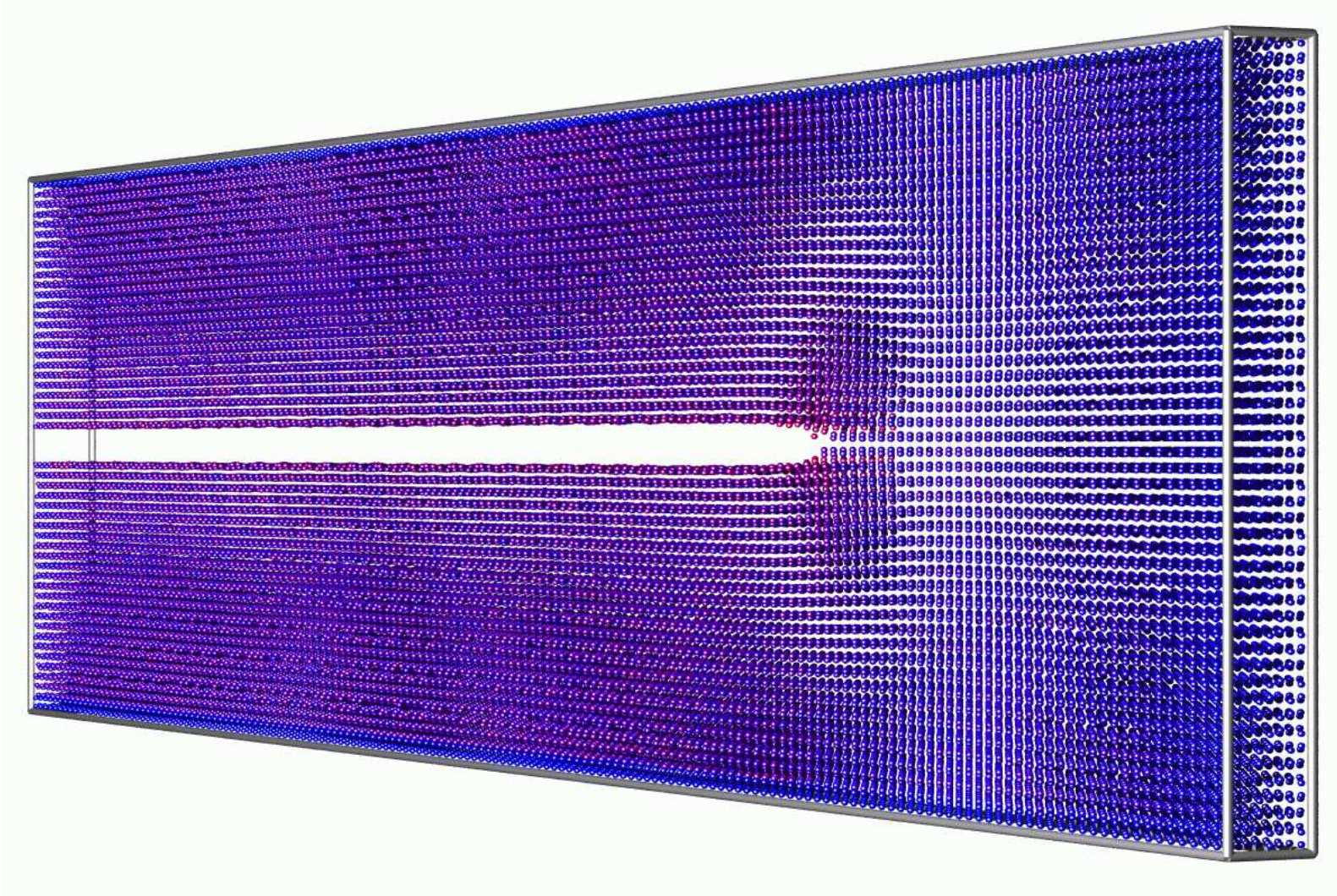


# Dislocations and Cracks



- 
- 
- ➡ Brittle
  - ➡ Ductile
  - ➡ Dislocation
  - ➡ Burgers Vector
  - ➡ Glide Plane
  - ➡ Frenkel–Kontorova Model
  - ➡ Hexatic Phases
  - ➡ Orientational Order, Mermin–Wagner Theorem
  - ➡ Kosterlitz–Thouless–Berezinskii Transition
  - ➡ Cracks
  - ➡ Conformal Mapping
  - ➡ Stress Intensity Factor

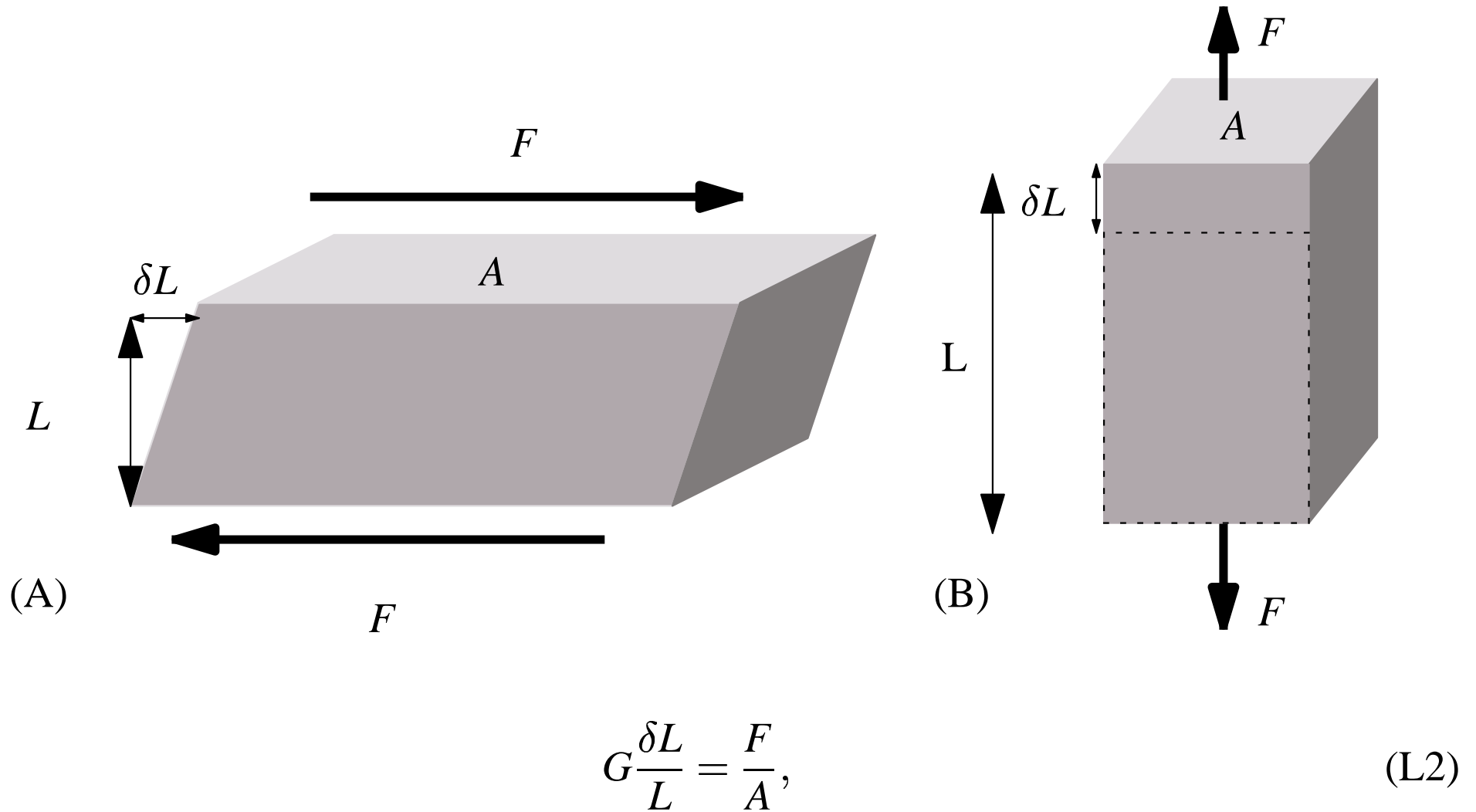




Given surface energy of  $\Gamma = 1 \text{ J/m}^2$ , height  $h$  at which it pays to split object in two is

$$h = \sqrt{\frac{4\Gamma}{\rho g}} \approx 1.4 \text{ cm.} \quad (\text{L1})$$

# Failure in Shear



$$\mathcal{S} = \frac{F}{A} = \begin{cases} \frac{G}{5} & \text{shear} \\ \frac{Y}{5} & \text{tension.} \end{cases} \quad (\text{L3})$$

Material	Shear modulus $G/5$ ( $10^{11}$ ergs $\text{cm}^{-3}$ )	Yield strength ( $10^{11}$ ergs $\text{cm}^{-3}$ )
Iron	1.0–1.6	0.02–1
Copper	1.0	0.005
Titanium	1.0	0.08

