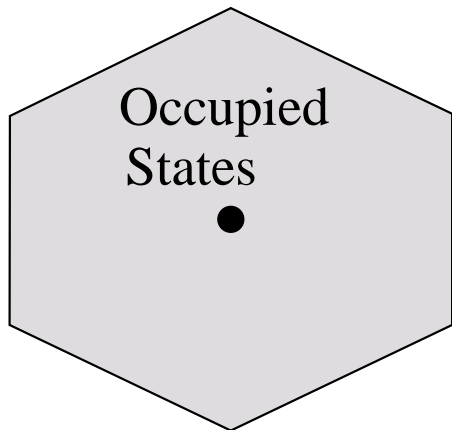
- ➡ **Insulator**: All bands are either full or empty.
- ➡ **Metal**: At least one band is partially occupied.
- ➡ **Semiconductor**: Insulator where energy gap is less than around 1 eV.
- ➡ **Semimetal**: Metal with very small population of conduction electrons.

Definition of Metals, Insulators, and Semiconductors

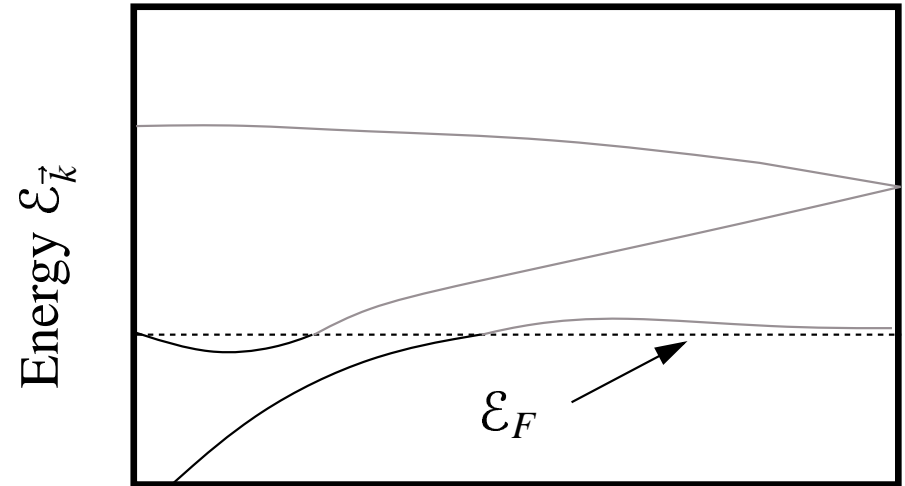
Insulator



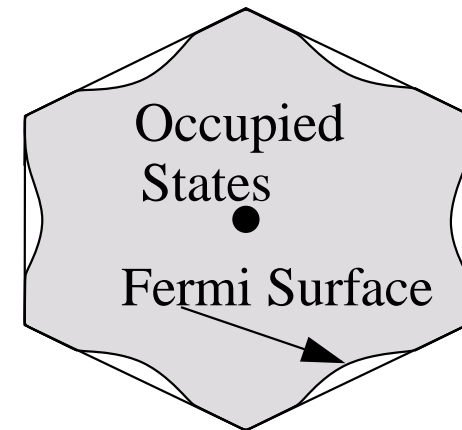
Wave vector k

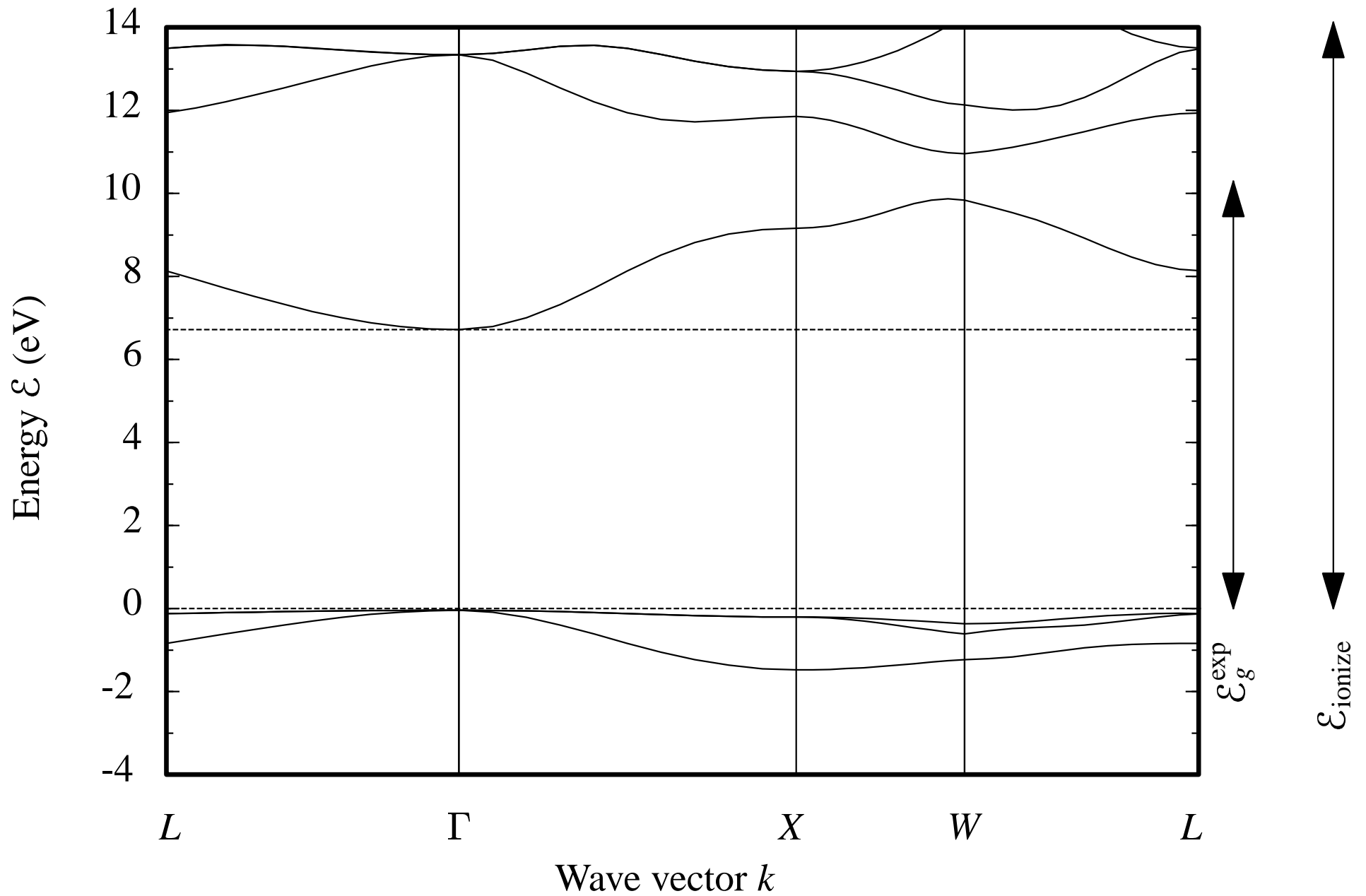


Metal



Wave vector k





Nearly Free Electron Metals