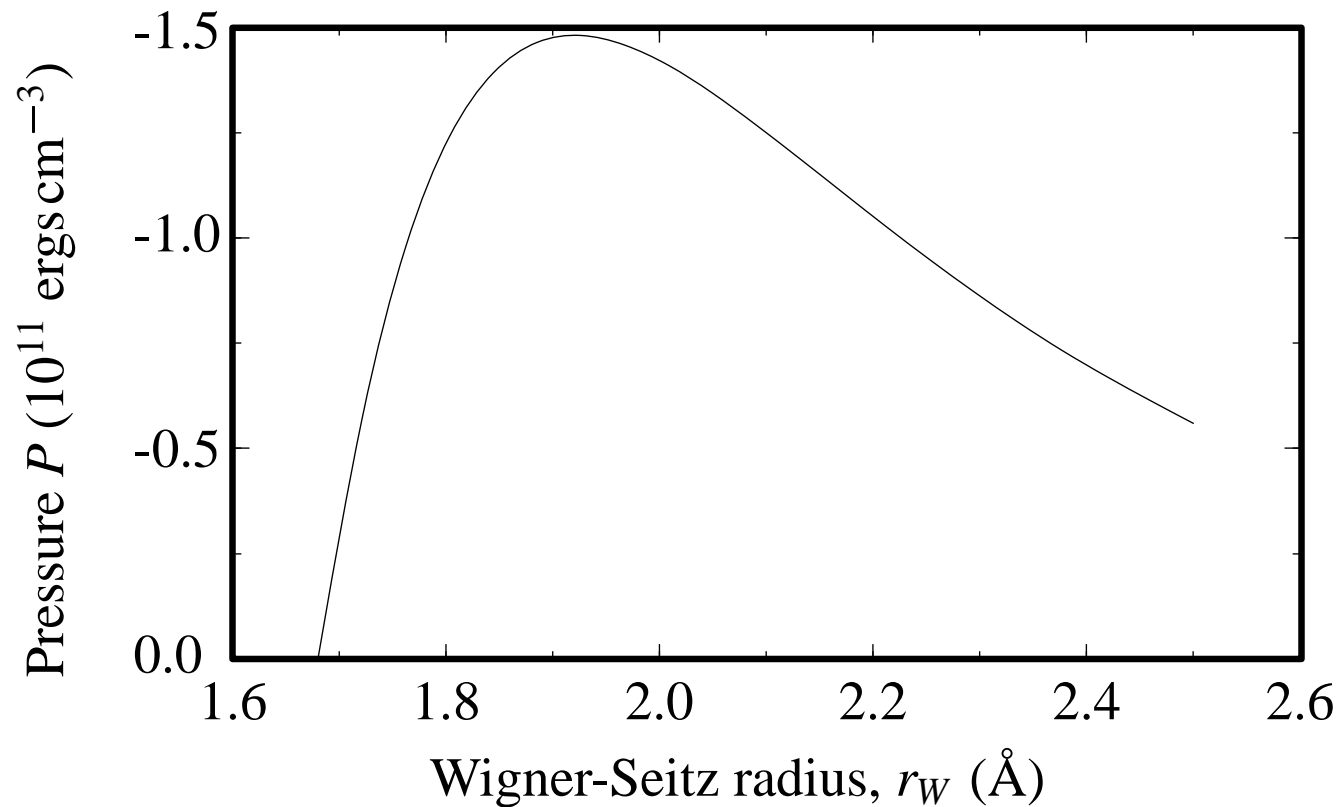
$$\mathcal{S} = \frac{F}{A} = \begin{cases} \frac{G}{5} & \text{shear} \\ \frac{Y}{5} & \text{tension.} \end{cases} \quad (\text{L4})$$

$$Y \frac{\delta L}{L} = \frac{F}{A}. \quad (\text{L5})$$

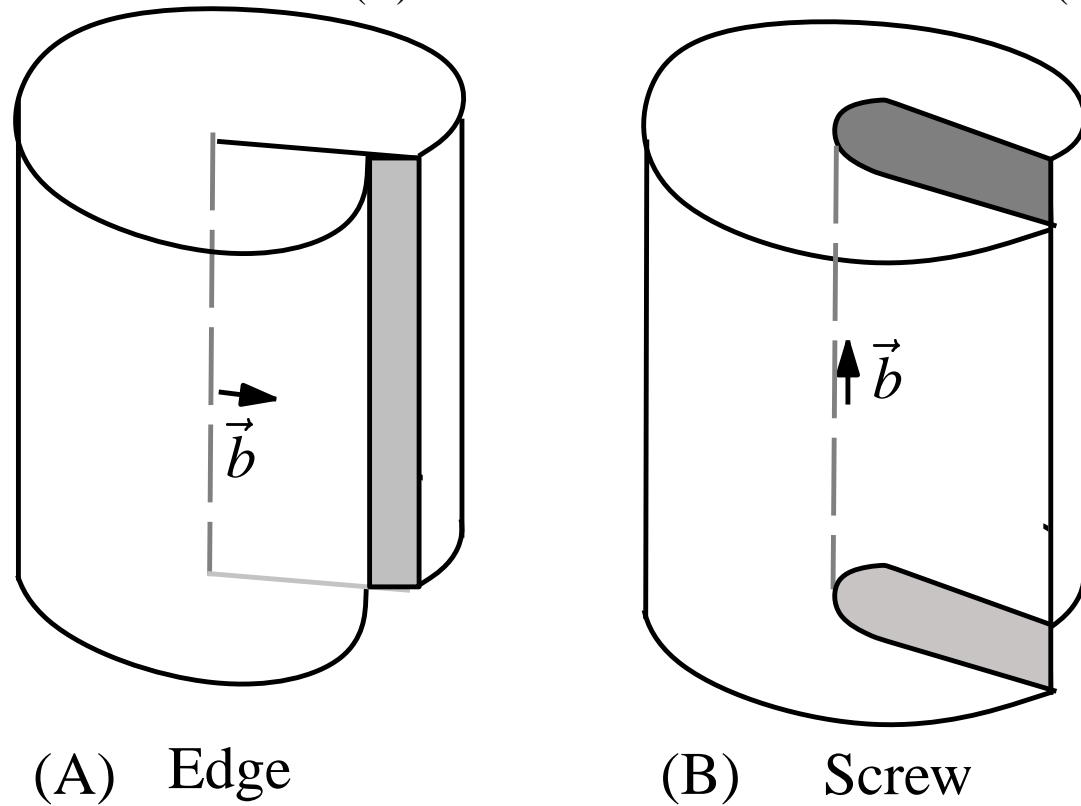
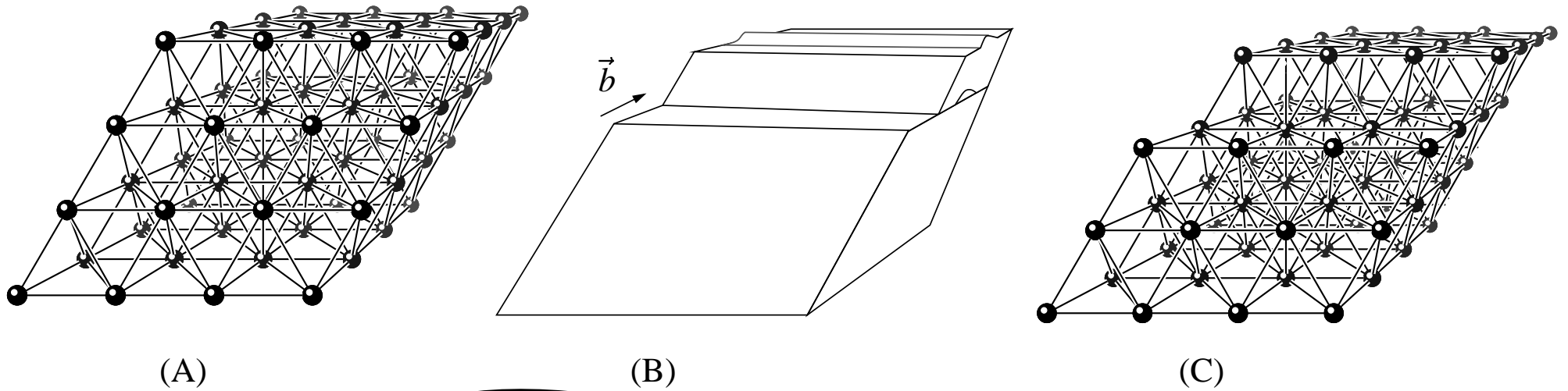
Material	Young's Modulus $Y/5$ ( $10^{11}$ ergs $\text{cm}^{-3}$ )	Theoretical Strength ( $10^{11}$ ergs $\text{cm}^{-3}$ )	Practical Strength ( $10^{11}$ ergs $\text{cm}^{-3}$ )	Ratio
Iron	4.0	4	0.03	0.008
Titanium	2.2	3.1	0.03	0.009
Silicon	3.2	1.5	0.07	0.05
Glass	1.4	4	0.04	0.01

# Complete Cohesive Energy Curve

$$-P = \frac{1}{4\pi r_W^2} \frac{\partial \mathcal{E}(r_W)}{\partial r_W}. \quad (\text{L6})$$

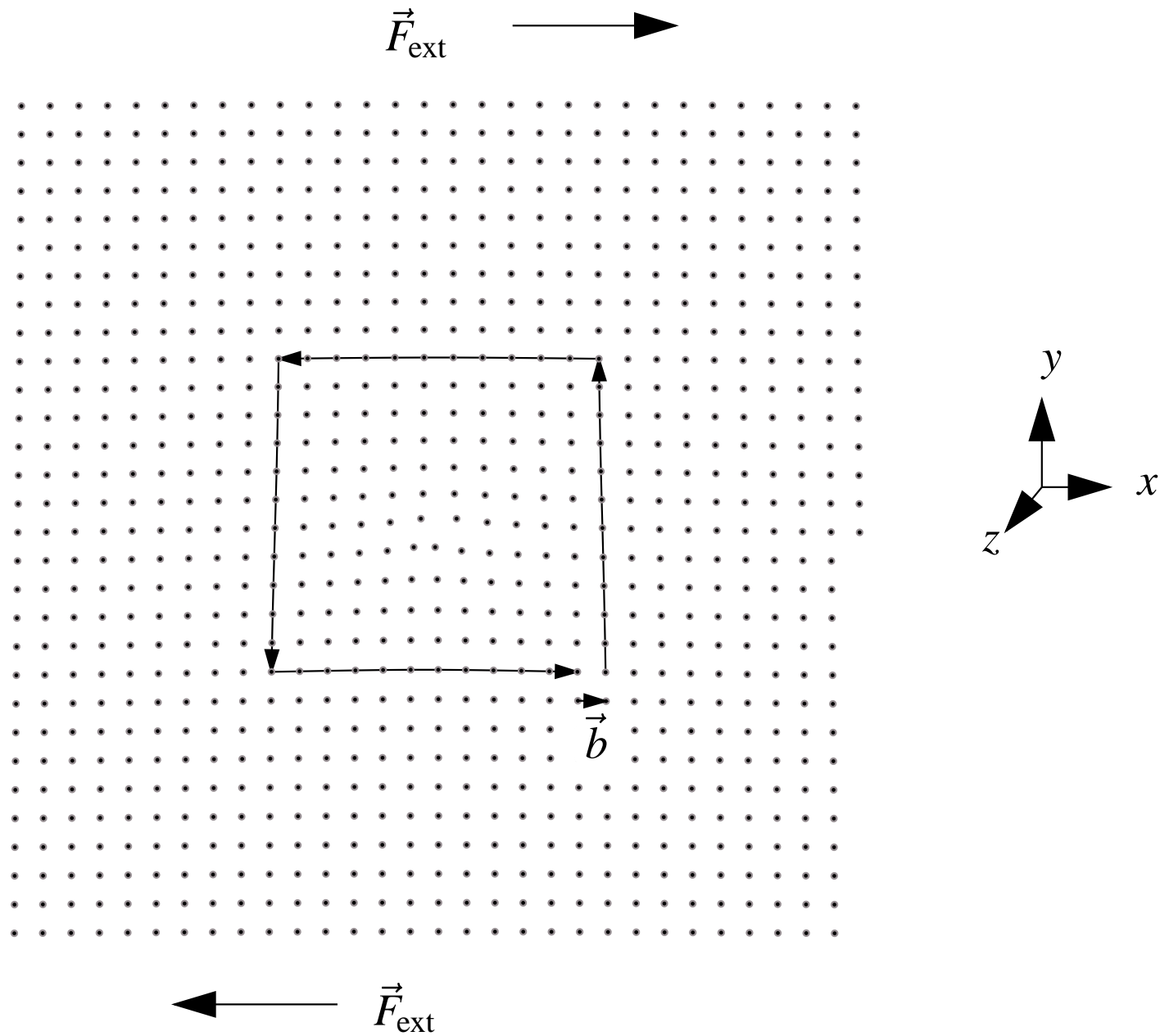


# Dislocations

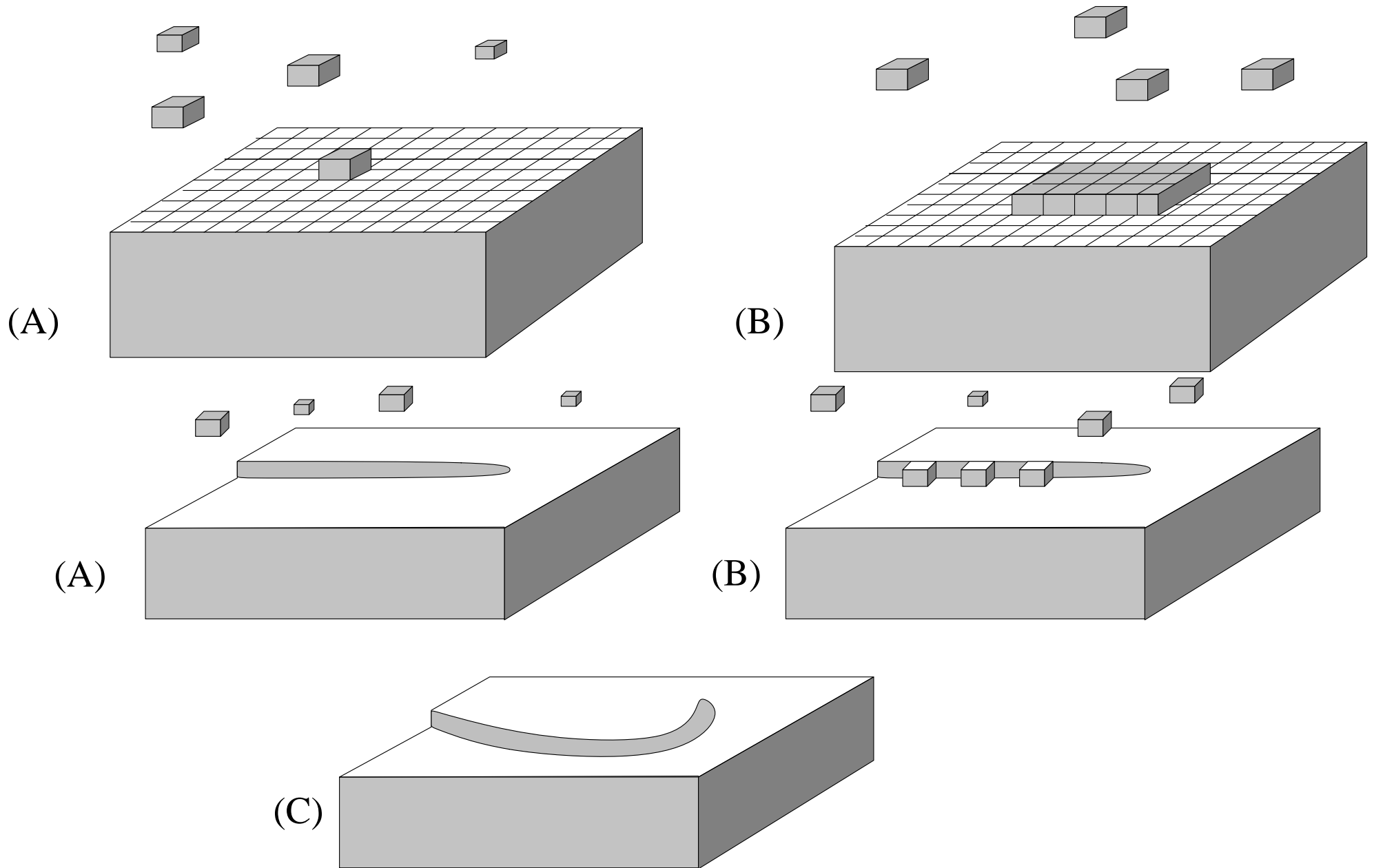




# Burgers Vector



# Experimental Observations of Dislocations 10



# Experimental Observations of Dislocations 11

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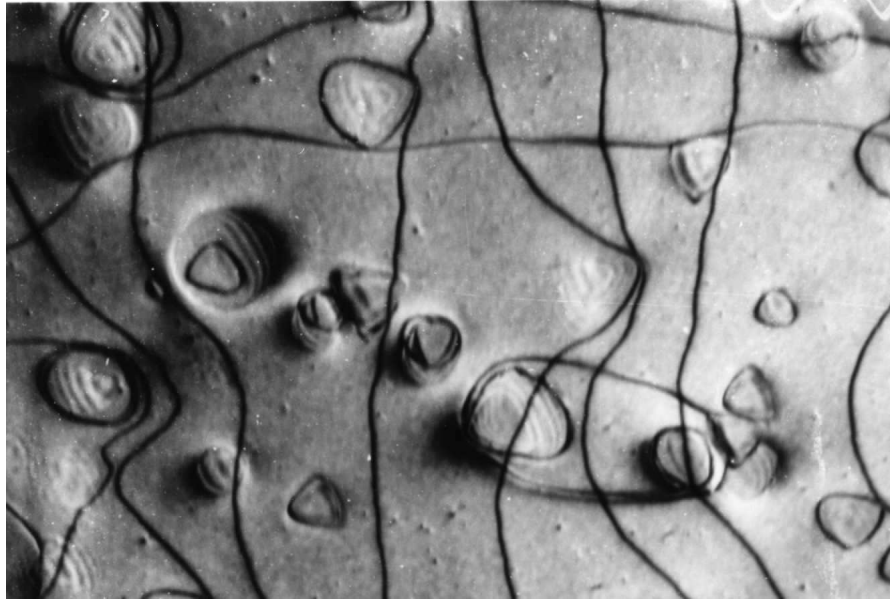


[Source: [Amelinckx \(1964\)](#)]

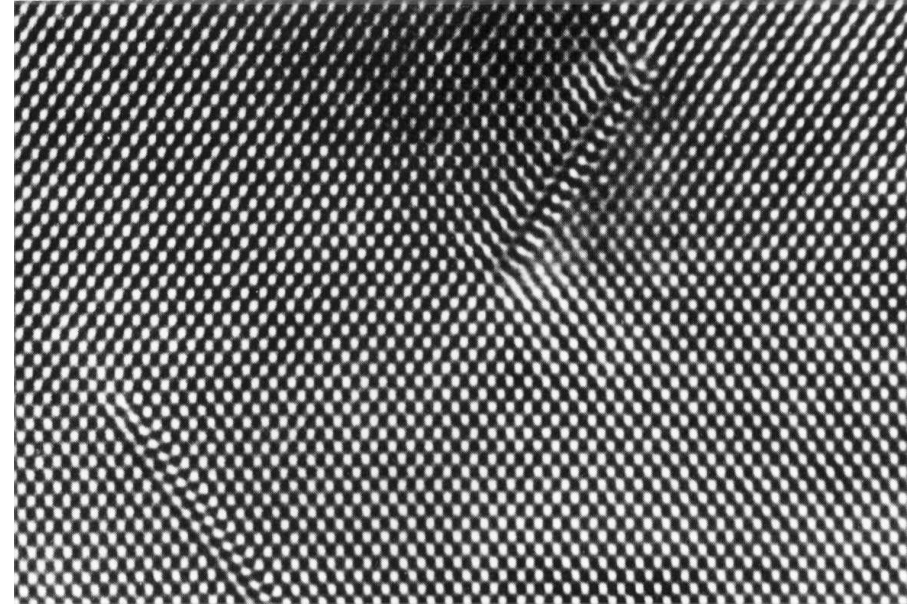
# Experimental Observations of Dislocations 12

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(A)



(B)

(A) Courtesy of J. Humphreys, Manchester University.)

[(B) Cullis et al. (1985)]

$$f_x = \sigma_{xy} b_x, \quad (\text{L7})$$

$$\sigma_{xy} = \frac{F_{\text{ext}}}{Na^2} \quad (\text{L8})$$

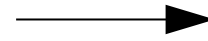
$$\vec{f} = (\sigma \cdot \vec{b}) \times \hat{L}. \quad (\text{L9})$$

Peach–Kohler force

# One-Dimensional Dislocations: Frenkel–Kontorova Model

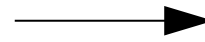
14

$$F = F_{\text{crit}}/2$$



(A)

$$F = F_{\text{crit}}$$



(B)

$-5a$   $-4a$   $-3a$   $-2a$   $-a$   $0$   $a$   $2a$   $3a$   $4a$   $5a$

Find force needed to move dislocation in simple one–dimensional model.

21st April 2003

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# One-Dimensional Dislocations: Frenkel–Kontorova Model

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$$U(x) = \frac{1}{2}\mathcal{K}\left[x - a \operatorname{int}\left(x/a + \frac{1}{2}\right)\right]^2 - fx, \quad (\text{L10})$$

$$f_n = k[x_{n+1} - x_n - a] + k[x_{n-1} - x_n + a] - \frac{\partial U}{\partial x}. \quad (\text{L11})$$

$$f_n = \begin{cases} k[x_{n+1} - x_n - a] + k[x_{n-1} - x_n + a] + f - \mathcal{K}[x_n - (n-1)a] & \text{for } n \leq 0 \\ k[x_{n+1} - x_n - a] + k[x_{n-1} - x_n + a] + f - \mathcal{K}[x_n - na] & \text{for } n > 0. \end{cases} \quad (\text{L12})$$

$$x_n = f/\mathcal{K} + a(n-1) + A_l e^{qn}, \quad (\text{L13})$$

$$k(e^q - 2 + e^{-q}) - \mathcal{K} = 0 \quad (\text{L14})$$

$$x_n = f/\mathcal{K} + an + A_r e^{-qn}. \quad (\text{L15})$$

# One-Dimensional Dislocations: Frenkel–Kontorova Model

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$$-a + A_l = A_r \quad (\text{L16a})$$

$$A_l e^q = a + A_r e^{-q}, \quad (\text{L16b})$$

$$A_l = \frac{a}{e^q + 1} \quad (\text{L17a})$$

$$A_r = \frac{-a}{e^{-q} + 1}. \quad (\text{L17b})$$

$$x_0 = -\frac{a}{2} = \frac{f_c}{\mathcal{K}} - a + A_l \quad (\text{L18})$$

$$\Rightarrow f_c = \frac{a\mathcal{K}}{2} \tanh \frac{q}{2}. \quad (\text{L19})$$

$$q \approx \sqrt{\frac{\mathcal{K}}{k}} \quad (\text{L20})$$

# One-Dimensional Dislocations: Frenkel–Kontorova Model

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$$f_c \approx \frac{a\mathcal{K}}{4} \sqrt{\frac{\mathcal{K}}{k}}. \quad (\text{L21})$$

# Impossibility of Crystalline Order in Two Dimensions

18

Peierls and Landau showed that two-dimensional crystals are destroyed by thermal fluctuations.

$$U = \int d^2r \frac{1}{2} C \sum_{\alpha\beta} \frac{\partial u_\alpha}{\partial r_\beta} \frac{\partial u_\alpha}{\partial r_\beta}. \quad (\text{L22})$$

$$u_\alpha(\vec{r}) = \sum_{\vec{k}} e^{i\vec{r}\cdot\vec{k}} u_\alpha(\vec{k}). \quad (\text{L23})$$

$$\vec{u}(\vec{k}) = 0 \text{ for } k > 1/\mathcal{D}. \quad (\text{L24})$$

$$U = \int d^2r \frac{1}{2} C \sum_{\beta\alpha\vec{k}\vec{k}'} k_\beta k'_\beta e^{i(\vec{k}-\vec{k}')\cdot\vec{r}} u_\alpha(\vec{k}) u_\alpha^*(\vec{k}') \quad (\text{L25})$$

$$= \frac{\mathcal{V}C}{2} \sum_{\alpha\vec{k}} k^2 |u_\alpha(\vec{k})|^2 \quad (\text{L26})$$

# Impossibility of Crystalline Order in Two Dimensions

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$$\langle u^2 \rangle = \left\langle \int \frac{d^2r}{V} \sum_{\beta} u_{\beta}(\vec{r}) u_{\beta}(\vec{r}) \right\rangle \quad (\text{L27})$$

$$= \sum_{\beta \vec{k}} \langle |u_{\beta}(\vec{k})|^2 \rangle \quad (\text{L28})$$

$$\vec{u}(\vec{k}) = \vec{u}^*(-\vec{k}) \quad (\text{L29})$$

$$\begin{aligned} & \langle |u_{\beta}(\vec{k})|^2 \rangle \\ = & \frac{\int \prod_{\alpha \vec{k}'} du_{\alpha}(\vec{k}') |u_{\beta}(\vec{k})|^2 e^{-\beta \frac{V_C}{2} \sum_{\alpha \vec{k}'} k'^2 |u_{\alpha}(\vec{k}')|^2}}{\int \prod_{\alpha \vec{k}'} du_{\alpha}(\vec{k}') e^{-\beta \frac{V_C}{2} \sum_{\alpha \vec{k}'} k'^2 |u_{\alpha}(\vec{k}')|^2}} \end{aligned} \quad (\text{L30})$$

$$= \frac{\int du_{\beta}(\vec{k}) |u_{\beta}(\vec{k})|^2 e^{-\beta V_C k^2 |u_{\beta}(\vec{k})|^2}}{\int du_{\beta}(\vec{k}) e^{-\beta V_C k^2 |u_{\beta}(\vec{k})|^2}} \quad (\text{L31})$$

$$= \frac{\int du^r du^i [(u^r)^2 + (u^i)^2] e^{-\beta V_C k^2 [(u^r)^2 + (u^i)^2]}}{\int du^r du^i e^{-\beta V_C k^2 [(u^r)^2 + (u^i)^2]}} \quad (\text{L32})$$

# Impossibility of Crystalline Order in Two Dimensions

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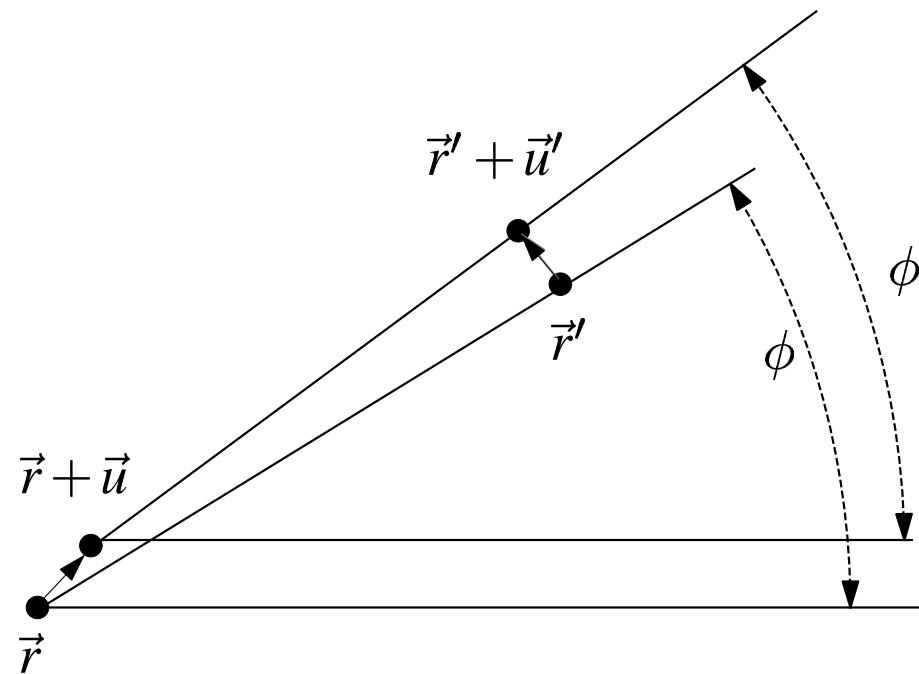
$$= \frac{k_B T}{\mathcal{V} C k^2}. \quad (\text{L33})$$

$$\langle u^2 \rangle = \sum_{\alpha \vec{k}} \frac{k_B T}{\mathcal{V} C k^2} \quad (\text{L34})$$

$$= 2 \int \frac{d^2 k}{(2\pi)^2} \frac{k_B T}{C k^2} \quad (\text{L35})$$

$$= 2 \int_0^{1/\mathcal{D}} \frac{dk}{2\pi k} \frac{k_B T}{C} \rightarrow \infty. \quad (\text{L36})$$





$$(dx, dy) = \vec{r}' - \vec{r}. \quad (\text{L37})$$

$$\phi = \tan^{-1}(dy/dx). \quad (\text{L38})$$

$$\vec{r} + \vec{u}(\vec{r}) \quad \text{and} \quad \vec{r}' + \vec{u}(\vec{r}'). \quad (\text{L39})$$

$$\phi' = \tan^{-1} \left( \frac{dy + \partial u_y / \partial x dx + \partial u_y / \partial y dy}{dx + \partial u_x / \partial x dx + \partial u_x / \partial y dy} \right) \quad (\text{L40})$$

$$\approx \phi + \frac{dx dy}{dx^2 + dy^2} \left[ \left( \frac{\partial u_y}{\partial y} - \frac{\partial u_x}{\partial x} \right) + \frac{dx}{dy} \frac{\partial u_y}{\partial x} - \frac{dy}{dx} \frac{\partial u_x}{\partial y} \right] \quad (\text{L41})$$

$$\Rightarrow \phi' - \phi = \cos \phi \sin \phi \left( \frac{\partial u_y}{\partial y} - \frac{\partial u_x}{\partial x} \right) + \cos^2 \phi \frac{\partial u_y}{\partial x} - \sin^2 \phi \frac{\partial u_x}{\partial y}. \quad (\text{L42})$$

$$\delta \phi(\vec{r}) = \frac{1}{2} \left( \frac{\partial u_y}{\partial x} - \frac{\partial u_x}{\partial y} \right). \quad (\text{L43})$$

$$\delta \phi(\vec{r}) = \frac{1}{2} \sum_{\vec{k}} (ik_x u_y(\vec{k}) - ik_y u_x(\vec{k})) e^{i\vec{k} \cdot \vec{r}}. \quad (\text{L44})$$

$$\int \frac{d^2 r}{\mathcal{V}} \langle \delta \phi(\vec{r}) \delta \phi(\vec{r}) \rangle \quad (\text{L45})$$

$$= \frac{1}{4} \sum_{\vec{k}} k_x^2 \langle |u_x(\vec{k})|^2 \rangle + k_y^2 \langle |u_y(\vec{k})|^2 \rangle - k_x k_y \langle (u_x(\vec{k}) u_y^*(\vec{k}) + u_y(\vec{k}) u_x^*(\vec{k})) \rangle \quad (\text{L46})$$

$$= \frac{1}{4} \sum_{\vec{k}} \frac{k_B T}{C \mathcal{V} k^2} (k_x^2 + k_y^2) \quad (\text{L47})$$

$$= \frac{k_B T}{4C} \int_0^{2\pi} d\theta \int_0^{1/\mathcal{D}} \frac{dk k}{(2\pi)^2} = \frac{k_B T}{16\pi \mathcal{D}^2 C}. \quad (\text{L48})$$

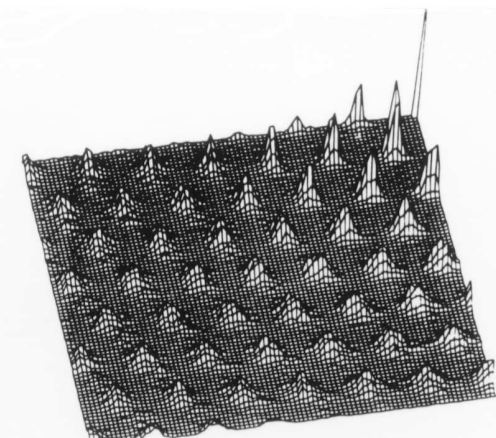
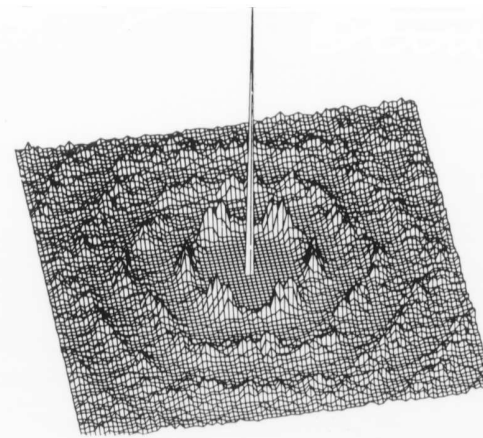
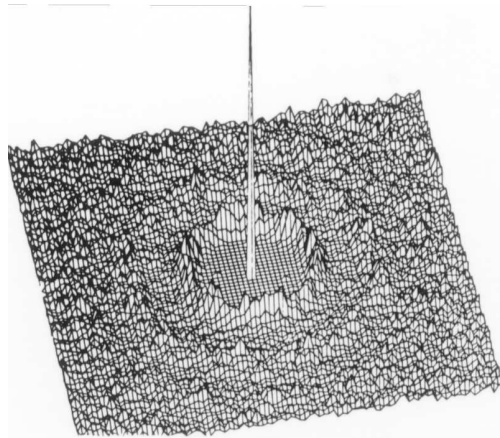
# Kosterlitz–Thouless–Berezinskii Transition <sup>24</sup>

Liquid

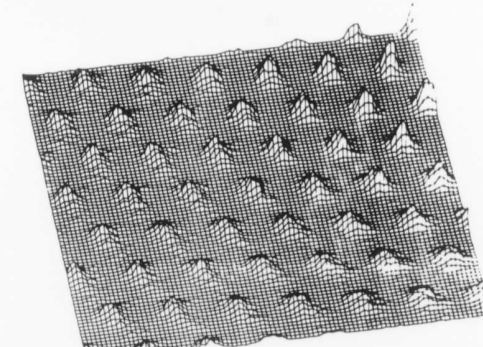
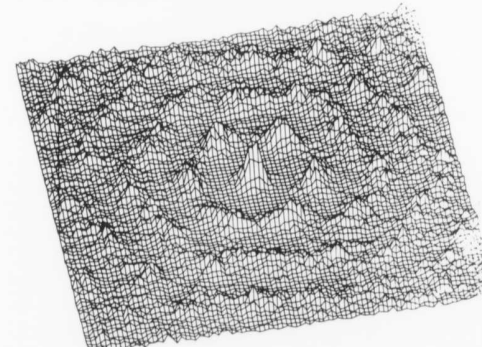
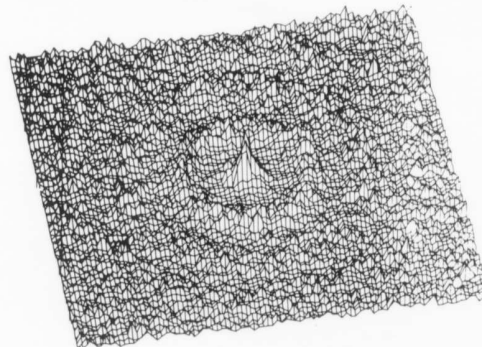
Hexatic

Crystal

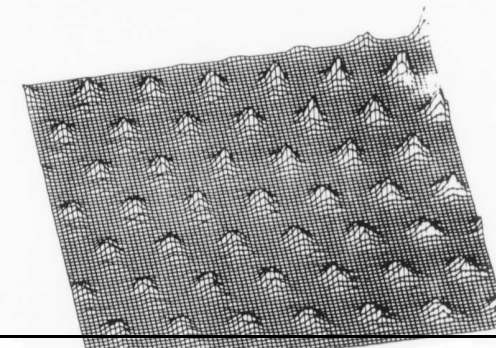
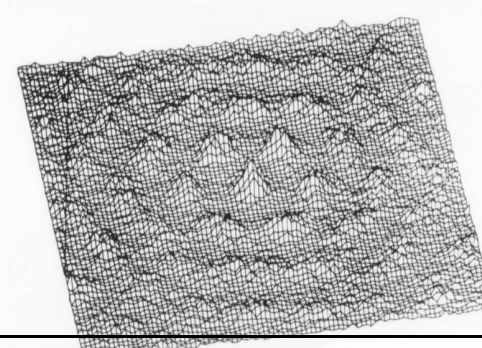
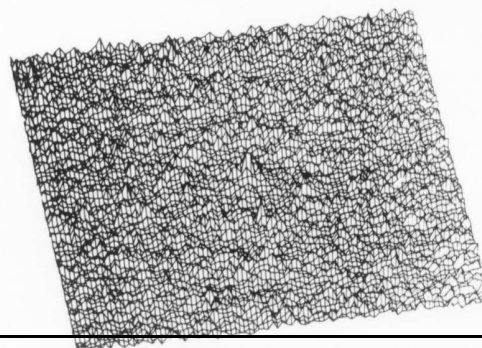
$t = 0$  s



$t = 0.05$  s



$t = 0.1$  s



21st April 2003

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# Kosterlitz–Thouless–Berezinskii Transition<sup>25</sup>

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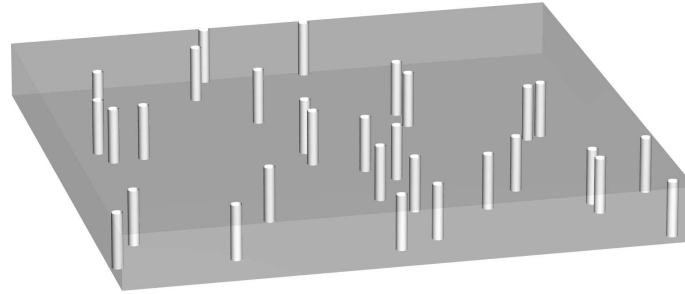
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[Murray and Grier (1996)]

# Kosterlitz–Thouless–Berezinskii Transition <sup>26</sup>

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$$u_x = 0, \quad u_y = 0, \quad u(x, y) = u_z(x, y). \quad (\text{L49})$$

$$U = \frac{a\mu}{2} \int d^2r (\nabla u)^2 \quad (\text{L50})$$

$$\nabla^2 u = 0. \quad (\text{L51})$$

$$u(x, y) = \frac{a}{2\pi} \text{Im} \ln[x + iy]. \quad (\text{L52})$$

$$U = \frac{a\mu}{2} \left( \frac{a}{2\pi} \right)^2 \int d^2r \left[ \frac{-y}{x^2 + y^2} \right]^2 + \left[ \frac{x}{x^2 + y^2} \right]^2 \quad (\text{L53})$$



# Kosterlitz–Thouless–Berezinskii Transition 27

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$$= \frac{a\mu}{2} \left(\frac{a}{2\pi}\right)^2 \int_a^R dr 2\pi r \frac{1}{r^2} \quad (\text{L54})$$

$$\rightarrow \frac{1}{4\pi} (a^3 \mu) \ln\left(\frac{R}{a}\right) + w. \quad (\text{L55})$$

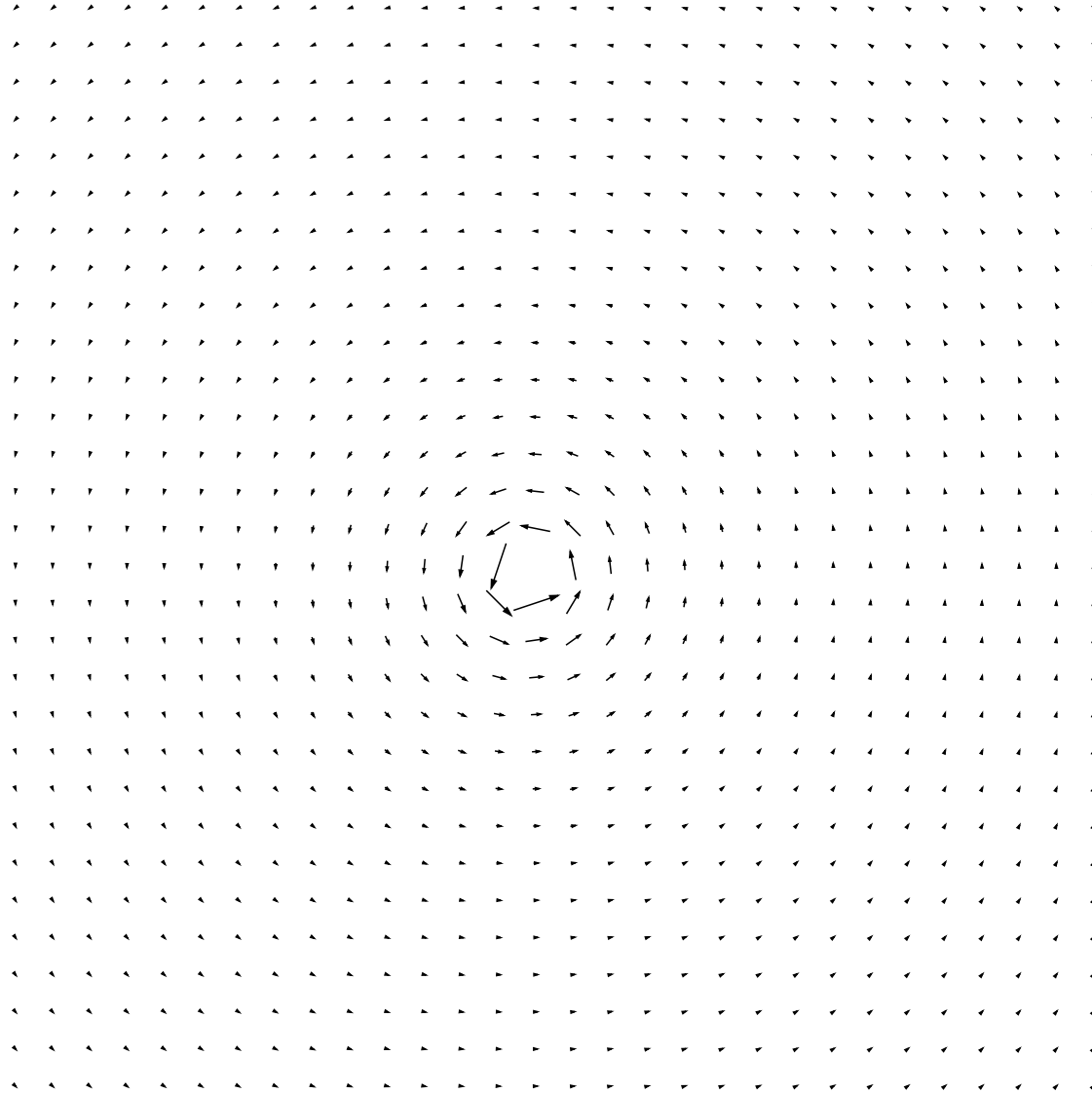
$$u(x, y) = \frac{a}{2\pi} \text{Im} \{ \ln[x + iy] - \ln[x - x_0 + iy] \}. \quad (\text{L56})$$

$$2q^2 \ln\left(\frac{x_0}{a}\right) + 2w \quad \text{with } q^2 = \frac{a^3 \mu}{4\pi}. \quad (\text{L57})$$

# Kosterlitz–Thouless–Berezinskii Transition <sup>28</sup>

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$$S = 2k_B \ln(L/a),$$

(L58)

# Kosterlitz–Thouless–Berezinskii Transition 29

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$$\mathcal{E} = q^2 \ln(L/a). \quad (\text{L59})$$

$$k_B T_c = \frac{q^2}{2}. \quad (\text{L60})$$

$$\mathcal{H} = \frac{1}{2} \sum_{i \neq j} U(|\vec{r}_i - \vec{r}_j|) + 2w \quad \text{with} \quad U(r) = 2q^2 \ln(r/a) \quad (\text{L61})$$

$$\langle r^2 \rangle = \frac{\int_a^\infty dr 2\pi r r^2 e^{-\beta U(r)}}{\int_a^\infty dr 2\pi r e^{-\beta U(r)}} \quad (\text{L62})$$

$$= a^2 \left[ \frac{\beta q^2 - 1}{\beta q^2 - 2} \right]. \quad (\text{L63})$$

$$Z_{\text{gr}} = 1 + \sum_{\vec{r}_1 \vec{r}_2} e^{-\beta U(|\vec{r}_1 - \vec{r}_2|) - 2\beta w} + \dots \quad (\text{L64})$$

# Kosterlitz–Thouless–Berezinskii Transition 30

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$$n(r)dr = \frac{dr}{R^2} \langle \delta_{r,|\vec{r}_1-\vec{r}_2|} \rangle = \frac{dr}{R^2} \sum_{\vec{r}_1\vec{r}_2} e^{-\beta U(|\vec{r}_1-\vec{r}_2|)-2\beta w} \delta_{r,|\vec{r}_1-\vec{r}_2|} + \dots \quad (\text{L65})$$

$$\approx \frac{1}{a^2} \frac{2\pi r^2 dr}{a^2} e^{-\beta U(r)-2\beta w}. \quad (\text{L66})$$

$$\vec{p} = \alpha \vec{E} = rq \langle (\cos \theta, \sin \theta) \rangle \quad (\text{L67})$$

$$= \int \frac{d\theta}{2\pi} e^{-\beta U(r)-2\beta w + \beta Eqr \cos \theta} rq (\cos \theta, \sin \theta) \quad (\text{L68})$$

$$= \frac{1}{2} \beta q^2 r^2 \vec{E}. \quad (\text{L69})$$

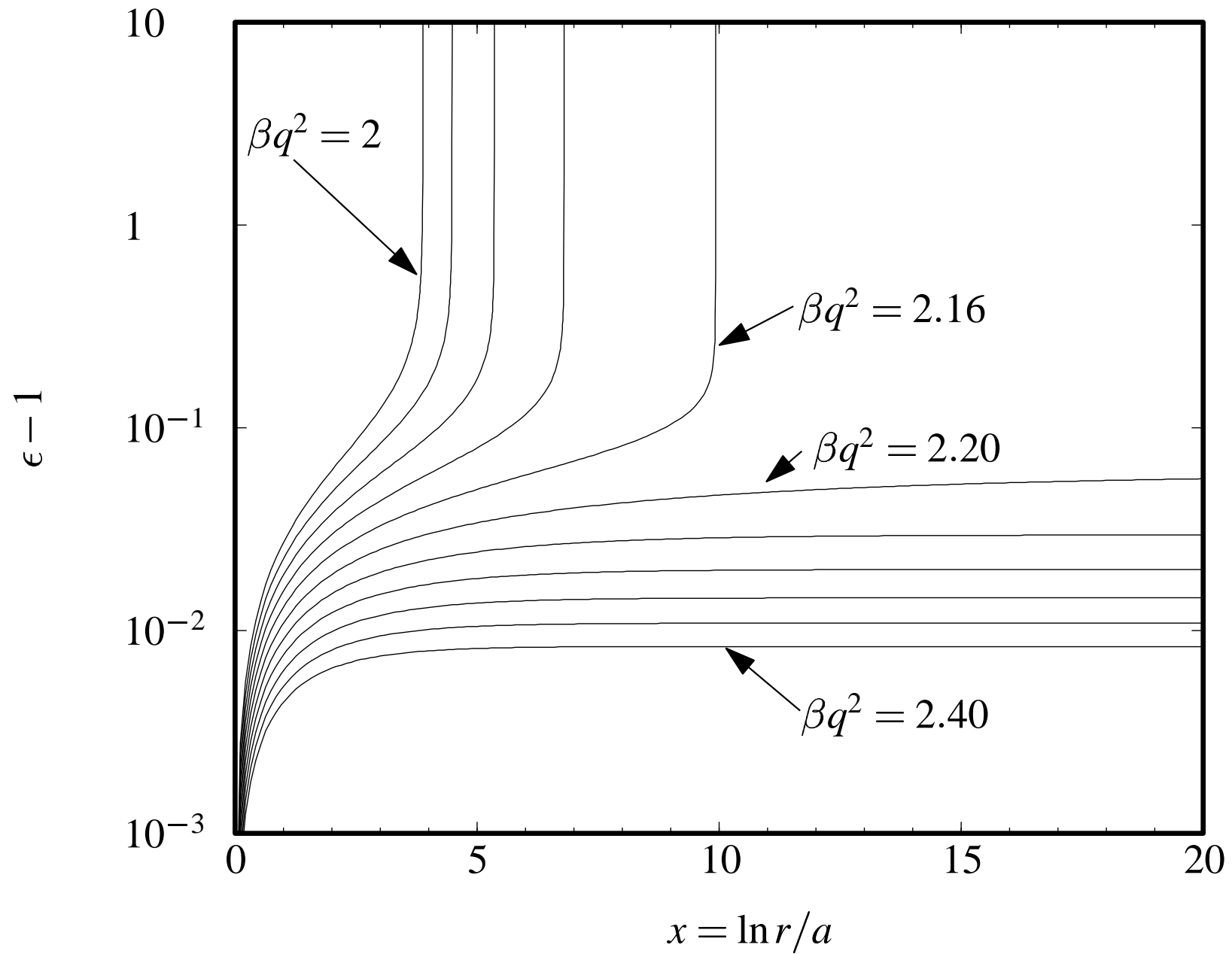
$$d\chi(r) = n(r) dr \alpha(r) = \frac{1}{2} \beta q^2 \left(\frac{r}{a}\right)^2 \frac{2\pi r dr}{a^2} e^{-\beta U/\epsilon(r)-2\beta w}. \quad (\text{L70})$$

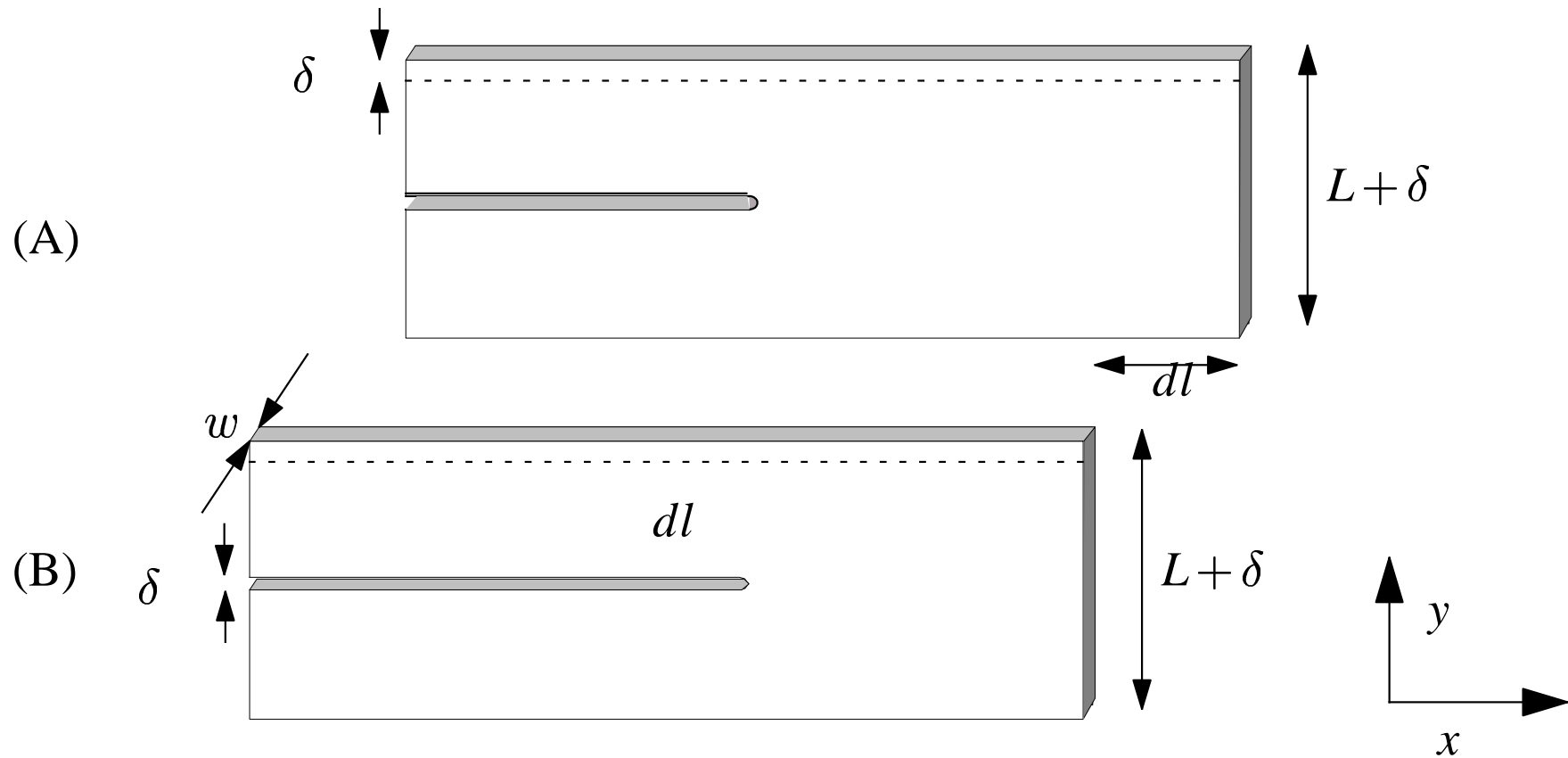
$$\epsilon(r) = 1 + 4\pi \int_a^r d\chi = 1 + 4\pi \int_a^r dr' n(r') \alpha(r') \quad (\text{L71})$$

$$\Rightarrow \frac{d\epsilon(r)}{dr} = 4\pi^2 \beta q^2 \frac{r^3}{a^4} e^{-\beta U(r)/\epsilon(r) - 2\beta w} \quad (\text{L72})$$

$$\Rightarrow \frac{d\epsilon(x)}{dx} = 4\pi^2 \beta q^2 x^{3-2\beta q^2/\epsilon(x)} e^{-2\beta w}. \quad (\text{L73})$$

# Kosterlitz–Thouless–Berezinskii Transition <sup>32</sup>





$$U = \frac{1}{2} \delta^2 w \frac{Y}{L}, \quad (\text{L74})$$

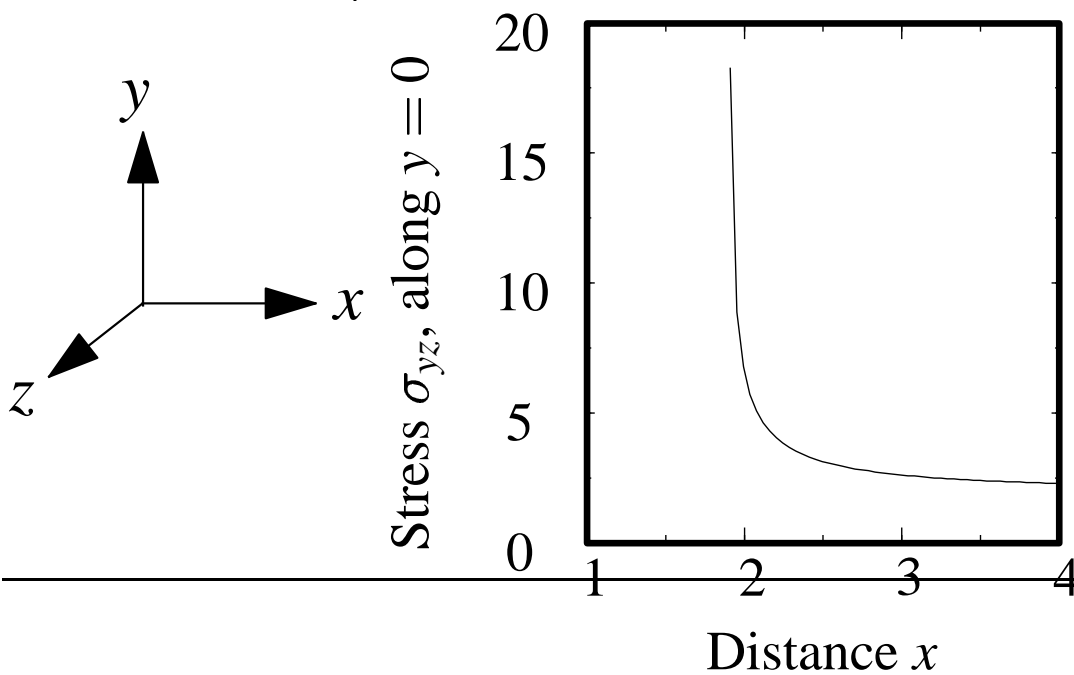
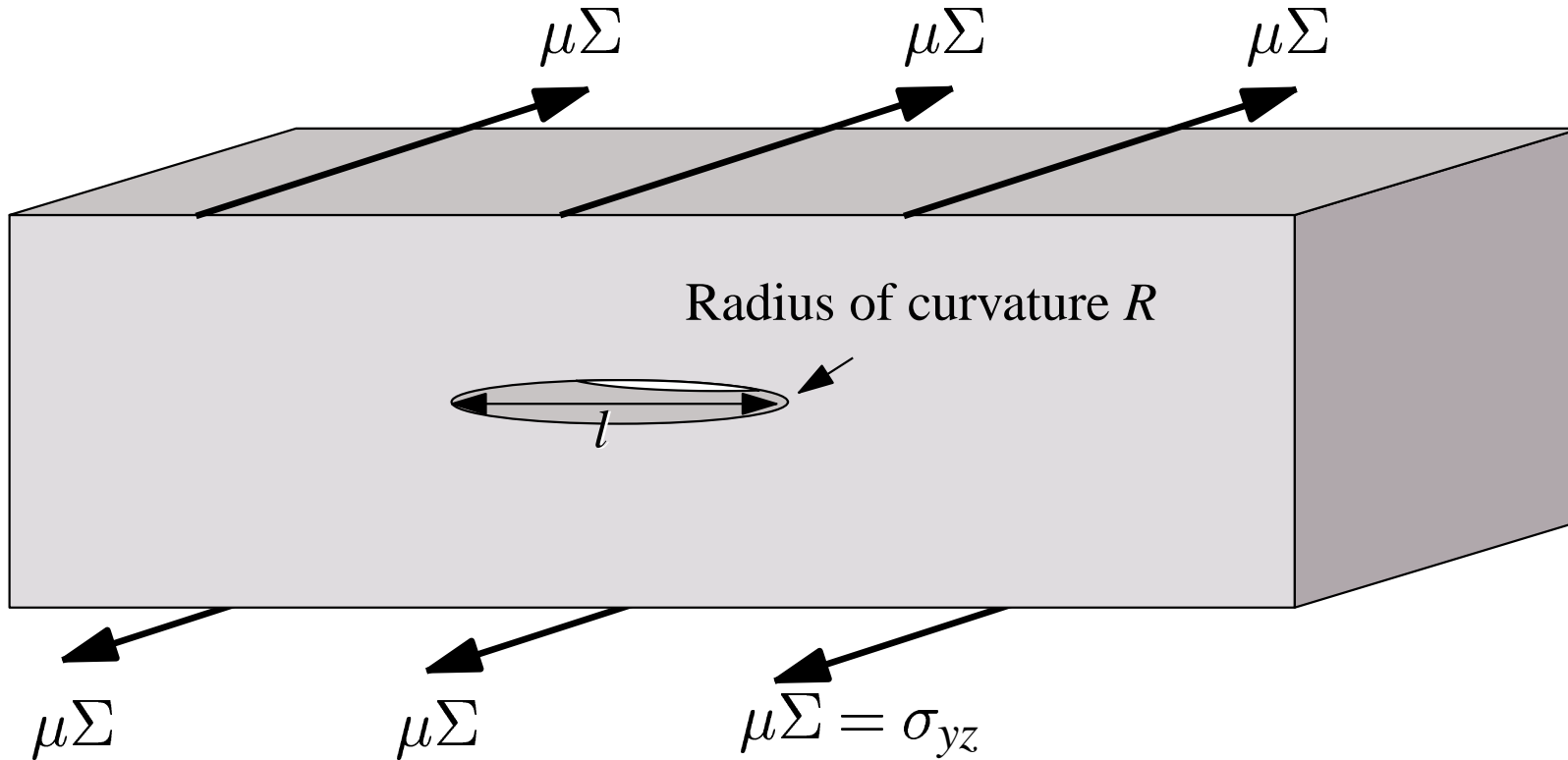
$$dU = dl \frac{1}{2} \delta^2 w \frac{Y}{L}. \quad (\text{L75})$$

$$\Gamma w dl = dl \frac{1}{2} \delta^2 w \frac{Y}{L} \quad (\text{L76})$$

$$\Rightarrow \delta = \sqrt{\frac{2\Gamma L}{Y}} \quad \text{and} \quad \sigma_{yy} = Y \frac{\delta}{L} = \sqrt{\frac{2\Gamma Y}{L}}. \quad (\text{L77})$$



# Fracture of a Strip



$$\frac{\text{Maximum stress}}{\text{applied stress}} \propto \sqrt{\frac{l}{R}}, \quad (\text{L78})$$

$$\nabla^2 u = 0. \quad (\text{L79})$$

$$u = \frac{\phi(\zeta) + \overline{\phi(\zeta)}}{2}, \quad (\text{L80})$$

$$\sigma_{yz} = \mu \frac{\partial u}{\partial y} = \frac{\mu}{2} [i\phi'(x+iy) - \overline{i\phi'(x+iy)}] \quad (\text{L81})$$

$$\Rightarrow \phi'(\zeta) \rightarrow -i\Sigma \text{ for } \zeta \rightarrow \infty. \quad (\text{L82})$$

$$(x(t), y(t)) \quad (\text{L83})$$

$$\vec{T} = \left( \frac{\partial x}{\partial t}, \frac{\partial y}{\partial t} \right) \quad \text{and} \quad \vec{N} = \left( -\frac{\partial y}{\partial t}, \frac{\partial x}{\partial t} \right) \quad (\text{L84})$$

$$(\sigma_{xz}, \sigma_{yz}) \cdot \vec{N} = 0 \quad (\text{L85})$$

$$\Rightarrow \mu \left( \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y} \right) \cdot \left( -\frac{\partial y}{\partial t}, \frac{\partial x}{\partial t} \right) = 0 \Rightarrow \frac{\partial u}{\partial y} \frac{\partial x}{\partial t} - \frac{\partial u}{\partial x} \frac{\partial y}{\partial t} = 0 \quad (\text{L86})$$

$$\Rightarrow \left( -\frac{\partial \phi}{\partial ix} + \frac{\partial \bar{\phi}}{\partial ix} \right) \frac{\partial x}{\partial t} = \left( \frac{\partial \phi}{\partial iy} - \frac{\partial \bar{\phi}}{\partial iy} \right) \frac{\partial y}{\partial t} \quad (\text{L87})$$

$$\Rightarrow \frac{\partial \phi}{\partial t} = \frac{\partial \bar{\phi}}{\partial t} \quad (\text{L88})$$

$$\Rightarrow \phi(\zeta) = \overline{\phi(\bar{\zeta})} \quad (\text{L89})$$

$$\zeta = \omega + \frac{p}{\omega}, \quad (\text{L90})$$

$$\omega = e^{i\theta}, \quad (\text{L91})$$

$$\phi(\omega) = \overline{\phi(\omega)} = \bar{\phi}\left(\frac{1}{\omega}\right), \quad (\text{L92})$$

$$\omega = \frac{\zeta + \sqrt{\zeta^2 - 4p}}{2}. \quad (\text{L93})$$

$$\phi(\omega) \rightarrow -i\Sigma\omega \quad \text{for } \omega \rightarrow \infty. \quad (\text{L94})$$

$$\bar{\phi}(1/\omega) \rightarrow -i\Sigma\omega \quad \text{for } \omega \rightarrow \infty, \quad (\text{L95})$$

$$\bar{\phi}(\omega) \rightarrow \frac{-i\Sigma}{\omega} \quad \text{for } \omega \rightarrow 0 \quad (\text{L96})$$

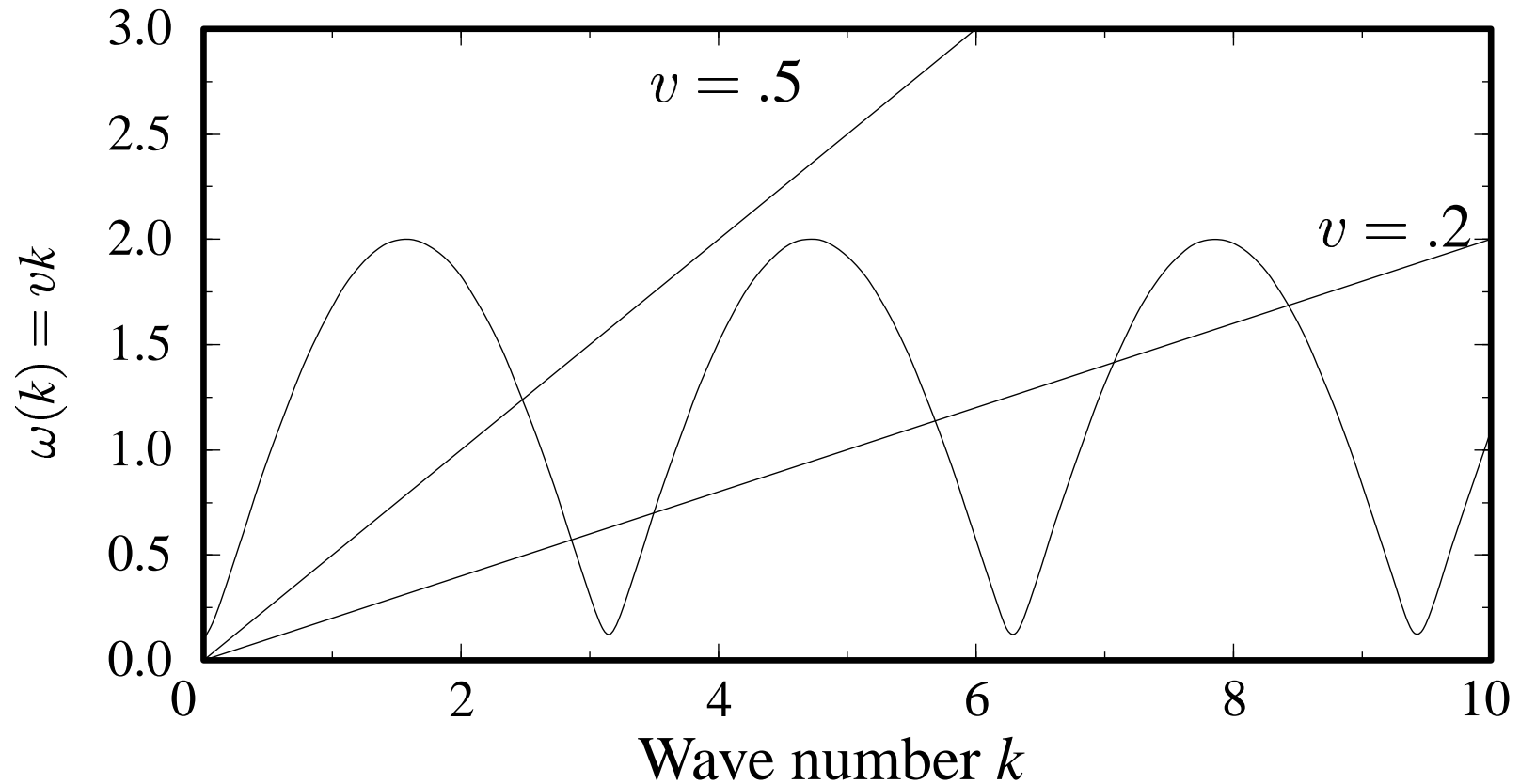
$$\phi(\omega) \rightarrow \frac{i\Sigma}{\omega} \quad \text{for } \omega \rightarrow 0. \quad (\text{L97})$$

$$\phi(\omega) = -i\Sigma\omega + i\frac{\Sigma}{\omega} \quad (\text{L98})$$

$$\Rightarrow \phi(\zeta) = -i\Sigma\frac{\zeta}{2}(1 + \sqrt{1 - 4p/\zeta^2}) + i\Sigma\frac{\zeta}{2p}(1 - \sqrt{1 - 4p/\zeta^2}). \quad (\text{L99})$$

$$\sigma_{yz} = \mu \frac{\partial u}{\partial y} = \frac{\mu \Sigma x}{\sqrt{x-2}\sqrt{x+2}} \rightarrow \frac{\mu \Sigma}{\sqrt{x-2}} \quad \text{as } x \rightarrow 2. \quad (\text{L100})$$

$$K = \lim_{r \rightarrow 0} \sqrt{2\pi r} \sigma_{yz}. \quad (\text{L101})$$



$$\vec{r} \rightarrow \vec{r} + \vec{R} \quad (\text{L102a})$$

$$t \rightarrow t + a/v. \quad (\text{L102b})$$

$$e^{i\vec{k}\cdot(\vec{r}+\vec{R})-i\omega(t+a/v)} = e^{i\vec{k}\cdot\vec{r}-i\omega t} \quad (\text{L103})$$

$$\Rightarrow e^{i\vec{k}\cdot\vec{R}-i\omega a/v} = 1 \quad (\text{L104})$$

$$\Rightarrow (\vec{k} + \vec{K}) \cdot \vec{v} = \omega(\vec{k}) \quad (\text{L105})$$

$$\Rightarrow \vec{k} \cdot \vec{v} = \omega(\vec{k}). \quad (\text{L106})$$